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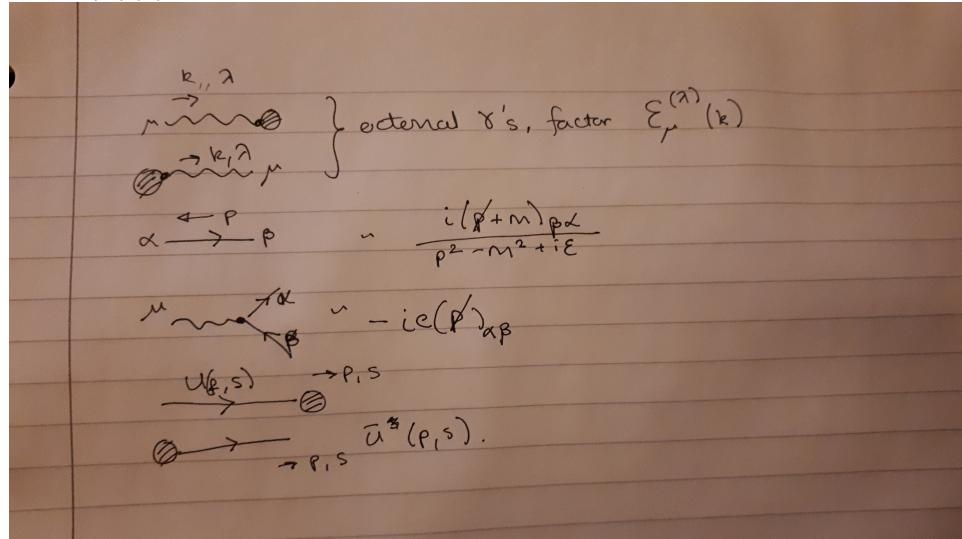
- (i) Write down the defining relation of the Clifford algebra of γ^μ in 1+3 dimensions.
- (ii) Write down the Dirac equation for a positive-frequency solution momentum space Dirac spinor $u(s, p^\mu)$ of spin s and 4-momentum p^μ .
- (iii) Write down the Lagrangian density for an electron e^- of mass m and charge e coupled to an electromagnetic field A_μ .
- (iv) Consider Compton scattering $\gamma(q, \epsilon_{in}(\lambda))e^-(p, s) \rightarrow \gamma(q', \epsilon_{out}(\lambda'))e^-(p', s')$ in Lorentz gauge.
 - (a) Write down/draw the Feynman rules needed for the calculation of the amplitude.
 - (b) Draw Feynman diagram(s) representing the total amplitude.
 - (c) Use the Feynman rules and diagrams to derive the total amplitude for Compton scattering.
 - (d) Calculate the modulus squared of the spin/polarisation averaged/summed amplitude in the limit $m \rightarrow 0$ in terms of the Mandelstam variables $s = (p+q)^2$ and $u = (p-q')^2$. Derive any properties of γ^μ you use from the Clifford algebra. You may assume that

$$\sum_s u(s, p)\bar{u}(s, p) = (\not{p} + m), \quad \sum_\lambda \epsilon_\mu^*(\lambda)\epsilon_\nu(\lambda) = -\eta_{\mu\nu}.$$

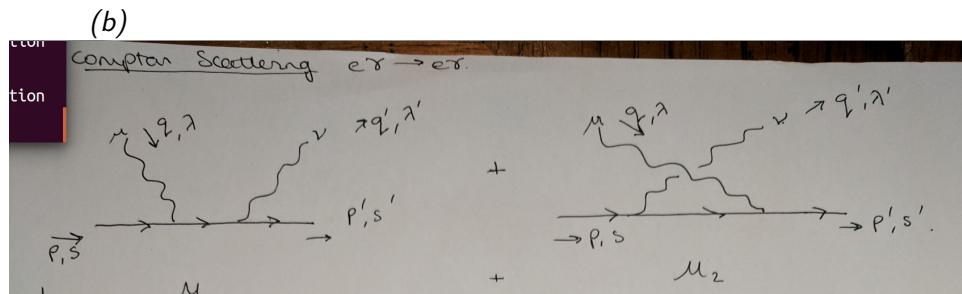
for a photon polarisation $\epsilon(\lambda)$.

Marks	Comments
[1]	(i) $\{\gamma^\mu, \gamma^\nu\} = 2\eta^{\mu\nu}I$
[1]	(ii) $(\not{p} - m)u(s, p^\mu) = 0$
[1]	(iii) $\mathcal{L} = \bar{\psi}(i\not{\partial} - m)\psi - \frac{1}{4}F_{\mu\nu}F^{\mu\nu} - e\bar{\psi}\not{A}\psi,$
[2]	where $F_{\mu\nu} = \partial_\mu A_\nu - \partial_\nu A_\mu$.

(iv)(a)



[3]



[2]

(c)

$$\mathcal{M} = \mathcal{M}_1 + \mathcal{M}_2 = (-ie)^2 \bar{u}(s', p') \epsilon_{out\nu}^*(\lambda') \gamma^\nu \frac{i(p+q)}{(p+q)^2} \epsilon_{in\mu}(\lambda) \gamma^\mu u(s, p) + \left[\begin{array}{l} q \leftrightarrow q' \\ \epsilon_{in} \leftrightarrow \epsilon_{out}^* \end{array} \right]$$

[1]

(d)

$$\begin{aligned} |\bar{\mathcal{M}}_1|^2 &= \frac{e^4}{4s^2} \sum_{\lambda, \lambda', s, s'} u(s', p') \bar{u}(s', p') \epsilon_{out}^*(\lambda') (p + q) \epsilon_{in}(\lambda) u(s, p) \bar{u}(s, p) \epsilon_{in}^*(\lambda) \mu (p + q) \epsilon_{out}(\lambda') \\ &= \frac{e^4}{4s^2} \text{Tr}[p' \gamma_\nu (p + q) \gamma_\mu p \gamma^\mu (p + q) \gamma^\nu] \end{aligned}$$

Now, we need

$$\gamma_\mu \not{a} \gamma^\mu = a^\rho \gamma_\mu (-\gamma^\mu \gamma_\rho + 2\delta_\rho^\mu) = a^\rho (-4\gamma_\rho + 2\gamma_\rho) = -2q.$$

$$\gamma_\mu \gamma^\nu \gamma^\rho \gamma^\mu = \gamma_\mu \gamma^\nu (-\gamma^\mu \gamma^\rho + 2\eta^{\rho\mu}) = 2\gamma^\nu \gamma^\rho + 2\gamma^\rho \gamma^\nu = 2\{\gamma^\nu, \gamma^\rho\} = 4\eta^{\nu\rho}.$$

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$$\begin{aligned}\gamma_\mu \gamma^\nu \gamma^\rho \gamma^\sigma \gamma^\mu &= \gamma_\mu \gamma^\nu \gamma^\rho (-\gamma^\mu \gamma^\sigma + 2\eta^{\mu\sigma}) = -4\eta^{\nu\rho} \gamma^\sigma + 2\gamma^\sigma \gamma^\nu \gamma^\rho \\ &= -4\eta^{\nu\rho} \gamma^\sigma + 2\gamma^\sigma (-\gamma^\rho \gamma^\nu + 2\eta^{\nu\rho}) = -2\gamma^\sigma \gamma^\rho \gamma^\nu.\end{aligned}$$

$\text{Tr}(\gamma^\mu \gamma^\nu) = \text{Tr}(\{\gamma^\mu, \gamma^\nu\})/2 = 4\eta^{\mu\nu}$. Thus

$$\begin{aligned}\text{Tr}(\gamma^\mu \gamma^\nu \gamma^\rho \gamma^\sigma) &= 2\eta^{\mu\nu} \text{Tr}(\gamma^\rho \gamma^\sigma) - \text{Tr}(\gamma^\nu \gamma^\mu \gamma^\rho \gamma^\sigma) = 8\eta^{\mu\nu} \eta^{\rho\sigma} - 2\eta^{\mu\rho} \text{Tr}(\gamma^\nu \gamma^\sigma) + \text{Tr}(\gamma^\nu \gamma^\rho \gamma^\mu \gamma^\sigma) \\ &= 8\eta^{\mu\nu} \eta^{\rho\sigma} - 8\eta^{\mu\rho} \eta^{\nu\sigma} + 2\eta^{\mu\sigma} \text{Tr}(\gamma^\nu \gamma^\rho) - \text{Tr}(\gamma^\nu \gamma^\rho \gamma^\sigma \gamma^\mu)\end{aligned}$$

Hence (by cyclicity of trace)

$$\text{Tr}(\gamma^\mu \gamma^\nu \gamma^\rho \gamma^\sigma) = 4\eta^{\mu\nu} \eta^{\rho\sigma} - 4\eta^{\mu\rho} \eta^{\nu\sigma} + 4\eta^{\mu\sigma} \eta^{\nu\rho}.$$

Also, using cyclicity of the trace,

$$\begin{aligned}|\bar{\mathcal{M}}_1|^2 &= \frac{e^4}{s^2} \text{Tr}[\not{p}'(\not{p} + \not{q})\not{p}(\not{p} + \not{q})], \quad \not{p}\not{p} = p^2 = 0, \Rightarrow |\bar{\mathcal{M}}_1|^2 = \frac{e^4}{s^2} \text{Tr}[\not{p}'\not{q}\not{p}\not{q}] \\ &= \frac{4e^4}{s^2} (2\not{p}' \cdot q \not{p} \cdot q) = -2e^4 s/u.\end{aligned}$$

For $|\bar{\mathcal{M}}_2|^2$, $s = (p + q)^2 \leftrightarrow (p - q')^2 = u$ so $|\bar{\mathcal{M}}_2|^2 = -2e^4 s/u$.

$$\begin{aligned}\mathcal{M}_1 \mathcal{M}_2^\dagger &= \frac{e^4}{2su} \text{Tr}[\not{p} \gamma^\mu (\not{p} + \not{q}) \gamma_\nu \not{p}' \gamma_\mu (\not{p} - \not{q}') \gamma^\nu] \\ &= -\frac{e^4}{2su} \text{Tr}[\gamma_\mu \not{p}' \not{p} \gamma^\mu (\not{p} + \not{q})(\not{p} - \not{q}')] \\ &= -\frac{2e^4}{su} \not{p} \cdot \not{p}' \text{Tr}[(\not{p} + \not{q})(\not{p} - \not{q}')]\end{aligned}$$

The trace is $4(p^2 - p \cdot q' + q \cdot p - q \cdot q')$ but $p = p' - q' - q$ so trace is $4(q \cdot (p' + q' - q) - p \cdot q' - q \cdot q') = 0$ in the massless limit since $q \cdot p' = p \cdot q'$.

Hence

$$|\bar{\mathcal{M}}|^2 = -2e^4 \left(\frac{u}{s} + \frac{s}{u} \right).$$

[10]

[20]