

Chapter 3

Chapter 1: the phenomena.

Chapter 2: measuring intrinsic properties, e.g. viscosity and elasticity.

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Now need to encapsulate those properties in governing equations.

Conservation equations – true all materials

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- ▶ Need Constitutive (material dependent) Relation between stress $\boldsymbol{\sigma}$ and flow \mathbf{u} .

Constitutive equations

Phenomenology

- 'Simple' materials

- Perfectly elastic material

- Time derivatives

Exact approximations

- Linear viscoelasticity

- Second-order fluid

Semi-empirical models

- Generalised Newtonian

- Oldroyd-B

- K-BKZ

'Simple' materials

Lagrangian description

$$\mathbf{X} \rightarrow \mathbf{x}(\mathbf{X}, t)$$

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Deformation of line element (for micro-lengths \ll macro-lengths)

$$\delta \mathbf{X} \rightarrow \delta \mathbf{x} = \mathbf{A} \cdot \delta \mathbf{X}, \quad A_{ij} = \frac{\partial x_i}{\partial X_j}$$

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This functional dependence not useful, except for fast elastic limit and slow viscous limits (each with single parameter)

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Require $\sigma \{A\}$ to obey this identity for all $Q(t)$.

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In incompressible material with isotropic in rest state,

$$f(U^2) = U^2 f_1 + U^{-2} f_2, \quad f_i \text{ scalar functions of invariants of } U,$$

so

$$\sigma = A A^T f_1 + A^{-1 T} A^{-1} f_2$$

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Virtual work and σ co-diagonal with U gives

$$\sigma_1 = \frac{1}{\lambda_2 \lambda_3} \left(\frac{\partial w}{\partial \lambda_1} = \lambda_1 \frac{\partial w}{\partial \alpha} - \lambda_1^{-3} \frac{\partial w}{\partial \beta} \right),$$

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Better for data fitting – Ogden model: $w(\lambda_1^n + \lambda_2^n + \lambda_3^n)$

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$$\text{strain-rate } E' = QEQ^T, \quad \text{vorticity (tensor) } \Omega' = Q\Omega Q^T - \dot{Q}Q^T$$

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This is the rate of change of σ seen by an observer rotating with the vorticity, and so is universal.

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Recall rotation frames

$$\overset{\circ}{\dot{\mathbf{x}}} = \dot{\mathbf{x}} + \Omega \mathbf{x}$$

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Recall stretching material line element

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Recall stretching material line element

$$\dot{\delta l} = \delta l \cdot \nabla u$$

so for second-order tensor

$$\delta \dot{l} \delta l = \nabla u^T \cdot \delta l \delta l + \delta l \delta l \cdot \nabla u$$

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For a Newtonian viscous fluid $G(s) = \delta(s)$ and for an elastic solid $G(s) = 1$.

Linear viscoelasticity 2

Student Exercise: If $G(t)$ has a single exponential decay,

$$G(t) = G_0 e^{-t/\tau}$$

show that a polar plot of $Re(G^*)$ versus $Im(G^*)$ as (real) ω varies is part of a circle.

Linear viscoelasticity 3

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Hence recoil after stop steady shear flow $\dot{\gamma}_0$

$$-\dot{\gamma}_0 \frac{\int_0^{\infty} sG(s) ds}{\int_0^{\infty} G(s) ds}$$

Student Exercise

Second-order fluid

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Dangerous in stability analyses and numerical studies, where bad behaviour can occur outside limitation of weak and slowly varying.

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- ▶ $\mu_{\text{ext}} = \mu + \left(\alpha + \frac{1}{4}\beta\right) \dot{\epsilon}$

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- ▶ but must keep last term small

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- ▶ Carreau, Yasuda & Cross

$$\mu = \mu_{\infty} + (\mu_0 - \mu_{\infty}) (1 + (\tau\dot{\gamma})^a)^{(n-1)/a}$$

with plateaux at high and low $\dot{\gamma}$.

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- ▶ **Bingham**

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► Herchel-Buckley

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$$\mu_{\text{ext}} = \mu_0 \frac{1 - \lambda_2 \dot{\epsilon} - 2\lambda_1 \lambda_2 \dot{\epsilon}^2}{(1 - 2\lambda_1 \dot{\epsilon})(1 + \lambda_1 \dot{\epsilon})}$$

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- ▶ becomes negative just above $\dot{\epsilon} = 1/2\lambda_1$!!!!

Variants of Oldroyd-B

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- ▶ **Multi-mode** versions of above

Molecular reformulation of Oldroyd-B

also for better numerics

Microstructure A :

$$\overset{\nabla}{A} + \frac{f}{\tau} (A - I) = 0$$

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Oldroyd-B $f = 1$

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FENE modification, for nice behaviour in extensional flow

$$f = \frac{L^2}{L^2 - \text{trace } A}$$

K-BKZ model fluid

Kay-Bernstein-Kearsley-Zappa

History dependence through time integrals

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Kay-Bernstein-Kearsley-Zappa

History dependence through time integrals

Merging of **linear viscoelasticity** and **nonlinear elasticity**

$$\sigma = \int_0^\infty \dot{G}(s) \left[\frac{\partial w}{\partial \alpha} (\tilde{A}\tilde{A}^T - I) - \frac{\partial w}{\partial \beta} (\tilde{A}^{-1T}\tilde{A}^{-1} - I) \right] ds$$

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functions of combinations α and β eigenvalues of \tilde{A} .

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In extensional flow

$$\mu_{\text{ext}} = \int_0^{\infty} \dot{G}(s) [\phi_1 (e^{2\dot{\epsilon}s} - e^{-\dot{\epsilon}s}) + \phi_2 (e^{\dot{\epsilon}s} - e^{-2\dot{\epsilon}s})] s \, ds / \dot{\epsilon}$$

Wagner model

$$\phi_2 = 0 \quad \text{so } N_2 = 0$$

and

$$\phi_1 = \exp\left(-k\sqrt{\alpha - 3 + \theta(\beta - \alpha)}\right)$$

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In shear shear

$$\phi_1 = \exp(-k\dot{\gamma}(t - s))$$