## Part II continued - more details on general issues

Last time – Finite Differences

Higher orders - central, 1-sided, non-equispaced

Compact 4th order Poisson solver

Upwinding

Grids - non-Cartesian, stretched, staggered

Conservative

This time - Finite Elements

### Finite Elements = Two ideas

- 1. Simple representation for unknown function over the finite element
  - not point data of FD
- 2. Weak formulation of the governing equations
  - variational statement

### Finite Elements

Good for engineering problems with complex geometries

- 'just' need to triangulate domain

Good for elliptic, OK for parabolic, poor for hyperbolic

Good for accuracy & conservative

Poor difficult programming on unstructured grid

Poor no efficient Poisson solver on unstructured grid

Poor difficult presenting results on unstructured grid

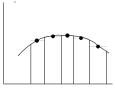
Use packages, do not program yourself

# Representations in 1D

#### a. Constant elements

$$f(x) = f_i$$

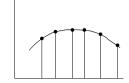
in 
$$x_{i-1} \le x < x_i$$



#### b. Linear elements

$$f(x) = f_{i-1} \frac{x_i - x}{x_i - x_{i-1}} + f_i \frac{x - x_{i-1}}{x_i - x_{i-1}}$$

in 
$$x_{i-1} \le x < x_i$$



## More representations in 1D

First map element to unit interval

$$x(\xi) = x_{i-1} + (x_i - x_{i-1})\xi$$
 for  $0 \le \xi \le 1$ 

#### c. Quadratic elements

$$f(x) = f_{i-1}(1-\xi)(1-2\xi) + f_{i-\frac{1}{2}}4\xi(1-\xi) + f_i\xi(2\xi-1)$$

NB: f' discontinuous at boundaries

#### d. Cubic elements

Obvious generalisation, but better:

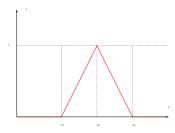
$$f(x) = f_{i-1}(1-\xi)^2(1+2\xi) + f'_{i-1}(1-\xi)^2\xi + f_i\xi^2(3-2\xi) + f'_i\xi^2(1-\xi),$$

Now only f'' discontinuous at boundaries – see splines later

### Basis functions for linear elements

$$\phi_{i}(x) = \begin{cases} \frac{x - x_{i-1}}{x_{i} - x_{i-1}} & \text{in } x_{i-1} \le x \le x_{i} \\ \frac{x_{i+1} - x}{x_{i+1} - x_{i}} & \text{in } x_{i} \le x \le x_{i+1} \\ 0 & \text{otherwise,} \end{cases}$$

with obvious modifications for the end elements.



### basis functions

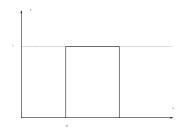
In all cases, write:

$$f(x) = \sum f_i \phi_i(x)$$

 $f_i$  amplitudes  $\phi_i(x)$  basis functions, nonzero only in a few elements

For the constant elements, the basis functions are

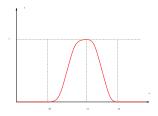
$$\phi_i(x) = \begin{cases} 1 & \text{in } x_{i-1} \le x < x_i \\ 0 & \text{otherwise.} \end{cases}$$

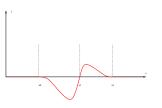


## Basis functions for cubic elements

$$\phi_i(x) = \begin{cases} \frac{(x_{i+1} - x)^2 (x_{i+1} + 2x - 3x_i)}{(x_{i+1} - x_i)^3} & \text{in } x_i \le x < x_{i+1} \\ \frac{(x_{i+1} - x_i)^2 (3x_i - 2x - x_{i-1})}{(x_i - x_{i-1})^3} & \text{in } x_{i-1} \le x < x_i \\ 0 & \text{otherwise,} \end{cases}$$

$$\tilde{\phi}_i(x) = \begin{cases} \frac{(x-x_i)(x_{i+1}-x)^2}{(x_{i+1}-x_i)^2} & \text{in} \quad x_i \le x < x_{i+1} \\ \frac{(x-x_i)(x-x_{i-1})^2}{(x_i-x_{i-1})^2} & \text{in} \quad x_{i-1} \le x < x_i \\ 0 & \text{otherwise.} \end{cases}$$





## Representations in 2D

Mostly triangles, sometimes rectangles

a. Constant elements

$$f(x) = f_i$$
 in each triangle i.

b. Linear elements Need  $\ell_{12}(\mathbf{x})$  vanishing on two vertices, unity on third

$$\ell_{12}(x,y) = \frac{(x-x_1)(y_2-y_1)-(x_2-x_1)(y-y_1)}{(x_3-x_1)(y_2-y_1)-(x_2-x_1)(y_3-y_1)}.$$

Then

$$f(\mathbf{x}) = f_1 \ell_{23}(\mathbf{x}) + f_2 \ell_{31}(\mathbf{x}) + f_3 \ell_{12}(\mathbf{x}).$$

Representation continuous over domain

## more representations in 2D

#### d. Cubic elements

Cubic in 2D has 10 degrees of freedom:

1 constant + 2 linear + 3 quadratic + 4 cubic.

Can fit f and  $\nabla f$  at vertices, plus value in centre = the 'bubble'.

#### e. Basis functions

In all cases, write:

$$f(x) = \sum f_i \phi_i(x)$$

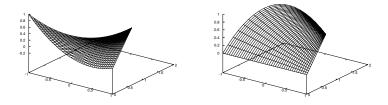
For linear elements,  $\phi_i$  is non-zero at only one vertex, vanishing on opposite sides of triangles, to form a several-sided pyramid.

Local nature  $\rightarrow$  sparse coupling matrices for PDEs

## more representations in 2D

c. Quadratic elements Values at vertices and mid-points

$$f(\mathbf{x}) = f_1 \ell_{23}(\mathbf{x})(2\ell_{23}(\mathbf{x}) - 1) + f_2 \ell_{31}(\mathbf{x})(2\ell_{31}(\mathbf{x}) - 1) + f_3 \ell_{12}(\mathbf{x})(2\ell_{12}(\mathbf{x}) - 1) + f_{23} 4\ell_{12}(\mathbf{x})\ell_{31}(\mathbf{x}) + f_{31} 4\ell_{23}(\mathbf{x})\ell_{12}(\mathbf{x}) + f_{12} 4\ell_{31}(\mathbf{x})\ell_{23}(\mathbf{x}).$$



## more representations in 2D

#### f. Rectangles

Obvious constant elements

Bilinear, taking values at vertices

$$f(\mathbf{x}) = f_1 \xi \eta + f_2 (1 - \xi) \eta + f_3 \xi (1 - \eta) + f_4 (1 - \xi) (1 - \eta).$$

Continuous over domain.

Biquadratic – sum of 9 terms, each product of quadratic in separate coordinates, taking values at vertices and midpoints.

Continuous and continuous tangential derivative at boundaries.

## Variational statement of Poisson problem

 $\nabla^2 f = \rho$  in volume V

with boundary condition, say f = g on surface S, with  $\rho(\mathbf{x})$  and  $g(\mathbf{x})$  given.

Rayleigh-Ritz variational formulation: out of all those functions  $f(\mathbf{x})$  that satisfy BCs, the one that minimises

$$I(f) = \int_V \left(\frac{1}{2} |\nabla f|^2 + \rho f\right) dV$$

also satisfies the Poisson problem.

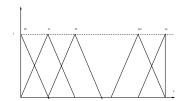
### Details in 1D

$$\frac{d^2f}{dx^2} = \rho \quad \text{in } a < x < b, \quad \text{with } f(a) = A \text{ and } f(b) = B,$$

where  $\rho(x)$ , A and B given.

Divide [a, b] into N equal segments h = (b - a)/N.

Use linear finite elements with basis functions



Unknown f(x) represented (BCs built in)

$$f(x) = A\phi_0(x) + B\phi_N(x) + \sum_{i=1}^{N-1} f_i \phi_i(x)$$

## Substitute FE representation

$$f(\mathbf{x}) = \sum f_i \phi_i(\mathbf{x})$$

Then

$$I(f) = \frac{1}{2} \sum_{ij} f_i f_j \underbrace{\int \nabla \phi_i \cdot \nabla \phi_j}_{\text{global stiffness } K_{ij}} + \sum_i f_i \underbrace{\int \rho \phi_i}_{\text{forcing } r_i}$$

Minimise over  $f_i$ 

$$K_{ij}f_j+r_i=0.$$

With these  $f_i$ , the f satisfies

$$-\int \nabla f \cdot \nabla \phi_i = \int \rho \phi_i \quad \text{for all } i,$$

i.e. satisfy PDE in all (finite)  $\phi_i$  directions.

The weak formulation of the PDE (f can be non- $C^2$ )

### more details in 1D

At interior pts

$$\mathcal{K}_{ij} = \int 
abla \phi_i \cdot 
abla \phi_j = \left\{ egin{array}{ll} 2/h & ext{if } i=j, \ -1/h & ext{if } i=j\pm 1, \ 0 & ext{otherwise.} \end{array} 
ight.$$

by  $\nabla \phi_i = 0, +1, -1, 0$ 

Take given  $\rho(x)$  to be piecewise constant, then forcing

$$r_i = \int \rho(x)\phi_i = h\rho_i.$$

So equation governing unknown amplitudes  $f_i$  becomes

$$\frac{1}{h}(-f_{i-1}+2f_i-f_{i+1})+h\rho_i=0 \quad \text{for } i=1,2,\ldots,N-1,$$

- same for the point values in the finite difference approach.

# more details in 1D

Remark If evaluate  $r_i$  more accurately

$$r_i = \int \rho(x)\phi_i(x) = \rho_i + \frac{h^3}{12}\rho_i'' + O(h^5).$$

So obtain  $f_i$  to  $O(h^4)$ .

Yet f(x) still only  $O(h^2)$  in interior of elements.

Remark For non-equispaced intervals, obtain

$$\frac{1}{h_{i-\frac{1}{2}}}\left(-f_{i-1}+f_{i}\right)+\frac{1}{h_{i+\frac{1}{2}}}\left(f_{i}-f_{i+1}\right)+\frac{h_{i-\frac{1}{2}}+h_{l+\frac{1}{2}}}{2}\rho_{i}=0.$$

i.e. FE approach naturally conservative.