## Mathematical Tripos Part III Black Holes: Examples Sheet 2

1. A general static, spherically symmetric metric can be written

$$ds^{2} = -A(r)dt^{2} + \frac{dr^{2}}{B(r)} + r^{2}d\Omega^{2},$$

where  $d\Omega^2$  is the metric on a unit 2-sphere. Assume that A(r) and B(r) are analytic functions of r such that both have a simple zero at  $r = r_+ > 0$  and are positive for  $r > r_+$ .

(a) Show that radial null geodesics are given by  $t \pm r^* = \text{constant}$ , where

$$r^* \equiv \int_{r_0}^r \frac{dx}{\sqrt{A(x)B(x)}}$$

with  $r_0 > r_+$  an arbitrary constant. Show that  $r^* \to -\infty$  as  $r \to r_+$ .

(b) Obtain the metric in ingoing Eddington-Finkelstein coordinates. Explain why this metric can be analytically continued through  $r = r_+$ .

2. Consider a particle with 4-velocity U in a stationary, asymptotically flat, space-time with timelike Killing vector field k.  $E = -k \cdot U$  has the interpretation of "energy per unit mass measured at infinity" if the particle moves on a geodesic. For non-geodesic motion, this equation is used to define "energy per unit mass measured at infinity."

(a) Consider a unit mass particle P following an orbit of k at radius  $r = r_P > 2M$  in the Schwarzschild spacetime. Assume that the force making this particle accelerate comes from a radial massless inelastic string, whose other end is held by an observer Q at infinity. If Q pulls the string through proper distance  $\delta S$  then what is the change  $\delta r_P$  in  $r_P$ ?

(b) What is the change  $\delta E$  in the energy of P measured by Q? This must equal the work  $F\delta S$  done by Q where F is the force that the string exerts on Q, i.e., the tension at Q. Calculate F. Show that  $F \to 1/(4M)$  as  $r_P \to 2M$ . What is the force measured by P as  $r_P \to 2M$ ?

- 3. Use isotropic coordinates to prove that a surface of constant t in the Schwarzschild spacetime is an asymptotically flat end with  $K_{ab} = 0$ .
- 4. (a) Let (M, g) be the 2d Einstein static Universe with metric  $ds^2 = -dt^2 + d\phi^2$  where  $\phi \sim \phi + 2\pi$ . Let S be the surface t = 0. Determine  $D^+(S)$  and  $J^+(S)$ . (b) Do the same where (M, g) is now the spacetime obtained by deleting the point  $t = \phi = 0$  from the Einstein static Universe. (c) Do the same for the Kruskal spacetime where S is the surface t = 0 in region I.
- 5. A perfect fluid has stress tensor  $T_{ab} = (\rho + p)u_a u_b + pg_{ab}$ , where  $\rho$  is the energy density, p the pressure, and  $u^a$  the 4-velocity of the fluid. Show that
  - (a) the dominant energy condition is obeyed if, and only if,  $\rho \ge |p|$ ;
  - (b) the weak energy condition is obeyed if, and only if,  $\rho \ge 0$  and  $\rho + p \ge 0$ ;
  - (c) the null energy condition is obeyed if, and only if,  $\rho + p \ge 0$ ;
  - (d) the strong energy condition is obeyed if, and only if,  $\rho + 3p \ge 0$  and  $\rho + p \ge 0$ .

A cosmological constant has  $p = -\rho$ . Which energy conditions does it violate? (Consider both signs for  $\rho$ .)

- 6. Consider two Lorentzian metrics on a manifold M related by a conformal transformation  $\bar{g} = \Omega^2 g$ where  $\Omega$  is a positive function on M.
  - (a) Show that g and  $\overline{g}$  have the same null geodesics.
  - (b) Show that the Ricci tensor of g is related to the Ricci tensor of  $\bar{g}$  by

$$R_{ab} = \bar{R}_{ab} + 2\Omega^{-1}\bar{\nabla}_a\bar{\nabla}_b\Omega + \bar{g}_{ab}\bar{g}^{cd}\left(\Omega^{-1}\bar{\nabla}_c\bar{\nabla}_d\Omega - 3\Omega^{-2}\partial_c\Omega\partial_d\Omega\right)$$

where  $\overline{\nabla}$  is the Levi-Civita connection associated with  $\overline{g}$ .

(c) Let  $\psi$  be a solution of the equation

$$g^{ab}\nabla_a\nabla_b\psi + \xi R\psi = 0$$

We say that the equation is *conformally covariant* if there exists a constant p such that  $\bar{\psi} \equiv \Omega^p \psi$  is a solution of the equation in a spacetime with metric  $\bar{g} = \Omega^2 g$  whenever  $\psi$  solves the equation in a spacetime with metric g. Determine the value of  $\xi$  for which this equation is conformally covariant.

7. The Robinson-Bertotti metric is

$$ds^{2} = -\lambda^{2}dt^{2} + M^{2}\left(\frac{d\lambda}{\lambda}\right)^{2} + M^{2}d\Omega^{2}$$

This is the product  $AdS_2 \times S^2$  where  $AdS_2$  denotes 2d anti-de Sitter spacetime. By replacing the time coordinate t by one of the radial null coordinates  $u = t + M/\lambda$ ,  $v = t - M/\lambda$  show that the singularity at  $\lambda = 0$  is merely a coordinate singularity. By introducing the new coordinates (U, V), defined by  $u = \tan(U/2)$ ,  $v = -\cot(V/2)$ , obtain the maximal analytic extension of the RB metric and deduce its Penrose diagram (more precisely: deduce the Penrose diagram of the  $AdS_2$  part of the RB metric). Is this spacetime globally hyperbolic?

8. Determine the Penrose diagram of de Sitter spacetime with metric

$$ds^{2} = -dt^{2} + H^{-2}\cosh^{2}(Ht)(d\chi^{2} + \sin^{2}\chi d\Omega^{2})$$

where H > 0 is a constant and  $0 \le \chi \le \pi$  (( $\chi, \theta, \phi$ ) parameterize a round 3-sphere). (*Hint.* Use a coordinate transformation  $t = t(\eta)$  to bring the metric to a form where it is manifestly conformal to the Einstein static Universe.)

9. Consider a vacuum spacetime that is asymptotically flat at null infinity. In lectures we introduced coordinates  $(u, \Omega, \theta, \phi)$  such that  $\mathcal{I}^+$  is  $\Omega = 0$  and the "unphysical" metric satisfies

$$\bar{g}|_{\Omega=0} = 2dud\Omega + d\theta^2 + \sin^2 d\phi^2$$

and, for small non-zero  $\Omega$ , the corrections to this are  $\mathcal{O}(\Omega)$  except for the uu,  $u\theta$  and  $u\phi$  components which are  $\mathcal{O}(\Omega^2)$ , and the  $\Omega\Omega$  component which vanishes everywhere. The physical metric is  $g = \Omega^{-2}\bar{g}$ . Introduce a new coordinate  $r = 1/\Omega$  and determine the form of the physical metric for large r, keeping track of the size of the subleading corrections. You should find that the uu component is  $\mathcal{O}(1)$ . Show that this component can be set to  $-1 + \mathcal{O}(1/r)$  by a shift  $r \to r + f(u, \theta, \phi)$ . Finally define "asymptotically inertial" coordinates (t, x, y, z) by t = u + rand (x, y, z) related to  $(r, \theta, \phi)$  as for spherical polars. Show that the spacetime metric becomes  $-dt^2 + dx^2 + dy^2 + dz^2$  with corrections that are  $\mathcal{O}(1/r)$  at large r.