# Small particles in a viscous fluid

Course in three parts

- 1. A quick course in micro-hydrodynamics
- 2. Sedimentation of particles
- 3. Rheology of suspensions

Good textbook for parts 1 & 2: A Physical Introduction to Suspension Dynamics by Elisabeth Guazzelli, Jeffrey F. Morris and Sylvie Pic (Cambridge Texts in Applied Mathematics 2012).

# Part 2. Sedimentation of particles

An isolated particle

Two particles

Finite clouds

Suspensions

Fluctuations in velocity of particles

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### An isolated particle

A sphere Stokes settling velocity A spherical viscous drop Flow past an ellipsoid

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Drag on a spherical drop  $2\pi\mu a \frac{2\mu+3\mu_i}{\mu+\mu_i}$ .

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For a disk  $a_1 \ll a_2 = a_3$   
 $F_1 \sim 16\pi\mu a_2 U_1, \quad F_2 \sim \frac{32}{3}\mu a_2 U_2, \quad G_i \sim \frac{8}{3}\mu a_2 \Omega_i$ 

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Important conclusion Drag  $\approx 6\pi\mu \times (\text{largest diameter}).$ 

# Part 2. Sedimentation of particles

#### An isolated particle

#### Two particles

Two equal spheres Two unequal spheres Three equal spheres Method of Reflections Zeroth approximation First reflection Second reflection Collecting terms

#### Finite clouds

Suspensions

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# Fall of two equal spheres

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Approximations for large separations

$$U = U_S\left(1 + \frac{3a}{2r} + \frac{a^3}{r^3}\right), \quad U_S\left(1 + \frac{3a}{4r} + \frac{a^3}{r^3}\right)$$

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# Method of Reflections

An iterative method for far-field interactions between particles (use lubrication for close particles).

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Can use computer algebra to go 100 reflections.

Zeroth approximation: each sphere falls as effectively isolated

$$\mathbf{u}^{0} = \mathbf{u}^{01} + \mathbf{u}^{02}$$
$$\mathbf{u}^{01}(\mathbf{x}) = \mathbf{U}_{1} \left(\frac{3a_{1}}{4r_{1}} + \frac{a_{1}^{3}}{4r_{1}^{3}}\right) + \frac{(\mathbf{U}_{1} \cdot \mathbf{r}_{1})\mathbf{r}_{1}}{r_{1}^{2}} \left(\frac{3a_{1}}{4r_{1}} - \frac{3a_{1}^{3}}{4r_{1}^{3}}\right)$$

with  $U_1$  isolated fall speed of 1, centred at  $x_1$ ,  $r_1 = x - x_1$ .

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And  $\mathbf{u}^{02}(\mathbf{x})$  similar.

# Method of Refections: first reflection

First refection: Sphere 2 sees  $\mathbf{u}^{01}(\mathbf{x})$  in neighbourhood of  $\mathbf{x} = \mathbf{x}_2$ 

$$\mathbf{u}^{01}(\mathbf{x}) = \mathbf{u}^{01}(\mathbf{x}_2) + \mathbf{r}_2 \cdot \nabla \mathbf{u}^{01}|_{\mathbf{x}_2} + \mathbf{r}_2 \mathbf{r}_2 : \nabla \nabla \mathbf{u}^{01}|_{\mathbf{x}_2} + \dots$$
(1)
(2)
(3)

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(1) dominates, a uniform flow. Sphere 2 (now force free) moves with this, i.e. additions to fall speed  $O(U_1a_1/r_{21}) + O(U_1a_1^3/r_{21}^3)$ .
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(2) smaller, a linear shear flow. Sphere 2 (couple free) rotates with vorticity (antisymmetric part of  $\nabla u^{01}$ ). Cannot deform with straining symmetric part, so reacts with a force-dipole  $O(a_2 \mu a_2 a_2 U_1 a_1/r_{21}^2)$ .

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(3) even smaller, a quadratic flow. Reacts with force quadrupole  $O(a_2^2 \mu a_2 a_2^2 U_1 a_1/r_{21}^3)$ . Addition to fall speed by Faxen  $\frac{1}{6}a_2^2 \nabla^2 u^{01} = O(a_2^2 U_1 a_1/r_{21}^3)$ .

### Method of Reflections: second reflection, collecting terms

Second reflection: Sphere 1 sees  $\mathbf{u}^{12}(\mathbf{x})$  from dipole, quadrupole,... from sphere 2 as uniform + linear + quadratic flows Second reflection: Sphere 1 sees  $\mathbf{u}^{12}(\mathbf{x})$  from dipole, quadrupole,... from sphere 2 as uniform + linear + quadratic flows Additional fall speed due to uniform flows

$$O\left(U_{1}\frac{a_{1}a_{2}^{3}}{r_{21}^{2}}\frac{1}{r_{12}^{2}}\right) + O\left(U_{1}\frac{a_{1}a_{2}^{5}}{r_{21}^{3}}\frac{1}{r_{12}^{3}}\right)$$

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Collecting contributions to fall speed

$$U_1 + U_2 \frac{a_2}{r_{12}} + U_2 \frac{a_2^3}{r_{12}^3} + U_2 \frac{a_1^2 a_2}{r_{12}^3} + U_1 \frac{a_1 a_2^2}{r_{21}^2 r_{12}^2} + U_1 \frac{a_1 a_2^5}{r_{21}^3 r_{12}^3}$$

 $+1/r^8 + \dots$  from 2nd reflection  $+1/r^7 + \dots$  from 3rd + more.

# Part 2. Sedimentation of particles

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Finite clouds Spherical cloud Boycott effect

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Cloud radius *R*. Particles radius *a*, density difference  $\Delta \rho$ , volume fraction  $\phi$ , viscosity of suspension  $\mu^*(\phi)$ .

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$$\frac{2\Delta\!\rho\,\mathsf{a}^2g}{9\mu} + \frac{2\Delta\!\rho\,\phi R^2g}{9\mu}\frac{2(\mu+\mu^*)}{2\mu+3\mu^*}$$

First term from fall of individual particles is much smaller.

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This loss produces a toroid. In a very tall tank, breaks up into two drops, process repeating (Metzner, Nicolas & Guazzelli 2007).

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http://youtu.be/6I5PNdAHILI

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#### Suspensions

Hindered settling Richardson-Zaki Batchelor renormalisation Summing interactions in a dilute suspension Divergence in naive sum Batchelor's 1972 renormalisation Polydispersity in sizes

#### Fluctuations in velocity of particles

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Richardson-Zaki (1954) empirical correlation, for volume fraction  $\phi$ 

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Note near maximum packing reduced to  $\frac{1}{20}$  – like porous media.

# Summing interactions with other particles in a dilute suspension

Settling velocity of test sphere due to 2nd at distance r

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with Stokes velocity for isolated sphere  $U_0 = 2\Delta\rho g a^2/9\mu$ 

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Far-field form from reflections

$$\Delta U(r)/U_0 = \frac{a}{r} + \frac{a^3}{r^3}$$
 1st reflection  
+  $\frac{a^4}{r^4} + \frac{a^6}{r^6} + \dots$  2nd reflection  
+  $\frac{a^7}{r^7} + \frac{a^9}{r^9} + \dots$  3rd reflection  
+  $\dots$ 

$$\langle \Delta U \rangle = \int_{r=2a}^{R} U_0 \left(\frac{a}{r} + \ldots\right) n \, dV$$

$$\langle \Delta U \rangle = \int_{r=2a}^{R} U_0 \left( \frac{a}{r} + \ldots \right) n \, dV = O \left( U_0 \phi \frac{R^2}{a^2} \right)$$

 $\phi = \frac{4\pi}{3} na^3$  the volume fraction.

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Does mean settling velocity depend on size of domain?

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- Does mean settling velocity depend on size of domain?
- Or is it an intrinsic property independent of domain?
   i.e. is pairwise addition naive?

$$\Delta U = \left(1 + \frac{a^2}{6} \nabla^2\right) u(x) \Big|_{\text{test sphere}} + \text{higher reflections}$$

Batchelor's renormalisation:

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Pairwise sum of  $O(U_0 a^4/r^4)$  higher reflections is convergent

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$$\langle \frac{a^2}{6} \nabla^2 u \rangle_{\text{test sphere}} = \frac{1}{2} U_0 \phi$$

Batchelor's renormalisation:

$$\Delta U = \left(1 + \frac{a^2}{6}\nabla^2\right) u(x)\big|_{\text{test sphere}} + \text{higher reflections}$$
Pairwise sum of  $O(U_0 a^4 / r^4)$  higher reflections is convergent  
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Batchelor's renormalisation:

$$\begin{split} \Delta U &= \left(1 + \frac{a^2}{6} \nabla^2\right) u(x) \big|_{\text{test sphere}} + \text{higher reflections} \\ \text{Pairwise sum of } O(U_0 a^4 / r^4) \text{ higher reflections is convergent} \\ \text{Now} \quad \langle u \rangle_{\text{everywhere}} = 0, \quad \text{so} \\ &\quad \langle u \rangle_{\text{test sphere}} = -\frac{11}{2} U_0 \phi \\ &\quad \langle \frac{a^2}{6} \nabla^2 u \rangle_{\text{test sphere}} = \frac{1}{2} U_0 \phi \end{split}$$

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Hence

$$\langle U \rangle = U_0(1-6.55\phi)$$
## Polydispersity in sizes

Batchelor's result assumes uniform distribution of equal spheres

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Hence practical result (Salin 1986)

$$\langle U \rangle = U_0(1-5.6\phi)$$

## Part 2. Sedimentation of particles

An isolated particle

Two particles

Finite clouds

Suspensions

Fluctuations in velocity of particles