#### Computational Methods in Fluid Mechanics

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GIAN 171002M01 at IIT Bombay

#### Course structure

#### Three parts:

- ► Simple Navier-Stokes problem by simple method
  - accuracy, stability, pressure

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  - discretisation, time-stepping, linear algebra

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- Simple Navier-Stokes problem by simple method
  - accuracy, stability, pressure
- ▶ Better treatment of general issues
  - discretisation, time-stepping, linear algebra
- Collection of special topics
  - demo FreeFem, hyperbolic, fast multipoles, free surface

### 1. The driven cavity

Incompressible Navier-Stokes

$$\begin{split} \nabla \cdot \mathbf{u} &= 0, \\ \rho \left( \frac{\partial \mathbf{u}}{\partial t} + \mathbf{u} \cdot \nabla \mathbf{u} \right) &= -\nabla \rho + \mu \nabla^2 \mathbf{u}, \end{split}$$

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2D,  $L \times L$ -box

$$\mathbf{u} = 0$$
 on  $y = 0$  and  $0 < x < L$ , and on  $x = 0$  or  $L$  and  $0 < y < L$ , and  $\mathbf{u} = (U(x), 0)$  on  $y = L$  and  $0 < x < L$ .

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$$\mathbf{u}=0$$
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To find the force on the lid

$$F = \int_0^L \mu \left. \frac{\partial u}{\partial y} \right|_{y=1} dx$$

Before writing any code, need to think about physics

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$$\frac{\partial \mathbf{u}}{\partial t} + \mathbf{u} \cdot \nabla \mathbf{u}$$
  
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 info diffuses, diffusivity  $\nu = \mu/\rho$ , i.e.  $\delta x = \sqrt{\nu \delta t}$ .

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- $\rho \frac{\partial \mathbf{u}}{\partial t} = -\nabla p \quad \text{with} \quad \nabla \cdot \mathbf{u} = 0$  info at  $\infty$  in 0 time, i.e. speed of sound  $= \infty$ .

Before writing any code, need to think about physics Converse, thinking about coding can deepen understanding of physics

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$$\rho \frac{\partial \mathbf{u}}{\partial t} = -\nabla p$$
 with  $\nabla \cdot \mathbf{u} = 0$  info at  $\infty$  in 0 time, i.e. speed of sound  $= \infty$ .

Re ≪ 1 must resolve fast diffusion of vorticity,
 Re ≫ 1 must resolve thin boundary layers,
 we study Re = 10.

What is well-posed?

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▶  $\frac{\partial \phi}{\partial t} + u(x,t) \frac{\partial \phi}{\partial x} = f(x,t)$  — first order hyperbolic Well posed with

IC  $\phi(x,0)$  and

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- $\begin{array}{l} \blacktriangleright \frac{\partial \phi}{\partial t} + u(x,t) \frac{\partial \phi}{\partial x} = f(x,t) \quad \text{ first order hyperbolic} \\ \text{Well posed with} \\ \text{IC } \phi(x,0) \text{ and} \\ \text{inflow BC, e.g. at } x = a \text{ need } \phi(a,t) \text{ if } u(a,t) > 0. \end{array}$
- $\frac{\partial^2 \phi}{\partial t^2} = c^2 \frac{\partial^2 \phi}{\partial x^2}$  second order hyperbolic Well posed with IC  $\phi(x,0)$  and  $\phi_t(x,0)$  and BC at both ends either  $\phi$  or  $\phi_x$  or mixed.

 $lackbox{} 
abla^2\phi=
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Well posed with

BC  $\phi$  or  $\partial \phi/\partial n$  or mixed

 $\nabla^2 \phi = \rho$  - Laplace/Poisson equation, elliptic Well posed with

BC 
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 $ightharpoonup \frac{\partial \phi}{\partial t} = D \frac{\partial^2 \phi}{\partial x^2}$  – Diffusion equation, parabolic

BC at both ends either  $\phi$  or  $\phi_x$  or mixed.

IC  $\phi(x,0)$  and

▶ Naming from quadratic forms

$$ax^2 + bxy + cy^2 + dx + fy + g =$$

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$$a\frac{\partial^{2}\phi}{\partial x^{2}} + b\frac{\partial^{2}\phi}{\partial x \partial y} + c\frac{\partial^{2}\phi}{\partial y^{2}} + d\frac{\partial\phi}{\partial x} + e\frac{\partial\phi}{\partial y} + f\phi = 0$$

Naming from quadratic forms

$$2x^2 + bxy + cy^2 + dy + fy + g$$

$$ax^2 + bxy + cy^2 + dx + fy + g$$

hyperbolic – tough ► elliptic – costly parabolic – safest

Numerically

$$ax^2 + bxy + cy^2 + dx + fy + g$$

 $ax^{2} + bxy + cy^{2} + dx + fy + g = 0$ 

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# Special physics – the corner

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# Special physics – the corner

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- ▶ Better  $\mathbf{u} = (U_0 \sin \pi x / L, 0)$  $\rightarrow \sigma \propto \ln r \rightarrow F$  difficult numerically
- ► Therefore we take  $\mathbf{u} = (U_0 \sin^2 \pi x / L, 0)$

#### Non-dimensionalisation

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Scale u on  $U_0$ , x and y on L, t on  $L/U_0$  and p on  $\rho U_0^2$ . Then

$$Re = \frac{\text{inertial terms } \rho U_0^2 / L}{\text{viscous terms } \mu U_0 / L^2} = \frac{U_0 L}{\nu}.$$

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with BCs

$$\mathbf{u}=0$$
 on  $y=0$  and  $0 < x < 1$ , and on  $x=0$  or  $1$  and  $0 < y < 1$  and  $\mathbf{u}=(\sin^2(\pi x),0)$  on  $y=1$  and  $0 < x < 1$ .

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$${f u} = 0 \quad {\rm on} \ y = 0 \ {\rm and} \ 0 < x < 1, \ {\rm and} \ {\rm on} \ x = 0 \ {\rm or} \ 1 \ {\rm and} \ 0 < y < 1$$
 and  ${f u} = (\sin^2(\pi x), 0) \quad {\rm on} \ y = 1 \ {\rm and} \ 0 < x < 1.$ 

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Finally the force, scaled by  $\mu U_0$ 

$$F = \int_0^1 \frac{\partial u}{\partial y} \bigg|_{y=1} dx.$$

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Methods for relaxing to SS  $\equiv$  pseudo time-stepping.

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#### Two options:

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  - primitive variable formulation

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  - primitive variable formulation
- Eliminate p by forming the vorticity equation
  - streamfunction-vorticity formulation

### 2. Streamfunction-vorticity formulation

Automatically satisfy constraint  $\nabla \cdot \mathbf{u} = 0$  by using the streamfunction representation  $\psi(\mathbf{x}, \mathbf{y})$ 

$$u = \frac{\partial \psi}{\partial y}$$
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In 2D flow vorticity is

$$\omega = \frac{\partial \mathbf{v}}{\partial \mathbf{x}} - \frac{\partial \mathbf{u}}{\partial \mathbf{v}} = -\nabla^2 \psi.$$

Take curl of momentum equation to eliminate p

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No stretching in 2D (first term on RHS)

$$\mathbf{u} \cdot \nabla \omega = \psi_y \omega_x - \psi_x \omega_y = \frac{\partial(\omega, \psi)}{\partial(x, y)}$$

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BC1: 
$$\mathbf{u} \cdot \mathbf{n} = 0$$
 all sides  $\rightarrow$  sides = streamline  $\rightarrow \psi = 0$ .

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$$\frac{\partial \psi}{\partial y} = \sin^2 \pi x$$
 on top  $y = 1$ ,  $0 < x < 1$ 

$$\frac{\partial \psi}{\partial y} = 0$$
 on bottom  $y = 0$ ,  $0 < x < 1$ 

$$\frac{\partial \psi}{\partial x} = 0$$
 on sides  $x = 0$  and 1,  $0 < y < 1$ 

# Solve as decoupled pair

1. At each t given  $\omega$ , find  $\psi$ :

$$\nabla^2 \psi = -\omega$$

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2. With  $\omega$  and now  $\psi$  known at t, find  $\omega$  at  $t + \Delta t$ :

$$\frac{\partial \omega}{\partial t} = -\frac{\partial (\omega, \psi)}{\partial (x, y)} + \frac{1}{Re} \nabla^2 \omega$$

with  $\omega$  on boundary so  $\frac{\partial \psi}{\partial n}$  correct

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with  $\omega$  on boundary so  $\frac{\partial \psi}{\partial n}$  correct  $\rightarrow$  not quite decoupled.