

G. K. Batchelor 8 March 1920 - 30 April 2000

A perspective of GKB's Micro-hydrodynamics

John Hinch

DAMTP, Cambridge

September 10, 2010

Understand, explain and exploit Kolmogorov

- Understand, explain and exploit Kolmogorov
- Townsend's experiments

- Understand, explain and exploit Kolmogorov
- Townsend's experiments
- Troubled cannot solve Navier-Stokes

- Understand, explain and exploit Kolmogorov
- Townsend's experiments
- Troubled cannot solve Navier-Stokes
- Marseille 1961 IUTAM/IUGG Congress
 - Stewart's 3 decades of $k^{-5/3}$ at $Re = 10^8$
 - Kolmogorov: ERROR (intermittency)

- Understand, explain and exploit Kolmogorov
- Townsend's experiments
- Troubled cannot solve Navier-Stokes
- Marseille 1961 IUTAM/IUGG Congress
 - Stewart's 3 decades of $k^{-5/3}$ at $Re = 10^8$
 - Kolmogorov: ERROR (intermittency)
- ► Time off: JFM, Collected papers of G.I.Taylor, textbook

- Understand, explain and exploit Kolmogorov
- Townsend's experiments
- Troubled cannot solve Navier-Stokes
- Marseille 1961 IUTAM/IUGG Congress
 - Stewart's 3 decades of $k^{-5/3}$ at $Re = 10^8$
 - Kolmogorov: ERROR (intermittency)
- ► Time off: JFM, Collected papers of G.I.Taylor, textbook
- 1967 start of second wave of research (no more turbulence) in low-Reynolds-number suspensions of particles producing 8 of his top 10 most cited papers

Before



What followed

10

- Before
 - Stokes $6\pi\mu aV$

Batchelor's research

Before

- Stokes $6\pi\mu aV$
- Einstein viscosity $\mu(1+\frac{5}{2}c)$, diffusivity $D = kT/6\pi\mu a$

Batchelor's research

Before

- Stokes $6\pi\mu aV$
- Einstein viscosity $\mu(1+\frac{5}{2}c)$, diffusivity $D = kT/6\pi\mu a$
- GI drops and viscosity of emulsion

Batchelor's research

Before

- Stokes $6\pi\mu aV$
- Einstein viscosity $\mu(1+\frac{5}{2}c)$, diffusivity $D = kT/6\pi\mu a$
- GI drops and viscosity of emulsion
- Brenner, Cox, Mason, Giesekus, Saffman, Bretherton
- Batchelor's research

Before

- Stokes 6πμaV
- Einstein viscosity $\mu(1+\frac{5}{2}c)$, diffusivity $D = kT/6\pi\mu a$
- GI drops and viscosity of emulsion
- Brenner, Cox, Mason, Giesekus, Saffman, Bretherton
- textbook and GI collected papers
- Batchelor's research

Before

- ► Stokes 6πµaV
- Einstein viscosity $\mu(1+\frac{5}{2}c)$, diffusivity $D = kT/6\pi\mu a$
- GI drops and viscosity of emulsion
- Brenner, Cox, Mason, Giesekus, Saffman, Bretherton
- textbook and GI collected papers \rightarrow opportunities
- Batchelor's research

Before

- Stokes 6πμaV
- Einstein viscosity $\mu(1+\frac{5}{2}c)$, diffusivity $D = kT/6\pi\mu a$
- GI drops and viscosity of emulsion
- Brenner, Cox, Mason, Giesekus, Saffman, Bretherton
- textbook and GI collected papers \rightarrow opportunities
- Batchelor's research
 - Hydrodynamic corrections to Stokes & Einstein

Before

- Stokes 6πμaV
- Einstein viscosity $\mu(1+\frac{5}{2}c)$, diffusivity $D = kT/6\pi\mu a$
- GI drops and viscosity of emulsion
- Brenner, Cox, Mason, Giesekus, Saffman, Bretherton
- textbook and GI collected papers \rightarrow opportunities
- Batchelor's research
 - Hydrodynamic corrections to Stokes & Einstein
 - To start:
 - (i) bulk stress (ii) μ at c^2
- What followed

A start

Stress system in a suspension of force-free particles JFM 1970

Suspension of small particles in a viscous fluid

- Iow Reynolds number flow about particles
- each particle sees a linear flow (local rheology)

Suspension of small particles in a viscous fluid

- Iow Reynolds number flow about particles
- each particle sees a linear flow (local rheology)

Everywhere

$$\sigma_{ij} = -p\delta_{ij} + 2\mu e_{ij} + \sigma_{ij}^+$$

with viscous solvent stress and extra non-zero only inside particles

Suspension of small particles in a viscous fluid

- Iow Reynolds number flow about particles
- each particle sees a linear flow (local rheology)

Everywhere

$$\sigma_{ij} = -p\delta_{ij} + 2\mu e_{ij} + \sigma^+_{ij}$$

with viscous solvent stress and extra non-zero only inside particles

Ensemble average (from turbulence research)

$$\langle \sigma_{ij}
angle = - \langle p
angle \delta_{ij} + 2 \mu \langle e_{ij}
angle + \langle \sigma_{ij}^+
angle$$

Switch to volume average. Use divergence theorem for force-free particles

$$\langle \sigma_{ij}^+ \rangle = n \int_{A} \sigma_{ik} n_k x_j - \mu (u_i n_j + u_j n_i) \, dA$$

with **n** number of particles per unit volume and **A** surface of typical particle.

Switch to volume average. Use divergence theorem for force-free particles

$$\langle \sigma_{ij}^+ \rangle = n \int_{\mathcal{A}} \sigma_{ik} n_k x_j - \mu (u_i n_j + u_j n_i) \, dA$$

with **n** number of particles per unit volume and **A** surface of typical particle.

Much cited result - still 30 pa

L&L 1959

Stress system in a suspension of force-free particles JFM 1970

Long discussion in paper, including:

Antisymmetric stress tensor when couple per unit volume

- Antisymmetric stress tensor when couple per unit volume
- Move integral from surface of particle to far field: Stresslet and Couplet

- Antisymmetric stress tensor when couple per unit volume
- Move integral from surface of particle to far field: Stresslet and Couplet
- Energy budget
 - connection to Einstein's dissipation approach

- Antisymmetric stress tensor when couple per unit volume
- Move integral from surface of particle to far field: Stresslet and Couplet
- Energy budget
 - connection to Einstein's dissipation approach
- Surface tension/energy

Stress system in a suspension of force-free particles JFM 1970

What followed?

▶ 640+ citations.

Stress system in a suspension of force-free particles JFM 1970

- ▶ 640+ citations.
- Evaluation of bulk stress in many rheologies, once solved micro-structure evolution

Stress system in a suspension of force-free particles JFM 1970

- 640+ citations.
- Evaluation of bulk stress in many rheologies, once solved micro-structure evolution
- Ensemble average needed for calculating non-local rheology

Stress system in a suspension of force-free particles JFM 1970

- 640+ citations.
- Evaluation of bulk stress in many rheologies, once solved micro-structure evolution
- Ensemble average needed for calculating non-local rheology
- Homogenisation (1980s) a step back,

Stress system in a suspension of force-free particles JFM 1970

What followed?

- ▶ 640+ citations.
- Evaluation of bulk stress in many rheologies, once solved micro-structure evolution
- Ensemble average needed for calculating non-local rheology
- ► Homogenisation (1980s) a step back,

except compute in periodic boxes

A preparation

Slender-body theory for particles of arbitrary cross-section in Stokes flow JFM 1970

A preparation

Slender-body theory for particles of arbitrary cross-section in Stokes flow JFM 1970

Previous studies: Burgers 1938, Tuck 1964, Cox 1970, Tillet 1970

Not quite matched asymptotics, but ...
Slender-body theory for particles of arbitrary cross-section in Stokes flow JFM 1970

Previous studies: Burgers 1938, Tuck 1964, Cox 1970, Tillet 1970

Not quite matched asymptotics, but ...

Outer represented by line-distribution of Stokeslets

Slender-body theory for particles of arbitrary cross-section in Stokes flow JFM 1970

Previous studies: Burgers 1938, Tuck 1964, Cox 1970, Tillet 1970

Not quite matched asymptotics, but ...

- Outer represented by line-distribution of Stokeslets
- Inner a two-dimensional potential problem gives

Slender-body theory for particles of arbitrary cross-section in Stokes flow JFM 1970

Previous studies: Burgers 1938, Tuck 1964, Cox 1970, Tillet 1970

Not quite matched asymptotics, but ...

- Outer represented by line-distribution of Stokeslets
- Inner a two-dimensional potential problem gives

Effective radius =
$$\begin{cases} a & \text{circle} \\ \frac{1}{2}(b+c) & \text{ellipse} \\ 2^{-n}a & n\text{-star} \end{cases}$$

by conformal transformation

Slender-body theory for particles of arbitrary cross-section in Stokes flow JFM 1970

Previous studies: Burgers 1938, Tuck 1964, Cox 1970, Tillet 1970

Not quite matched asymptotics, but ...

- Outer represented by line-distribution of Stokeslets
- Inner a two-dimensional potential problem gives

Effective radius =
$$\begin{cases} a & \text{circle} \\ \frac{1}{2}(b+c) & \text{ellipse} \\ 2^{-n}a & n\text{-star} \end{cases}$$

by conformal transformation

Evaluated force and couple

Stress generated in a non-dilute suspension of elongated particles by pure straining motion JFM 1971

Stress generated in a non-dilute suspension of elongated particles by pure straining motion JFM 1971

Rods of length ℓ , radius *b*, number per unit volume *n*. Hence average lateral spacing $h = (2n\ell)^{-1/2}$.

Regimes (separation *h* decreasing)

Stress generated in a non-dilute suspension of elongated particles by pure straining motion JFM 1971

Rods of length ℓ , radius *b*, number per unit volume *n*. Hence average lateral spacing $h = (2n\ell)^{-1/2}$.

Regimes (separation *h* decreasing)

• (Very) dilute: $\ell \ll h$

Stress generated in a non-dilute suspension of elongated particles by pure straining motion JFM 1971

Rods of length ℓ , radius *b*, number per unit volume *n*. Hence average lateral spacing $h = (2n\ell)^{-1/2}$.

Regimes (separation h decreasing)

- (Very) dilute: $\ell \ll h$
- Semi-dilute: $b \ll h \ll \ell$

Stress generated in a non-dilute suspension of elongated particles by pure straining motion JFM 1971

Rods of length ℓ , radius *b*, number per unit volume *n*. Hence average lateral spacing $h = (2n\ell)^{-1/2}$.

Regimes (separation h decreasing)

- (Very) dilute: $\ell \ll h$
- Semi-dilute: $b \ll h \ll \ell$
- (Nematic phase transition to aligned rods: $h = \sqrt{b\ell}$)
- ► Concentrated: h ~ b

Stress generated in a non-dilute suspension of elongated particles by pure straining motion JFM 1971

Use bulk-stress paper for formula for stress

Stress generated in a non-dilute suspension of elongated particles by pure straining motion JFM 1971

Use bulk-stress paper for formula for stress Use slender-body paper, with outer boundary condition at typical separation h in place of at length ℓ

Stress generated in a non-dilute suspension of elongated particles by pure straining motion JFM 1971

Use bulk-stress paper for formula for stress Use slender-body paper, with outer boundary condition at typical separation h in place of at length ℓ

Result: extensional viscosity

$$\mu_{\mathsf{ext}} = \mu \frac{4\pi n\ell^3}{9\log h/b}$$

Stress generated in a non-dilute suspension of elongated particles by pure straining motion JFM 1971

Use bulk-stress paper for formula for stress Use slender-body paper, with outer boundary condition at typical separation h in place of at length ℓ

Result: extensional viscosity

$$\mu_{\mathsf{ext}} = \mu \frac{4\pi n\ell^3}{9\log h/b}$$

One data point

Stress generated in a non-dilute suspension of elongated particles by pure straining motion JFM 1971

Use bulk-stress paper for formula for stress Use slender-body paper, with outer boundary condition at typical separation h in place of at length ℓ

Result: extensional viscosity

$$\mu_{\mathsf{ext}} = \mu \frac{4\pi n\ell^3}{9\log h/b}$$

One data point

Much larger than the viscosity of the solvent $\mu,$ which is shear viscosity of suspension

Stress generated in a non-dilute suspension of elongated particles by pure straining motion JFM 1971

What followed?

Correct outer with a Brinkman approach by Shaqfeh (1990)

Stress generated in a non-dilute suspension of elongated particles by pure straining motion JFM 1971

What followed?

- Correct outer with a Brinkman approach by Shaqfeh (1990)
- Anisotropic viscosity ($\mu_{\mathsf{ext}} \gg \mu_{\mathsf{shear}}$) produces anisotropic flow



Stress generated in a non-dilute suspension of elongated particles by pure straining motion JFM 1971

What followed?

- Correct outer with a Brinkman approach by Shaqfeh (1990)
- Anisotropic viscosity ($\mu_{\mathsf{ext}} \gg \mu_{\mathsf{shear}}$) produces anisotropic flow



Needed in polymer rheology

Sedimentation in a dilute suspension of spheres JFM 1972

Sedimentation in a dilute suspension of spheres JFM 1972

Settling velocity of test sphere due to 2nd at distance r

$$V(r) = V_0 + \Delta V(r)$$

with Stokes velocity for isolated sphere $V_0=2\Delta\rho ga^2/9\mu$

Sedimentation in a dilute suspension of spheres JFM 1972

Settling velocity of test sphere due to 2nd at distance r

$$V(r) = V_0 + \Delta V(r)$$

with Stokes velocity for isolated sphere $V_0=2\Delta
ho ga^2/9\mu$

Far field form from reflections

$$\Delta V(r)/V_0 = \frac{a}{r} + \frac{a^3}{r^3}$$
 1st reflection
+ $\frac{a^4}{r^4} + \frac{a^6}{r^6} + \dots$ 2nd reflection
+ $\frac{a^7}{r^7} + \frac{a^9}{r^9} + \dots$ 3rd reflection
+ \dots

Sedimentation in a dilute suspension of spheres JFM 1972

Naive pairwise addition of disturbances within large domain $r \leq R$, with *n* spheres per unit volume

$$\langle \Delta V \rangle = \int_{r=2a}^{R} V_0 \left(\frac{a}{r} + \ldots\right) n \, dV$$

Sedimentation in a dilute suspension of spheres JFM 1972

Naive pairwise addition of disturbances within large domain $r \leq R$, with *n* spheres per unit volume

$$\langle \Delta V \rangle = \int_{r=2a}^{R} V_0 \left(\frac{a}{r} + \dots\right) n \, dV = O\left(V_0 c \frac{R^2}{a^2}\right)$$
$$c = \frac{4\pi}{3} n a^3 \text{ the volume fraction.}$$

Sedimentation in a dilute suspension of spheres JFM 1972

Naive pairwise addition of disturbances within large domain $r \leq R$, with *n* spheres per unit volume

$$\langle \Delta V \rangle = \int_{r=2a}^{R} V_0 \left(\frac{a}{r} + \ldots \right) n \, dV = O\left(V_0 c \frac{R^2}{a^2} \right)$$

 $c = \frac{4\pi}{3}na^3$ the volume fraction.

The divergence problem:

- Does mean settling velocity depend on size of domain?
- Or is it an intrinsic property independent of domain?
 i.e. is pairwise addition naive?

Sedimentation in a dilute suspension of spheres JFM 1972

Batchelor's renormalization:

$$\Delta V = \left(1 + rac{a^2}{6}
abla^2\right) u(x) \big|_{\text{test sphere}} + \text{higher reflections}$$

Sedimentation in a dilute suspension of spheres JFM 1972

Batchelor's renormalization:

$$\Delta V = \left(1 + rac{a^2}{6}
abla^2\right) u(x) \big|_{\text{test sphere}} + \text{higher reflections}$$

Pairwise sum of $O(V_0 a^4/r^4)$ higher reflections is convergent

Sedimentation in a dilute suspension of spheres JFM 1972

Batchelor's renormalization:

 $\Delta V = \left(1 + \frac{a^2}{6}\nabla^2\right) u(x)\big|_{\text{test sphere}} + \text{higher reflections}$ Pairwise sum of $O(V_0 a^4/r^4)$ higher reflections is convergent Now $\langle u \rangle_{\text{everywhere}} = 0$,

Sedimentation in a dilute suspension of spheres JFM 1972

Batchelor's renormalization:

$$\Delta V = \left(1 + \frac{a^2}{6}\nabla^2\right)u(x)\big|_{\text{test sphere}} + \text{higher reflections}$$

Pairwise sum of $O(V_0 a^4/r^4)$ higher reflections is convergent
Now $\langle u \rangle_{\text{everywhere}} = 0$, so
 $\langle u \rangle_{\text{test sphere}} = -\frac{11}{2}V_0c$ back flow

Sedimentation in a dilute suspension of spheres JFM 1972

Batchelor's renormalization:

$$\Delta V = \left(1 + \frac{a^2}{6}\nabla^2\right)u(x)\big|_{\text{test sphere}} + \text{higher reflections}$$

Pairwise sum of $O(V_0 a^4/r^4)$ higher reflections is convergent
Now $\langle u \rangle_{\text{everywhere}} = 0$, so
 $\langle u \rangle_{\text{test sphere}} = -\frac{11}{2}V_0c$ back flow
 $\langle \frac{a^2}{6}\nabla^2 u \rangle_{\text{test sphere}} = \frac{1}{2}V_0c$

Sedimentation in a dilute suspension of spheres JFM 1972

Batchelor's renormalization:

$$\Delta V = \left(1 + \frac{a^2}{6}\nabla^2\right) u(x)\big|_{\text{test sphere}} + \text{higher reflections}$$
Pairwise sum of $O(V_0 a^4/r^4)$ higher reflections is convergent
Now $\langle u \rangle_{\text{everywhere}} = 0$, so
 $\langle u \rangle_{\text{test sphere}} = -\frac{11}{2}V_0c$ back flow
 $\langle \frac{a^2}{6}\nabla^2 u \rangle_{\text{test sphere}} = \frac{1}{2}V_0c$
 $\langle \text{higher reflections} \rangle = -1.55V_0c$

Hence

$$\langle V \rangle = V_0(1-6.55c)$$

Sedimentation in a dilute suspension of spheres JFM 1972

Assumes uniform separation of pairs for result

$$\langle V \rangle = V_0(1-6.55c)$$

Sedimentation in a dilute suspension of spheres JFM 1972

Assumes uniform separation of pairs for result

$$\langle V \rangle = V_0(1-6.55c)$$

Coefficient smaller if more close pairs.

Sedimentation in a dilute suspension of spheres JFM 1972

Assumes uniform separation of pairs for result

$$\langle V \rangle = V_0(1-6.55c)$$

Coefficient smaller if more close pairs.

If crystal then

$$V = V_0 \left(1 - k c^{1/3} \right)$$

Sedimentation in a dilute suspension of spheres JFM 1972

Assumes uniform separation of pairs for result

$$\langle V \rangle = V_0(1-6.55c)$$

Coefficient smaller if more close pairs.

If crystal then

$$V = V_0 \left(1 - k c^{1/3} \right)$$

For sedimentation: given force, find average velocity For porous media: given velocity, find average force

$$\langle F \rangle = F_0 \left(1 + \frac{3}{\sqrt{2}} c^{1/2} + \frac{3}{2} c \right)$$

Sedimentation in a dilute suspension of spheres JFM 1972

What followed?

Sedimentation in a dilute suspension of spheres JFM 1972

What followed?

► 750+ citations

Sedimentation in a dilute suspension of spheres JFM 1972

What followed?

- ► 750+ citations
- Erroneous applications subtracting wrong infinity
... renormalization of hydrodynamic interactions

Sedimentation in a dilute suspension of spheres JFM 1972

What followed?

- ► 750+ citations
- Erroneous applications subtracting wrong infinity
- Alternative averaged equation approach recognising divergences as change of ρ and μ from solvent to suspension values

... renormalization of hydrodynamic interactions

Sedimentation in a dilute suspension of spheres JFM 1972

What followed?

- 750+ citations
- Erroneous applications subtracting wrong infinity
- Alternative averaged equation approach recognising divergences as change of ρ and μ from solvent to suspension values
- Much more from Batchelor, e.g.

Brownian diffusion of particle with hydrodynamic interactions JFM 1976

$$D = \frac{1 - 6.55c}{6\pi\mu a} \left[\frac{c}{1 - c} \frac{\partial \mu}{\partial c} = kT(1 + 8c + 30c^2) \right]$$

The determination of the bulk stress in a suspension of spherical particles to order $c^2\,$ JFM 1972

The determination of the bulk stress in a suspension of spherical particles to order $c^2\,$ JFM 1972

Another renormalization of hydrodynamic interactions, with J.T Green (PhD 1970)

The determination of the bulk stress in a suspension of spherical particles to order $c^2\,$ JFM 1972

Another renormalization of hydrodynamic interactions, with J.T Green (PhD 1970)

For pure straining, by trajectory calculation of nonuniform probability distribution of separation of pairs

$$\mu * = \mu \left(1 + \frac{5}{2}c + 7.6c^2 \right)$$

The determination of the bulk stress in a suspension of spherical particles to order c^2 JFM 1972

Another renormalization of hydrodynamic interactions, with J.T Green (PhD 1970)

For pure straining, by trajectory calculation of nonuniform probability distribution of separation of pairs

$$\mu * = \mu \left(1 + \frac{5}{2}c + 7.6c^2 \right)$$

For simple shear, problem of closed trajectories (k = 5.2?)

$$\mu * = \mu \left(1 + \frac{5}{2}c + kc^2 \right)$$

The determination of the bulk stress in a suspension of spherical particles to order $\rm c^2$ JFM 1972

Another renormalization of hydrodynamic interactions, with J.T Green (PhD 1970)

For pure straining, by trajectory calculation of nonuniform probability distribution of separation of pairs

$$\mu * = \mu \left(1 + \frac{5}{2}c + 7.6c^2 \right)$$

For simple shear, problem of closed trajectories (k = 5.2?)

$$\mu * = \mu \left(1 + \frac{5}{2}c + kc^2 \right)$$

Case of strong Brownian motion (JFM 1977)

$$\mu * = \mu \left(1 + \frac{5}{2}c + \frac{6.2c^2}{2} \right)$$

The determination of the bulk stress in a suspension of spherical particles to order $\rm c^2$ JFM 1972

Another renormalization of hydrodynamic interactions, with J.T Green (PhD 1970)

For pure straining, by trajectory calculation of nonuniform probability distribution of separation of pairs

$$\mu * = \mu \left(1 + \frac{5}{2}c + 7.6c^2 \right)$$

For simple shear, problem of closed trajectories (k = 5.2?)

$$\mu * = \mu \left(1 + \frac{5}{2}c + kc^2 \right)$$

Case of strong Brownian motion (JFM 1977)

$$\mu * = \mu \left(1 + \frac{5}{2}c + \frac{6.2c^2}{2} \right)$$

Note strain-hardening and shear-thinning rheology

80

Sedimentation in a dilute polydisperse system of interacting spheres JFM 1982, Parts I, II, Corrigendum

$$\langle V_i \rangle = V_{i0} (1 + S_{ij} c_j)$$

Sedimentation in a dilute polydisperse system of interacting spheres JFM 1982, Parts I, II, Corrigendum

$$\langle V_i \rangle = V_{i0} \left(1 + S_{ij} c_j \right)$$

If equal density and nearly equal sizes, then increase of near pairs and

$$\langle V \rangle = V_0 \left(1 - 5.6c \right)$$

Sedimentation in a dilute polydisperse system of interacting spheres JFM 1982, Parts I, II, Corrigendum

$$\langle V_i \rangle = V_{i0} \left(1 + S_{ij} c_j \right)$$

If equal density and nearly equal sizes, then increase of near pairs and

$$\langle V \rangle = V_0 \left(1 - 5.6c \right)$$

as dilute Richardson-Zaki.

Sedimentation in a dilute polydisperse system of interacting spheres JFM 1982, Parts I, II, Corrigendum

$$\langle V_i \rangle = V_{i0} \left(1 + S_{ij} c_j \right)$$

If equal density and nearly equal sizes, then increase of near pairs and

$$\langle V \rangle = V_0 \left(1 - 5.6c \right)$$

as dilute Richardson-Zaki.

Also

Diffusion in a dilute polydisperse system of interacting spheres JFM 1983

Brady's Stokesian Dynamics (1988) numerical simulations at moderate concentrations of hard spheres



Brady's Stokesian Dynamics (1988) numerical simulations at moderate concentrations of hard spheres



Boundary integral methods for emulsions (1996)

Brady's Stokesian Dynamics (1988) numerical simulations at moderate concentrations of hard spheres



Boundary integral methods for emulsions (1996) Ladd's Lattice Boltzmann simulations (1996)

More sedimentation

More sedimentation

Structure formation in bidisperse sedimentation JFM 1986 with van Rensburg



Light and heavy particles of same size, both at c = 0.2

Fast separation

What followed in sedimentation

Inclined settling (Boycott effect) Acrivos 1979

What followed in sedimentation

- Inclined settling (Boycott effect) Acrivos 1979
- Structural instability for fibres Koch & Shaqfeh 1989



What followed in sedimentation

- Inclined settling (Boycott effect) Acrivos 1979
- Structural instability for fibres Koch & Shaqfeh 1989
- ► Fluctuations depend on the size of box Guazzelli 2001 $\langle V'^2 \rangle = O\left(V_0^2 c \frac{R}{a}\right)$



Suspensions: what happened in parallel

Leal & Hinch 1972

Leal & Hinch 1972

Deformation and breakup of drops, emulsions Acrivos 1970

Leal & Hinch 1972

- Deformation and breakup of drops, emulsions Acrivos 1970
- Electrical Double Layers and VdW forces Russel 1985

Leal & Hinch 1972

- ► Deformation and breakup of drops, emulsions Acrivos 1970
- Electrical Double Layers and VdW forces Russel 1985
- Rough spheres Leighton 1989

What followed

Rheology, with experiments

- Rheology, with experiments
- Fluid dynamics of non-Newtonian fluids
 Normal stresses, stress relaxation, stress saturation, elastic boundary layers, anisotropic flow

- Rheology, with experiments
- Fluid dynamics of non-Newtonian fluids
 Normal stresses, stress relaxation, stress saturation, elastic boundary layers, anisotropic flow
- Electro- and Magneto- rheological fluids

- Rheology, with experiments
- Fluid dynamics of non-Newtonian fluids
 Normal stresses, stress relaxation, stress saturation, elastic boundary layers, anisotropic flow
- Electro- and Magneto- rheological fluids
- Microfluidic devices

drop production, mixing, slip at wall

- Rheology, with experiments
- Fluid dynamics of non-Newtonian fluids
 Normal stresses, stress relaxation, stress saturation, elastic boundary layers, anisotropic flow
- Electro- and Magneto- rheological fluids
- Microfluidic devices

drop production, mixing, slip at wall

Micro-rheology

Batchelor's fluidized beds

Batchelor's fluidized beds (not low Re)

- Low Reynolds-number bubbles in fluidized beds Arch Mech 1974
- Expulsion of particles from a buoyant blob in a fluidized bed JFM 1994
- A new theory of the instability of a uniform fluidized-bed JFM 1988
- Instability of stationary unbounded stratified fluid JFM 1991 with Nitsche
- Secondary instability of a gas-fluidized bed JFM 1993

Batchelor's fluidized beds (not low Re)

- Low Reynolds-number bubbles in fluidized beds Arch Mech 1974
- Expulsion of particles from a buoyant blob in a fluidized bed JFM 1994
- A new theory of the instability of a uniform fluidized-bed JFM 1988
- Instability of stationary unbounded stratified fluid JFM 1991 with Nitsche
- Secondary instability of a gas-fluidized bed JFM 1993
- So bubbles in gas-luidized, not liquid-fluidized (Jackson 2000)
Batchelor's fluidized beds (not low Re)

- Low Reynolds-number bubbles in fluidized beds Arch Mech 1974
- Expulsion of particles from a buoyant blob in a fluidized bed JFM 1994
- A new theory of the instability of a uniform fluidized-bed JFM 1988
- Instability of stationary unbounded stratified fluid JFM 1991 with Nitsche
- Secondary instability of a gas-fluidized bed JFM 1993

So bubbles in gas-luidized, not liquid-fluidized (Jackson 2000) But dense regions have higher drag so rise!

Batchelor's last paper

Break-up of a falling drop containing disperse particles JFM 1997 with Nitsche

Batchelor's last paper

Break-up of a falling drop containing disperse particles JFM 1997 with Nitsche



Batchelor's last paper

Break-up of a falling drop containing disperse particles JFM 1997 with Nitsche



But a little later

- video by Metzger, Nicolas & Guazzelli 2007

112

Batchelor's micro-hydrodynamics

Outline Bulk stress Slender-body Extensional viscosity of rods Renormalization, sedimentation Renormalization, effective viscosity Polydispersity Beyond c^2 More sedimentation More suspensions Fluidized bed Cloud