Resumé: Driven cavity, ψ – ω formulation

Poisson problem: $\nabla^2 \psi = -\omega$, SOR

Vorticity evolution:
$$\frac{\partial \omega}{\partial t} = -\frac{\partial(\omega,\psi)}{\partial(x,y)} + \frac{1}{Re}\nabla^2\omega$$

BC for ω

Timestep instability
$$\rightarrow \Delta t = \frac{1}{5} Re \Delta x^2$$

Check $O(\Delta x^2)$ accuracy

Results, force on lid

3. Primitive variable formulation, u, v, p

$$\frac{\partial u}{\partial t} = -u \frac{\partial u}{\partial x} - v \frac{\partial u}{\partial y} - \frac{\partial p}{\partial x} + \frac{1}{Re} \nabla^2 u$$

and similar for v.

With pressure so that

$$\frac{\partial u}{\partial x} + \frac{\partial v}{\partial y} = 0$$

Start from rest, with u and v given on the boundary.

p from where?

3.2 Pressure equation

$$abla \cdot \mathbf{u} = 0$$
 all time $ightarrow
abla \cdot \left(rac{\partial \mathbf{u}}{\partial t}
ight) = 0$

Taking the divergence of the momentum equation

$$\nabla \cdot \left(\frac{\partial \mathbf{u}}{\partial t}\right) = -\frac{\partial u_i}{\partial x_j} \frac{\partial u_j}{\partial x_i} - \mathbf{u} \cdot \nabla(\nabla \cdot \mathbf{u}) - \nabla^2 p + \frac{1}{Re} \nabla^2 (\nabla \cdot \mathbf{u})$$
0

i.e.

$$\nabla^2 p = -\frac{\partial u_i}{\partial x_j} \frac{\partial u_j}{\partial x_i}$$

NB: Poisson problem unavoidable

But boundary condition on p?

Boundary condition

Normal component of the momentum equation at a boundary, e.g. on x=0 where u=0 all y and t

$$\frac{\partial u}{\partial t} + u \frac{\partial u}{\partial x} + v \frac{\partial u}{\partial y} = -\frac{\partial p}{\partial x} + \frac{1}{Re} \left(\frac{\partial^2 u}{\partial x^2} + \frac{\partial^2 u}{\partial y^2} \right)$$

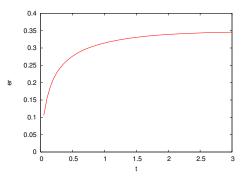
i.e.

$$\frac{\partial p}{\partial n} = \frac{1}{Re} \frac{\partial^2 u_n}{\partial n^2}$$

NB: pressure arbitrary to additive constant

Algorithm 1 (pressure equation) FAILS

Error is satisfying $\nabla \cdot \mathbf{u} = 0$ (N = 20) as function of t



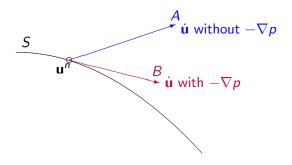
Also strange flow.

Coding error? Independent of Δt , small increase with N.

Pressure equation assumes $\nabla \cdot \mathbf{u}^n = 0$, and does not correct if untrue, so error accumulates.

3.3 Incompressibility as a constraint

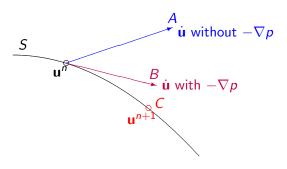
In space of all $\mathbf{u}(\mathbf{x},t)$, solution constrained to surface S where $\nabla \cdot \mathbf{u} = 0$



Role of ∇p is to project out component of $\partial \mathbf{u}/\partial t$ normal to surface.

... incompressibility as a constraint

Forward time stepping $O(\Delta t) o {\sf slow}$ drift away from surface



Avoid slow accumulation of errors by projecting ${\bf u}$ back onto surface at ${\bf C}$

Implemented by split time step.

3.4 Split time step

First part (no ∇p)

$$\mathbf{u}^* = \mathbf{u}^n + \Delta t \left(-\mathbf{u}^n \cdot \nabla \mathbf{u}^n + \frac{1}{Re} \nabla^2 \mathbf{u}^n \right).$$

Second projection part

$$\mathbf{u}^{n+1} = \mathbf{u}^* - \Delta t \nabla p,$$

where need $\nabla \cdot \mathbf{u}^{n+1} = 0$.

So solve

$$\nabla^2 p = \frac{1}{\Delta t} \nabla \cdot \mathbf{u}^*.$$

and then evaluate \mathbf{u}^{n+1} .

First part (no ∇p)

At interior points

$$\begin{array}{lcl} u_{ij}^{*} & = & u_{ij}^{n} \\ & + & \Delta t \left(-u_{ij}^{n} \frac{u_{i+1j}^{n} - u_{i-1j}^{n}}{2\Delta x} - v_{ij}^{n} \frac{u_{ij+1}^{n} - u_{ij-1}^{n}}{2\Delta x} \right) \\ & + & \frac{\Delta t}{Re\Delta x^{2}} \begin{pmatrix} 1 \\ 1 & -4 & 1 \\ 1 & 1 \end{pmatrix} u_{ij}^{n}, \end{array}$$

and a similar expression for v_{ij}^* .

Use BC on u and v.

Several algorithms for the projection step.

Projection step – algorithm 2

$$\frac{\Delta t}{\Delta x^2} \begin{pmatrix} 1 & 1 \\ 1 & -4 & 1 \\ 1 & 1 \end{pmatrix} p_{ij} = \frac{u_{i+1j}^* - u_{i-1j}^*}{2\Delta x} + \frac{v_{ij+1}^* - v_{ij-1}^*}{2\Delta x}$$

Does not quite give the desired $\nabla \cdot \mathbf{u}^{n+1} = 0$:

has a small error which tends to zero as $\Delta x \rightarrow 0$.

3.5 Algorithm 3 - exact $\nabla \cdot \mathbf{u}^{n+1} = 0$

Now with our central differencing

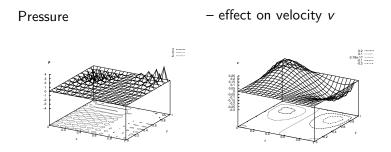
$$\frac{\partial u^{n+1}}{\partial x}\Big|_{ij} = \frac{u_{i+1j}^{n+1} - u_{i-1j}^{n+1}}{2\Delta x} \\
= \frac{\left(u_{i+1j}^* - \Delta t \frac{p_{i+2j} - p_{ij}}{2\Delta x}\right) - \left(u_{i-1j}^* - \Delta t \frac{p_{ij} - p_{i-2j}}{2\Delta x}\right)}{2\Delta x},$$

and similarly for $\partial v^{n+1}/\partial y$.

Hence pressure should satisfy $(recall \ f'' \neq (f')')$

$$\frac{\Delta t}{4\Delta x^2} \begin{pmatrix} 1 & 1 & 0 \\ 1 & 0 & -4 & 0 & 1 \\ 0 & 0 & 1 & 0 \end{pmatrix} p_{ij} = \left(\frac{u_{i+1j}^* - u_{i-1j}^*}{2\Delta x} + \frac{v_{ij+1}^* - v_{ij-1}^*}{2\Delta x}\right).$$

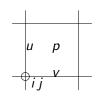
Problem – spurious pressure modes



Two uncoupled solutions for pressure on odd/even i+j Spurious mode + - + - has no ∇p Also errors 4 times larger from wide span molecule

3.6 Staggered grid – algorithm 4

New idea – a staggered grid with different variables at different locations



Write $u_{ij+\frac{1}{2}}$, $v_{i+\frac{1}{2}j}$ and $p_{i+\frac{1}{2}j+\frac{1}{2}}$

Good for central differencing

Boundary conditions

Boundaries coincide with mass flow BC

$$u_{0j+\frac{1}{2}} = u_{Nj+\frac{1}{2}} = 0$$
$$v_{i+\frac{1}{2}0} = v_{i+\frac{1}{2}N} = 0$$

for $j=0
ightarrow {\it N}-1$ and $i=0
ightarrow {\it N}-1$ respectively

The tangential component of velocity is held half a grid block away Use one point outside

$$\begin{split} v_{-\frac{1}{2}j} &= -v_{\frac{1}{2}j} \quad \text{and} \quad v_{N+\frac{1}{2}j} &= -v_{N-\frac{1}{2}j} \\ u_{i-\frac{1}{2}} &= -u_{i\frac{1}{2}} \quad \text{and} \quad u_{iN+\frac{1}{2}} &= 2\sin^2(i*\Delta x) - u_{iN-\frac{1}{2}} \end{split}$$

for $j=1 o {\it N}-1$ and $i=1 o {\it N}-1$ respectively

Momentum equation at $ij + \frac{1}{2}$

First part of split time step (without pressure)

$$\begin{array}{lll} u_{ij+\frac{1}{2}}^* & = & u_{ij+\frac{1}{2}}^n \\ & - & \Delta t \, u_{ij+\frac{1}{2}}^n \frac{u_{i+1j+\frac{1}{2}}^n - u_{i-1j+\frac{1}{2}}^n}{2\Delta x} \\ & - & \Delta t \frac{1}{4} \left(v_{i+\frac{1}{2}j}^n + v_{i-\frac{1}{2}j}^n + v_{i+\frac{1}{2}j+1}^n + v_{i-\frac{1}{2}j+1}^n \right) \frac{u_{ij+\frac{3}{2}}^n - u_{ij-\frac{1}{2}}^n}{2\Delta x} \\ & + & \frac{\Delta t}{Re\Delta x^2} \begin{pmatrix} 1 \\ 1 & -4 & 1 \\ 1 & \end{pmatrix} u_{ij+\frac{1}{2}}^n, \end{array}$$

Incompressibility at
$$i + \frac{1}{2}j + \frac{1}{2}$$

Compact

$$\frac{\Delta t}{\Delta x^2} \begin{pmatrix} 1 \\ 1 \\ -4 \\ 1 \end{pmatrix} p_{i+\frac{1}{2}j+\frac{1}{2}} = \frac{u_{i+1j+\frac{1}{2}}^* - u_{ij+\frac{1}{2}}^* - u_{i+\frac{1}{2}j+1}^* - v_{i+\frac{1}{2}j}^*}{2\Delta x}.$$

Pressure boundary condition at $O(\Delta x^2)$

$$p_{-\frac{1}{2}j+\frac{1}{2}} = p_{\frac{1}{2}j+\frac{1}{2}} + \frac{1}{Re} \left(\frac{-u_{3j+\frac{1}{2}} + 4u_{2,j+\frac{1}{2}} - 5u_{1j+\frac{1}{2}} + 2u_{0j+\frac{1}{2}}}{\Delta x} \right),$$

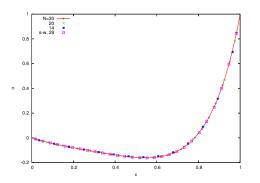
on left boundary and similar others

NB: On boundary need to advance normal component of \mathbf{u}^n to nonzero \mathbf{u}^* and the apply pressure projection to \mathbf{u}^{n+1} back to zero, to avoid erroneously making $\partial p/\partial n=0$

3.7 Results for algorithm 4

First check consistency of accuracy.

Steady horizontal velocity at $x = \frac{1}{2}$



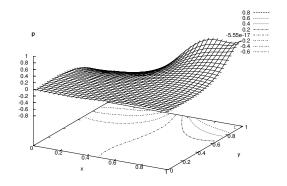
Re = 10 and N = 14, 20 and 30.

Also result from $\psi - \omega$ formulation at N = 28.

VERY IMPORTANT - agrees

... results

Pressure



N = 30

... results

Force on lid

$$F = \sum_{i=1}^{N-1} \frac{u_{i N + \frac{1}{2}} - u_{i N - \frac{1}{2}}}{\Delta x} \times \Delta x + O(\Delta x^{2}).$$

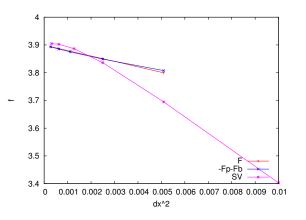
Similarly the viscous force on the bottom

Pressure force on sides from

$$\sum_{j=0}^{N-1} \frac{1}{2} \left(-p_{-\frac{1}{2}j+\frac{1}{2}} - p_{\frac{1}{2}j+\frac{1}{2}} + p_{N-\frac{1}{2}j+\frac{1}{2}} + p_{N+\frac{1}{2}j+\frac{1}{2}} \right) \times \Delta x + O(\Delta x^2).$$

... results

Steady force at Re = 10, as function of Δx^2



$$F = 3.8998 \pm 0.0002$$

and force on bottom is -0.254