

Levitation and locomotion on an air-table of plates with herringbone grooves.

John Hinch and H el ene de Maleprade

DAMTP-CMS, University of Cambridge

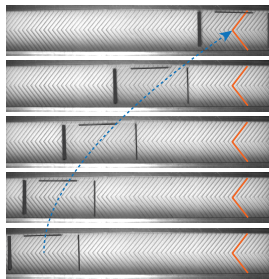
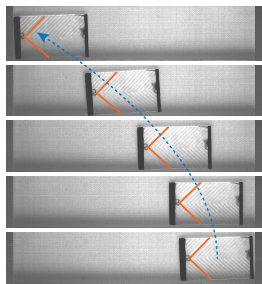
September 30, 2019

with Dan Soto & David Qu er e in Paris, and
Maximilian Sch ur, Steffen Hardt & Tobias Baier in Darmstadt

Experiments on an air table

Grooves on the floating body,

on the table

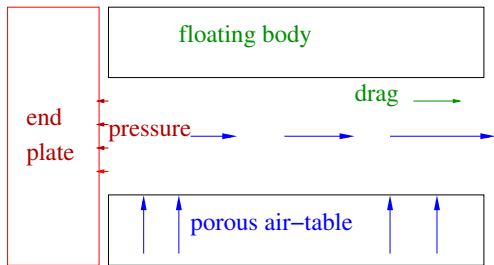


Accelerates to the left,

to the right.

Why different direction?

Left or right?



End plate attached to base (grooved table)

→ body **dragged** by flow to right

End plate attached to top (grooved body)

→ **pressure** pushes to left

Pressure also levitates floating body.

Some approximations

Grooves: width \gg height \rightarrow 2D flow, ignore spanwise

Grooves: length \gg height \rightarrow boundary layer equations
(lubrication theory with inertia)

Two more approximations later.

Study 2D flow down groove

Boundary layer equations with $p = p(x)$:

$$u_x + v_y = 0$$

$$\rho(u_t + uu_x + vu_y) = -p_x + \mu u_{yy}$$

BC on porous plate $y = 0$: $u = 0$ and

$$v = \frac{k}{\mu}(P_2 - p) \quad \text{with plate permeability} \quad k = \frac{\pi r^4}{8\ell d^2}$$

for holes of radius r , length ℓ , separation d

Nondimensionalise

Boundary layer equations with $p = p(x)$:

$$u_x + v_y = 0$$

$$Re(u_t + uu_x + vu_y) = -Kp_x + u_{yy}$$

BC on porous plate $y = 0$: $u = 0, \quad v = 1 - p$

$$K = \frac{\Delta p \text{ across porous plate}}{\Delta p \text{ down groove}} = \frac{8H^3 \ell d^2}{\pi L^2 r^4}$$

Forces:

$$\text{Propulsion } F_H = p(0) + \int_0^1 u_y|_{y=1} dx, \quad \text{Levitation } F_V = \int_0^1 p dx$$

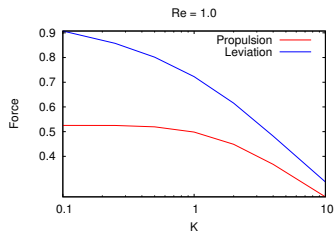
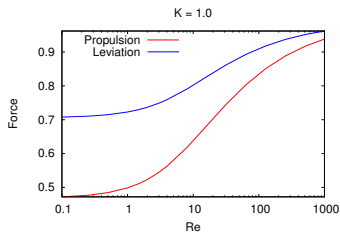
NB: different non-dimensionalisation, F_H smaller by area ratio H/L

Numerical solution

Integrate $v_y = -u_x$ from $y = 1$ to $y = 0 \rightarrow v(x, y = 0)$

Porous plate BC $p(x) = 1 - v(x, 0)$ into momentum equation.

Time step momentum to steady state.



NB: two forces (in different non-dimensionalisations) within a factor of two

NB: small change with Re , but decrease at large K (short groove)

Asymptotics

- ▶ Short groove, $K \gg 1$,

pressure drop mostly across porous plate, so $v(x, 0) \approx 1$

- ▶ $Re \ll 1$

- ▶ $Re \gg 1$

- ▶ Long groove, $K \ll 1$,

so $p = 1$ in groove except very near outlet

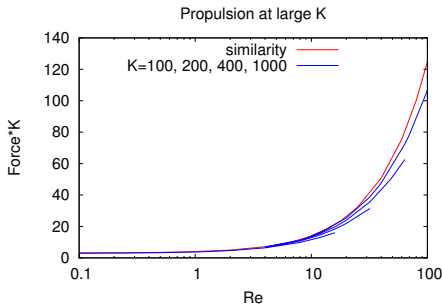
Short groove, $K \gg 1$, Re arbitrary

Most pressure drop across porous plate, so $p \sim 0$ in groove, so $v(x, y = 0) = 1$, hence similarity solution

$$u(x, y) = -xg'(y), \quad v(x, y) = g(y), \quad p = \frac{B}{K}(1 - x^2)$$

Momentum equation then

$$Re(g'^2 - gg'') = 2B - g'''$$



need $K \gg Re$

$$\text{also } F_V = \frac{2}{3}p(0)$$

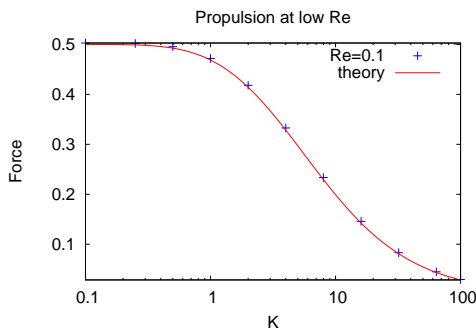
Low Reynolds number, K arbitrary

Lubrication theory

$$q = \int_0^1 u(x, y) dy = -\frac{K}{12} p_x, \quad q_x = v(x, 0) = 1 - p$$

so

$$F_H = \frac{1}{2}(1 - \operatorname{sech} \sqrt{12/K}), \quad F_V = 1 - \sqrt{K/12} \tanh \sqrt{12/K}$$



At $Re \ll 1$:
Drag on top =
Drag on bottom =
 $-\frac{1}{2}$ Pr force on end

High Reynolds number, K arbitrary

Separable solution with streamfunction

$$\psi(x, y) = -f(x)g(y).$$

Inviscid, vorticity constant along streamlines. Try

$$\omega(\psi) = k^2\psi, \quad \omega = -\psi_{yy}.$$

Hence

$$g(y) = \cos \frac{\pi}{2}y.$$

$$u = -f(x)g'(y), \quad v = f'(x)g(y)$$

Bernoulli integral

$$\beta^2 f^2 = p_0 - p, \quad \text{with } \beta^2 = \pi^2 Re/8K, \quad p_0 = p(0)$$

High Reynolds number, inviscid

Bernoulli integral

$$\beta^2 f^2 = p_0 - p.$$

Porous plate $f' = 1 - p$, so

$$f(x) = \frac{\sqrt{1 - p_0}}{\beta} \tan\left(\beta x \sqrt{1 - p_0}\right)$$

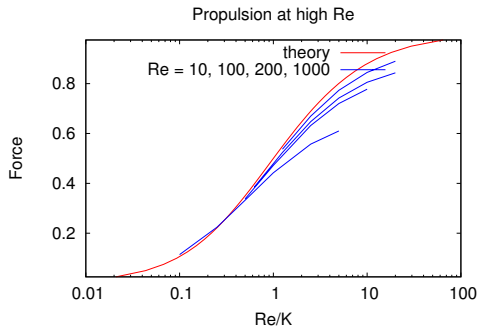
$p_0 = p(0)$ is determined by $p(1) = 0$:

$$\tan^2\left(\beta \sqrt{1 - p_0}\right) = \frac{p_0}{1 - p_0}$$

High Reynolds number, inviscid

Finally forces

$$F_H = -p_0, \quad F_V = 1 - \sqrt{p_0/\beta}$$

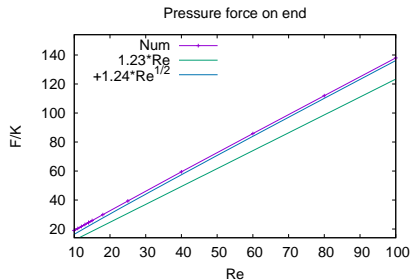
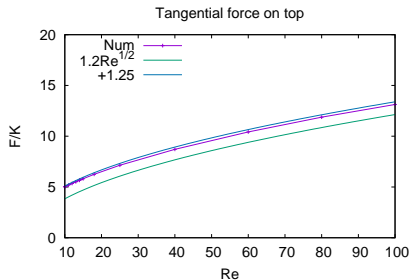


Need boundary layer corrections

Boundary layer corrections

For short grooves, there is a $Re^{1/2}$ similarity boundary layer

The boundary layer exerts a drag $1.2133Re^{1/2}$ on the top plate,
and has a displacement thickness $1.0366Re^{-1/2}$,
which leads to an enhanced pressure on the left end

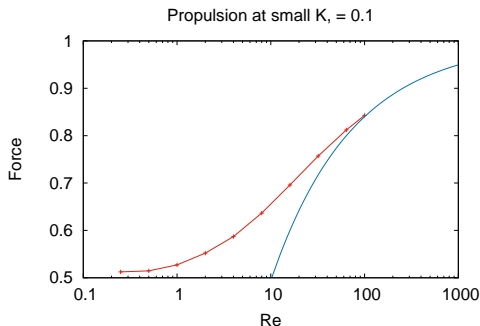


Long groove, $K \ll 1$, Re arbitrary

Most of the groove is at $p = 1$, the pressure under the air-table.

Hence pressure part of Propulsion and Levitation ~ 1 .

Pressure drop and flow only near exit of groove, so tough numerics.



But frictional drag halves the Propulsion at low Re .

Double limits

	Propulsion F_H	Levitation F_V
$K, Re \ll 1$	0.5	$1 - \sqrt{K/12}$
$K \ll 1 \ll Re$	1	$1 - \sqrt{8K/\pi^2 Re}$
$Re \ll 1 \ll K$	$3/K$	$4/K$
$1 \ll Re \ll K$	$\pi^2 Re/8K$	$\pi^2 Re/12K$
$1 \ll K \ll Re$	$1 - 2K/Re$	$1 - \sqrt{8K/\pi^2 Re}$

At $Re \gg 1$ and $K \ll 1$, uniform levitation pressure = propulsion pressure, i.e. $F_H = F_V$

At $Re \ll 1$, frictional drag reduces F_H by 50%.

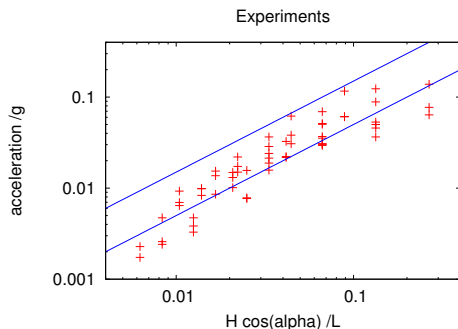
At $K \gg 1$, pressure is parabolic, reduces F_V by 2/3.

Hence $F_H/F_V = 0.5$ to 1.5 (different non-dimensionalisations)

Experiments

Reinstating the different dimensional factors, and resolving force along direction that body moves, predict

$$\text{Propulsion/weight} = (0.5 \text{ to } 1.5)h \cos \alpha / \ell$$



Grooves 5 heights & 4 lengths, Plates 3 thickness,
($K = 0.15 \rightarrow 2800$, $Re = 1 \rightarrow 200$) $a = 0.02 \rightarrow 1.4 \text{ m/s}^2$