### Resumé: Driven cavity, $\psi$ – $\omega$ formulation

Poisson problem:  $\nabla^2 \psi = -\omega$ , SOR

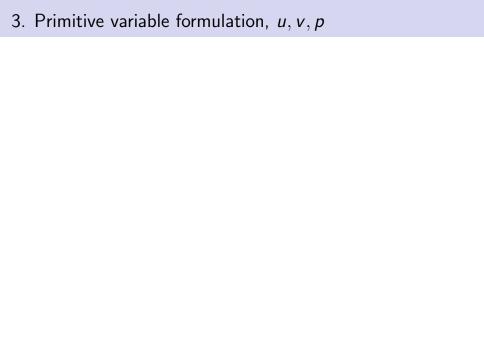
Vorticity evolution: 
$$\frac{\partial \omega}{\partial t} = -\frac{\partial(\omega,\psi)}{\partial(x,y)} + \frac{1}{Re}\nabla^2\omega$$

BC for  $\omega$ 

Timestep instability 
$$\rightarrow \Delta t = \frac{1}{5} Re \Delta x^2$$

Check  $O(\Delta x^2)$  accuracy

Results, force on lid



$$\frac{\partial u}{\partial t} = -u \frac{\partial u}{\partial x} - v \frac{\partial u}{\partial y} - \frac{\partial p}{\partial x} + \frac{1}{Re} \nabla^2 u$$

and similar for  $\nu$ .

$$\frac{\partial u}{\partial t} = -u \frac{\partial u}{\partial x} - v \frac{\partial u}{\partial y} - \frac{\partial p}{\partial x} + \frac{1}{Re} \nabla^2 u$$

and similar for v.

With pressure so that

$$\frac{\partial u}{\partial x} + \frac{\partial v}{\partial y} = 0$$

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p from where?

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 all time  $ightarrow 
abla \cdot \left( rac{\partial \mathbf{u}}{\partial t} 
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 all time  $\rightarrow \nabla \cdot \left( \frac{\partial \mathbf{u}}{\partial t} \right) = 0$ 

Taking the divergence of the momentum equation

$$\nabla \cdot \left(\frac{\partial \mathbf{u}}{\partial t}\right) = -\frac{\partial u_i}{\partial x_j} \frac{\partial u_j}{\partial x_i} - \mathbf{u} \cdot \nabla(\nabla \cdot \mathbf{u}) - \nabla^2 p + \frac{1}{Re} \nabla^2 (\nabla \cdot \mathbf{u})$$

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0

i.e.

$$\nabla^2 p = -\frac{\partial u_i}{\partial x_i} \frac{\partial u_j}{\partial x_i}$$

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NB: Poisson problem unavoidable

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But boundary condition on p?

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$$\frac{\partial u}{\partial t} + u \frac{\partial u}{\partial x} + v \frac{\partial u}{\partial y} = -\frac{\partial p}{\partial x} + \frac{1}{Re} \left( \frac{\partial^2 u}{\partial x^2} + \frac{\partial^2 u}{\partial y^2} \right)$$

Normal component of the momentum equation at a boundary, e.g. on x=0 where u=0 all y and t

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Normal component of the momentum equation at a boundary, e.g. on x=0 where u=0 all y and t

$$\begin{array}{cccc} \frac{\partial u}{\partial t} + & u & \frac{\partial u}{\partial x} + v & \frac{\partial u}{\partial y} & = -\frac{\partial p}{\partial x} + \frac{1}{Re} \left( & \frac{\partial^2 u}{\partial x^2} + & \frac{\partial^2 u}{\partial y^2} \right) \\ 0 & 0 & 0 & 0 \end{array}$$

i.e.

$$\frac{\partial p}{\partial n} = \frac{1}{Re} \frac{\partial^2 u_n}{\partial n^2}$$

Normal component of the momentum equation at a boundary, e.g. on x=0 where u=0 all y and t

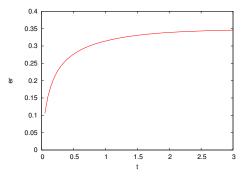
$$\frac{\partial u}{\partial t} + u \frac{\partial u}{\partial x} + v \frac{\partial u}{\partial y} = -\frac{\partial p}{\partial x} + \frac{1}{Re} \left( \frac{\partial^2 u}{\partial x^2} + \frac{\partial^2 u}{\partial y^2} \right)$$

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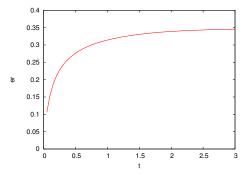
$$\frac{\partial p}{\partial n} = \frac{1}{Re} \frac{\partial^2 u_n}{\partial n^2}$$

NB: pressure arbitrary to additive constant

Error is satisfying  $\nabla \cdot \mathbf{u} = 0$  (N = 20) as function of t

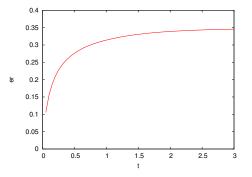


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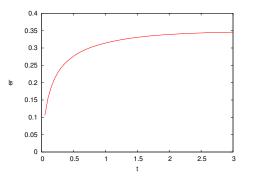
Also strange flow.

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Also strange flow. Coding error? Independent of  $\Delta t$ , small increase with N.

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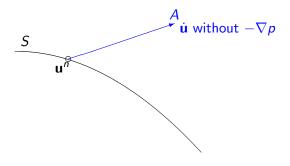
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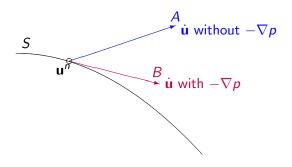
Pressure equation assumes  $\nabla \cdot \mathbf{u}^n = 0$ , and does not correct if untrue, so error accumulates.

In space of all  $\mathbf{u}(\mathbf{x}, t)$ , solution constrained to surface S where  $\nabla \cdot \mathbf{u} = 0$ 

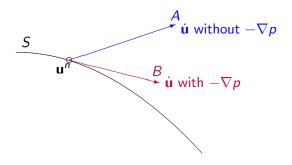
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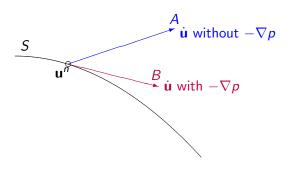


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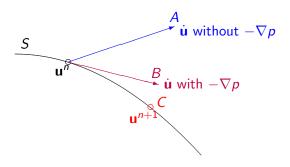


Role of  $\nabla p$  is to project out component of  $\partial \mathbf{u}/\partial t$  normal to surface.

Forward time stepping  $O(\Delta t) o$ slow drift away from surface

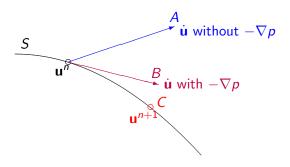


Forward time stepping  $O(\Delta t) o$ slow drift away from surface



Avoid slow accumulation of errors by *projecting* **u** *back onto surface at* **C** 

Forward time stepping  $O(\Delta t) o {\sf slow}$  drift away from surface



Avoid slow accumulation of errors by projecting  ${\bf u}$  back onto surface at  ${\bf C}$ 

Implemented by split time step.

First part (no  $\nabla p$ )

$$\mathbf{u}^* = \mathbf{u}^n + \Delta t \left( -\mathbf{u}^n \cdot \nabla \mathbf{u}^n + \frac{1}{Re} \nabla^2 \mathbf{u}^n \right).$$

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So solve

$$\nabla^2 p = \frac{1}{\Delta t} \nabla \cdot \mathbf{u}^*.$$

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and then evaluate  $\mathbf{u}^{n+1}$ .

### First part (no $\nabla p$ )

At interior points

$$\begin{array}{lll} u_{ij}^{*} & = & u_{ij}^{n} \\ & + & \Delta t \left( -u_{ij}^{n} \frac{u_{i+1j}^{n} - u_{i-1j}^{n}}{2\Delta x} - v_{ij}^{n} \frac{u_{ij+1}^{n} - u_{ij-1}^{n}}{2\Delta x} \right) \\ & + & \frac{\Delta t}{Re\Delta x^{2}} \begin{pmatrix} 1 \\ 1 & -4 & 1 \\ 1 \end{pmatrix} u_{ij}^{n}, \end{array}$$

and a similar expression for  $v_{ij}^*$ .

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Use BC on u and v.

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Several algorithms for the projection step.

#### Projection step - algorithm 2

$$\frac{\Delta t}{\Delta x^2} \begin{pmatrix} 1 & 1 \\ 1 & -4 & 1 \\ 1 & 1 \end{pmatrix} p_{ij} = \frac{u_{i+1j}^* - u_{i-1j}^*}{2\Delta x} + \frac{v_{ij+1}^* - v_{ij-1}^*}{2\Delta x}$$

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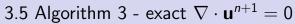
Does not quite give the desired  $\nabla \cdot \mathbf{u}^{n+1} = 0$ :

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Does not quite give the desired  $\nabla \cdot \mathbf{u}^{n+1} = 0$ :

has a small error which tends to zero as  $\Delta x \rightarrow 0$ .



#### 3.5 Algorithm 3 - exact $\nabla \cdot \mathbf{u}^{n+1} = 0$

Now with our central differencing

$$\left. \frac{\partial u^{n+1}}{\partial x} \right|_{::} = \frac{u_{i+1j}^{n+1} - u_{i-1j}^{n+1}}{2\Delta x}$$

$$\frac{\partial x}{\partial x}\Big|_{ij} = \frac{1}{2\Delta x}$$

$$\left(u_{i+1j}^* - \Delta t \frac{p_{i+2j} - p_{ij}}{2\Delta x}\right) - \left(u_{i-1j}^* - \Delta t \frac{p_{ij} - p_{ij}}{2\Delta x}\right)$$

$$= \frac{\left(u_{i+1j}^* - \Delta t \frac{p_{i+2j} - p_{ij}}{2\Delta x}\right) - \left(u_{i-1j}^* - \Delta t \frac{p_{ij} - p_{i-2j}}{2\Delta x}\right)}{2\Delta x},$$

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and similarly for  $\partial v^{n+1}/\partial y$ .

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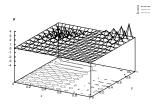
$$\frac{\partial u^{n+1}}{\partial x}\Big|_{ij} = \frac{u_{i+1j}^{n+1} - u_{i-1j}^{n+1}}{2\Delta x} \\
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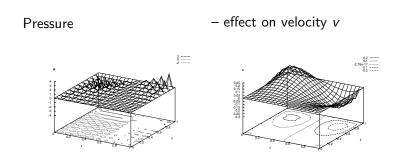
and similarly for  $\partial v^{n+1}/\partial y$ .

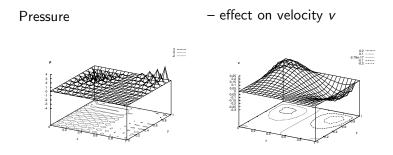
Hence pressure should satisfy  $(recall \ f'' \neq (f')')$ 

$$\frac{\Delta t}{4\Delta x^2} \begin{pmatrix} 1 & 1 & 0 \\ 1 & 0 & -4 & 0 & 1 \\ 0 & 0 & 1 & 0 \end{pmatrix} p_{ij} = \begin{pmatrix} u_{i+1j}^* - u_{i-1j}^* & v_{ij+1}^* - v_{ij-1}^* \\ 2\Delta x & 2\Delta x \end{pmatrix}.$$

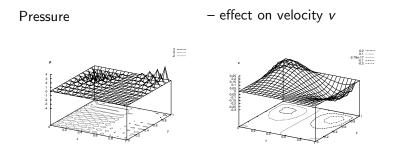
#### Pressure



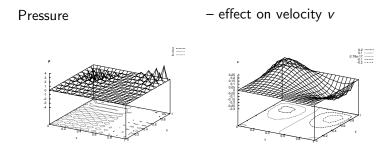




Two uncoupled solutions for pressure on odd/even i+j



Two uncoupled solutions for pressure on odd/even i+j Spurious mode + - + - + has no  $\nabla p$ 



Two uncoupled solutions for pressure on odd/even i+j Spurious mode + - + - has no  $\nabla p$  Also errors 4 times larger from wide span molecule

3.6 Staggered grid - algorithm 4

New idea -

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New idea – a staggered grid with different variables at different locations



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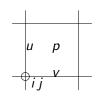
New idea – a staggered grid with different variables at different locations



Write  $u_{ij+\frac{1}{2}}$ ,  $v_{i+\frac{1}{2}j}$  and  $p_{i+\frac{1}{2}j+\frac{1}{2}}$ 

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New idea – a staggered grid with different variables at different locations



Write  $u_{ij+\frac{1}{2}}$ ,  $v_{i+\frac{1}{2}j}$  and  $p_{i+\frac{1}{2}j+\frac{1}{2}}$ 

Good for central differencing

Boundaries coincide with mass flow BC

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$$u_{0j+\frac{1}{2}} = u_{Nj+\frac{1}{2}} = 0$$
  
$$v_{i+\frac{1}{2}0} = v_{i+\frac{1}{2}N} = 0$$

for j=0 
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The tangential component of velocity is held half a grid block away Use one point outside

$$\begin{split} v_{-\frac{1}{2}j} &= -v_{\frac{1}{2}j} \quad \text{and} \quad v_{N+\frac{1}{2}j} &= -v_{N-\frac{1}{2}j} \\ u_{i-\frac{1}{2}} &= -u_{i\frac{1}{2}} \quad \text{and} \quad u_{iN+\frac{1}{2}} &= 2\sin^2(i*\Delta x) - u_{iN-\frac{1}{2}} \end{split}$$

for  $j=1 o {\it N}-1$  and  $i=1 o {\it N}-1$  respectively

$$u_{ij+\frac{1}{2}}^* = u_{ij+\frac{1}{2}}^n$$

$$\begin{array}{rcl} u_{ij+\frac{1}{2}}^* & = & u_{ij+\frac{1}{2}}^n \\ & - & \Delta t \, u_{ij+\frac{1}{2}}^n \frac{u_{i+1j+\frac{1}{2}}^n - u_{i-1j+\frac{1}{2}}^n}{2\Delta x} \end{array}$$

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## Incompressibility at $i + \frac{1}{2}j + \frac{1}{2}$

Compact

$$\frac{\Delta t}{\Delta x^2} \begin{pmatrix} 1 & 1 \\ 1 & -4 & 1 \\ 1 & 1 \end{pmatrix} p_{i+\frac{1}{2}j+\frac{1}{2}} = \frac{u_{i+1j+\frac{1}{2}}^* - u_{ij+\frac{1}{2}}^* - u_{i+\frac{1}{2}j+1}^* - v_{i+\frac{1}{2}j}^*}{\Delta x} + \frac{v_{i+\frac{1}{2}j+1}^* - v_{i+\frac{1}{2}j}^*}{\Delta x}.$$

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Pressure boundary condition at  $O(\Delta x^2)$ 

$$p_{-\frac{1}{2}j+\frac{1}{2}} = p_{\frac{1}{2}j+\frac{1}{2}} + \frac{1}{Re} \left( \frac{-u_{3j+\frac{1}{2}} + 4u_{2,j+\frac{1}{2}} - 5u_{1j+\frac{1}{2}} + 2u_{0j+\frac{1}{2}}}{\Delta x} \right),$$

on left boundary and similar others

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$$i + \frac{1}{2}j + \frac{1}{2}$$

Compact

$$\frac{\Delta t}{\Delta x^2} \begin{pmatrix} 1 & 1 \\ 1 & -4 & 1 \\ 1 & 1 \end{pmatrix} p_{i+\frac{1}{2}j+\frac{1}{2}} = \frac{u_{i+1j+\frac{1}{2}}^* - u_{ij+\frac{1}{2}}^* - u_{i+\frac{1}{2}j+1}^* - v_{i+\frac{1}{2}j}^*}{\Delta x}.$$

Pressure boundary condition at  $O(\Delta x^2)$ 

$$p_{-\frac{1}{2}j+\frac{1}{2}} = p_{\frac{1}{2}j+\frac{1}{2}} + \frac{1}{Re} \left( \frac{-u_{3j+\frac{1}{2}} + 4u_{2,j+\frac{1}{2}} - 5u_{1j+\frac{1}{2}} + 2u_{0j+\frac{1}{2}}}{\Delta x} \right),$$

on left boundary and similar others

NB: On boundary need to advance normal component of  $\mathbf{u}^n$  to nonzero  $\mathbf{u}^*$  and the apply pressure projection to  $\mathbf{u}^{n+1}$  back to zero, to avoid erroneously making  $\partial p/\partial n=0$ 

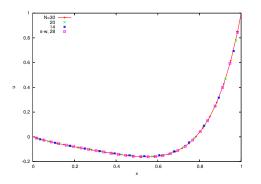
#### 3.7 Results for algorithm 4

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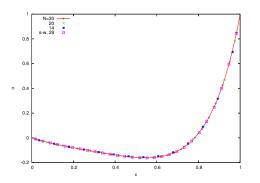


Re=10 and N=14, 20 and 30. Also result from  $\psi-\omega$  formulation at N=28.

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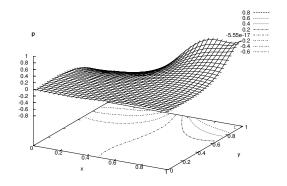


Re = 10 and N = 14, 20 and 30.

Also result from  $\psi - \omega$  formulation at N = 28.

VERY IMPORTANT - agrees

#### Pressure



N = 30

Force on lid

$$F = \sum_{i=1}^{N-1} \frac{u_{i N + \frac{1}{2}} - u_{i N - \frac{1}{2}}}{\Delta x} \times \Delta x + O(\Delta x^{2}).$$

Force on lid

$$F = \sum_{i=1}^{N-1} \frac{u_{iN+\frac{1}{2}} - u_{iN-\frac{1}{2}}}{\Delta x} \times \Delta x + O(\Delta x^2).$$

Similarly the viscous force on the bottom

Force on lid

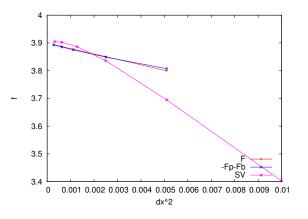
$$F = \sum_{i=1}^{N-1} \frac{u_{i N + \frac{1}{2}} - u_{i N - \frac{1}{2}}}{\Delta x} \times \Delta x + O(\Delta x^{2}).$$

Similarly the viscous force on the bottom

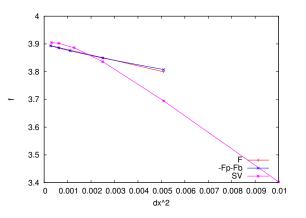
Pressure force on sides from

$$\sum_{j=0}^{N-1} \frac{1}{2} \left( -p_{-\frac{1}{2}j+\frac{1}{2}} - p_{\frac{1}{2}j+\frac{1}{2}} + p_{N-\frac{1}{2}j+\frac{1}{2}} + p_{N+\frac{1}{2}j+\frac{1}{2}} \right) \times \Delta x + O(\Delta x^2).$$

Steady force at Re = 10, as function of  $\Delta x^2$ 



Steady force at Re = 10, as function of  $\Delta x^2$ 



$$F = 3.8998 \pm 0.0002$$

and force on bottom is -0.254