

**Example Sheet 4: Ray Theory**

1. *The governing equation of sound waves in a stratified fluid.* A gas exactly satisfies the equation of state

$$p(\rho, T) = \lambda \rho^{\frac{1}{2}},$$

where  $\lambda$  is a constant independent of the temperature  $T$ . The gas fills  $z \geq 0$  and is at rest in a gravity field  $\mathbf{g} = (0, 0, -g)$ . Show that the density is given by

$$\rho = \frac{\lambda^2}{g^2(z - z_0)^2},$$

where  $z_0 < 0$  is a constant.

Show that linear sound waves propagating in the above gas satisfy

$$\frac{\partial^2 \tilde{\rho}}{\partial t^2} - \nabla \cdot (c_0^2 \nabla \tilde{\rho} - \frac{3}{2} \tilde{\rho} \mathbf{g}) = \frac{\partial^2 \tilde{\rho}}{\partial t^2} - \nabla \cdot \left( \rho_0^{3/2} c_0^2 \nabla \left( \rho_0^{-3/2} \tilde{\rho} \right) \right) = 0,$$

where  $\tilde{\rho}(\mathbf{x}, t)$  is the density perturbation,  $\rho_0(z)$  is the unperturbed density, and  $c_0(z)$  is the local speed of sound. Deduce that a slowly varying approximation for sound waves with wavenumber  $k$  is valid in  $z \geq 0$  if  $|kz_0| \gg 1$ .

2. *The wave-crest pattern near a shore line.* (Old Tripos: 87327). Surface waves on water have a dispersion relation  $\omega = \Omega(\kappa; x, z)$  where  $\kappa^2 = k_1^2 + k_3^2$ ,  $(x, z)$  are coordinates in the plane of the surface, and the medium is 'slowly varying' in the  $(x, z)$  coordinates. Assume the relation

$$\omega_t + (\mathbf{c}_g \cdot \nabla) \omega = 0,$$

where  $\mathbf{c}_g$  is the group velocity, and deduce

(a) that  $\omega$  is constant on rays,  $dz/dx = k_3/k_1$ ,

(b) that the wave crests at any instant are given by  $dz/dx = -k_1/k_3$ .

Surface tension effects are negligible, and the wave motion takes place over a sloping beach of depth  $h(x) = \alpha x^{1/2}$ , with  $\alpha$  a small positive constant. The dispersion relation for such waves may be assumed to be given by  $\Omega^2 = g\kappa \tanh \kappa h$ . Far from the shore-line  $x = 0$ , the waves are plane, have frequency  $\omega$ , and have angle  $\Phi$  between the crests and the shore-line. As the waves propagate towards the shore they become non-planar. Obtain the parametric equations

$$x = \frac{\lambda^2 g^2}{\alpha^2 \omega^4} \tanh^2 \lambda,$$

$$z - z_0 = \frac{g^2}{\alpha^2 \omega^4} \int_0^\lambda \frac{(1 - \tanh^2 p \sin^2 \Phi)^{1/2}}{\tanh p \sin \Phi} \frac{d}{dp} (p^2 \tanh^2 p) dp,$$

for the wave crest which passes through the shore-line at  $z = z_0$ . Show that near the shore-line the equation of the wave crest can be written explicitly as

$$(z - z_0)^4 \approx \left( \frac{4}{3 \sin \Phi} \right)^4 \frac{g^2}{\alpha^2 \omega^4} x^3.$$

3.\* *Wave breaking.* Ocean surface waves propagate obliquely from  $x = -\infty$  on water of depth  $h(x, z) = -\beta x$  towards a straight beach at  $x = 0$  where they break and are dissipated. For a slowly varying depth,  $\beta \ll 1$ , you may assume that the dispersion relation is

$$\Omega^2 = g\kappa \tanh \kappa h ,$$

where  $\kappa^2 = k_1^2 + k_3^2$  for the surface wavenumber  $(k_1, k_3)$ . As in question 2 the frequency  $\omega$ , and the component  $k_3$  of the wavenumber along the beach, remain constant. Deduce that the shorewards component of the wavenumber,  $k_1$ , increases, and  $k_1 \sim \omega(-g\beta x)^{-1/2}$  as  $x \rightarrow 0$ .

Find how the amplitude  $a$  of the waves varies, where  $2a$  is the difference in height between the crests and the troughs of the waves. Show that if the waves break when  $a\kappa = 0.1$  in a region where  $\kappa h < 1$ , then the point  $x_b$  at which they break is given by

$$a^2(\infty) \frac{k_1(\infty)}{\kappa(\infty)} \frac{\omega}{(-g\beta^3 x_b^3)^{1/2}} = 0.02 .$$

[*Hint:* Write down the solution for the free-surface, calculate the mean potential energy, and use equipartition of energy.]

4. *Sound rays in a slowly varying medium.* Deduce that for a time-independent, slowly varying, medium the frequency  $\omega$  is constant at a ‘ray point’ moving with the group velocity. If moreover the properties of the medium are independent of two Cartesian coordinates, say  $x$  and  $y$ , deduce Snell’s law that

$$\sin \alpha \propto c ,$$

where  $\alpha$  is the angle between the wavenumber  $\mathbf{k}$  and the  $z$ -axis, and  $c$  is the local phase speed for waves of wavenumber  $\mathbf{k}$ . For what type of dispersion relation is the direction of the ray parallel to  $\mathbf{k}$ ?

Consider a dispersion relation of the form  $\omega = A|\mathbf{k}|^z$ , where  $A$  is a constant, and let  $ds$  be an element of arc length along a ray. Show that in this case  $d\alpha/ds$  is constant along a ray, and hence that each ray is the arc of a circle. Show that a wave packet moving towards the plane  $z = 0$  takes an infinite time to reach it.

5. *The wave pattern generated by a duck swimming on a pseudo-fluid.* For a slowly varying, two-dimensional wave pattern of the form  $A(\mathbf{x}, t) \exp(i\varepsilon^{-1}\theta(\mathbf{x}, t))$ , and a local dispersion relation  $\omega = \Omega(\mathbf{k}, \mathbf{x}, t)$ , derive the ray-tracing equations

$$\dot{k}_i = -\frac{\partial\Omega}{\partial x_i}, \quad \dot{x}_i = \frac{\partial\Omega}{\partial k_i}, \quad \varepsilon^{-1}\dot{\theta} = -\Omega + k_i \frac{\partial\Omega}{\partial k_i}, \quad (i = 1, 2).$$

For a homogeneous, time-independent (but not necessarily isotropic) medium, show that all rays are straight lines. When the waves have zero frequency, deduce that if the point  $\mathbf{x}$  lies on a ray emanating from the origin in the direction given by a unit vector  $\hat{\mathbf{c}}_g$ , then

$$\theta(\mathbf{x}) = \theta(0) + \varepsilon \hat{\mathbf{c}}_g \cdot \mathbf{k} |\mathbf{x}|.$$

Consider a duck swimming steadily with velocity  $V$  in a homogeneous pseudo-fluid. In the duck's frame, the dispersion relation is found to be

$$\Omega(k_1, k_2) = \alpha(k_1^2 + k_2^2)^{\frac{1}{3}} - V k_1,$$

The duck generates a steady wave pattern. By writing  $(k_1, k_2) = \kappa(\cos \phi, \sin \phi)$ , show that the waves satisfy

$$\kappa = \frac{\alpha^3}{V^3 \cos^3 \phi},$$

and that the group velocity of these waves can be expressed as

$$\mathbf{c}_g = \frac{1}{3}V(-\cos^2 \phi - 3\sin^2 \phi, 2\sin \phi \cos \phi).$$

Deduce that the waves occupy a wedge of semi-angle  $\frac{1}{6}\pi$  about the negative  $x_1$ -axis. Find equation[s] describing the wave crests, and sketch the wave-crest pattern.

6. *'Reflection' and 'absorption' of internal gravity waves.* Two-dimensional internal gravity waves on a 'slowly varying' shear flow in the atmosphere satisfy the dispersion relation

$$\omega = \gamma y k + \frac{Nk}{(k^2 + \ell^2)^{1/2}},$$

where  $\gamma$  and  $N$  are positive constants, and  $k$  and  $\ell$  are the  $x$ - and  $y$ -components of the wave number respectively. Show that as a wave packet moves,  $\omega$  and  $k$  remain constant, while

$$\ell(t) = \ell_0 - \gamma k t,$$

where  $\ell_0$  is a constant. If  $\ell_0$  is positive, describe the motion in the [vertical]  $y$ -direction of a wave packet generated at the origin. Sketch the ray in the neighbourhood of

$$(a) \quad y = -\frac{N}{\gamma} \left( \frac{1}{k} - \frac{1}{(k^2 + \ell_0^2)^{1/2}} \right),$$

and

$$(b) \quad y = \frac{N}{\gamma(k^2 + \ell_0^2)^{1/2}}.$$