

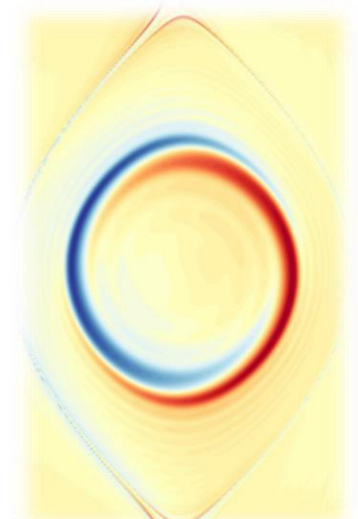
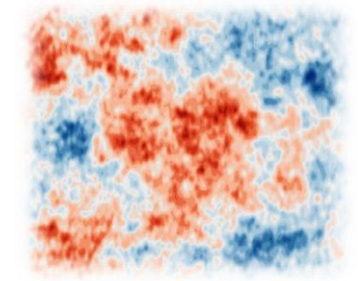
Data-Driven Spectral Measures and Generalized Eigenfunctions of Koopman Operators

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18/06/2024

- C., Townsend, "Rigorous data-driven computation of spectral properties of Koopman operators for dynamical systems" **Communications on Pure and Applied Mathematics**, 2024.
- C., "The mpEDMD Algorithm for Data-Driven Computations of Measure-Preserving Dynamical Systems," **SIAM Journal on Numerical Analysis**, 2023.
- C., Drysdale, Horning, "Rigged Dynamic Mode Decomposition: Data-Driven Generalized Eigenfunction Decompositions for Koopman Operators", arxiv preprint.
- C., "The Multiverse of Dynamic Mode Decomposition Algorithms," **Handbook of Numerical Analysis**, 2024.



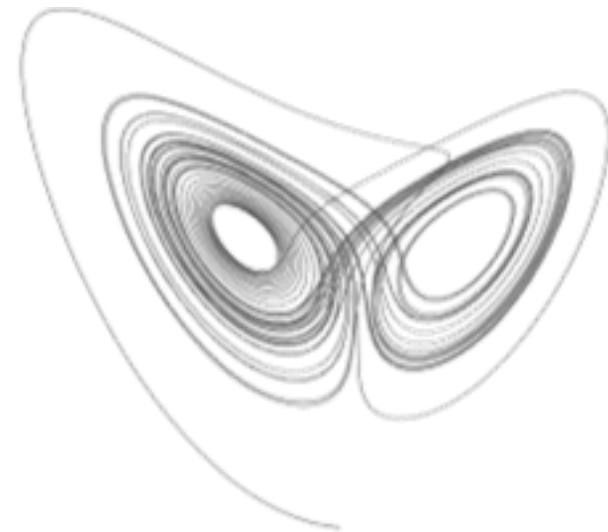
Data-driven dynamical systems

State $x \in \Omega \subseteq \mathbb{R}^d$.

Unknown function $F: \Omega \rightarrow \Omega$ governs dynamics: $x_{n+1} = F(x_n)$.

Goal: Learning from data $\{x^{(m)}, y^{(m)} = F(x^{(m)})\}_{m=1}^M$.

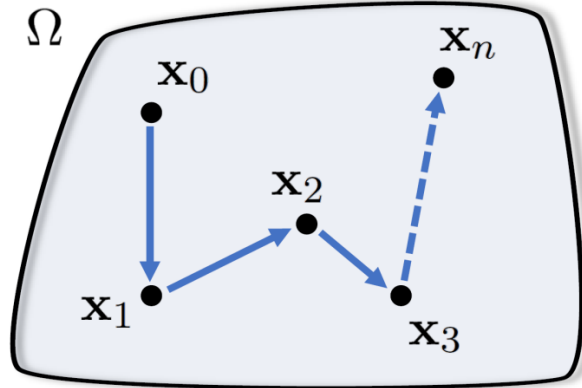
Applications: chemistry, climatology, control, electronics, epidemiology, finance, fluids, molecular dynamics, neuroscience, plasmas, robotics, video processing, etc.



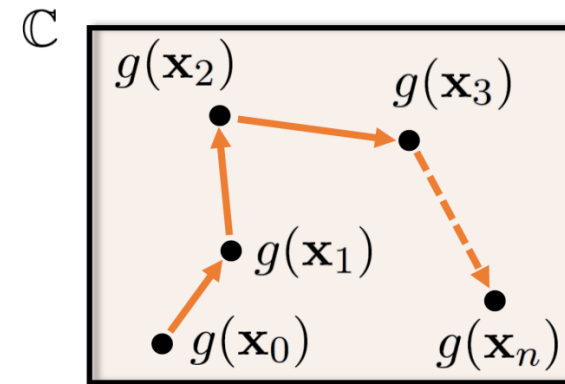
Surveys:

- Brunton, Budišić, Kaiser, Kutz, “*Modern Koopman theory for dynamical systems*,” SIAM Review, 2022.
- Budišić, Mohr, Mezić, “*Applied Koopmanism*,” Chaos, 2012.
- C., “*The Multiverse of Dynamic Mode Decomposition Algorithms*,” Handbook of Numerical Analysis, 2024.

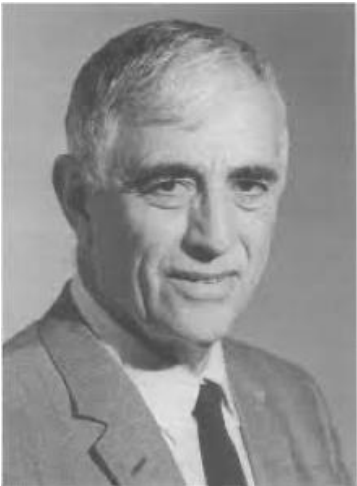
Koopman Operator \mathcal{K} : A global linearization



$g: \Omega \rightarrow \mathbb{C}$
 "observable"



Koopman

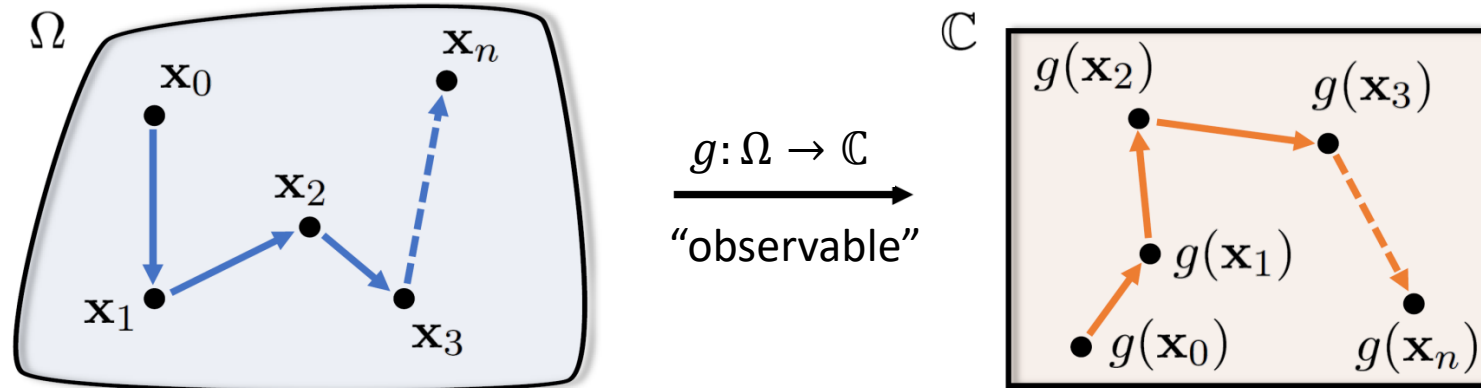


von Neumann



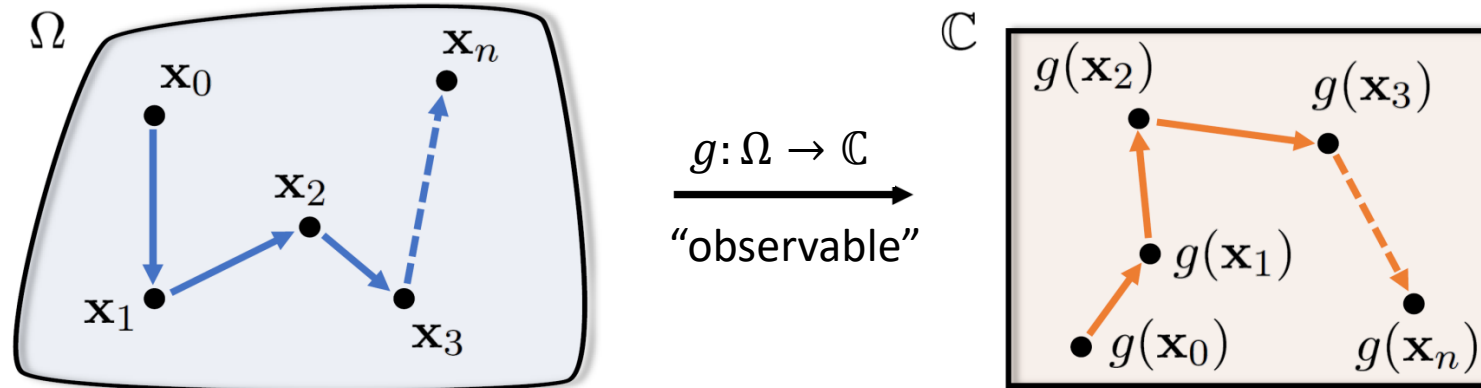
- Koopman, "Hamiltonian systems and transformation in Hilbert space," *Proc. Natl. Acad. Sci. USA*, 1931.
- Koopman, v. Neumann, "Dynamical systems of continuous spectra," *Proc. Natl. Acad. Sci. USA*, 1932.

Koopman Operator \mathcal{K} : A global linearization

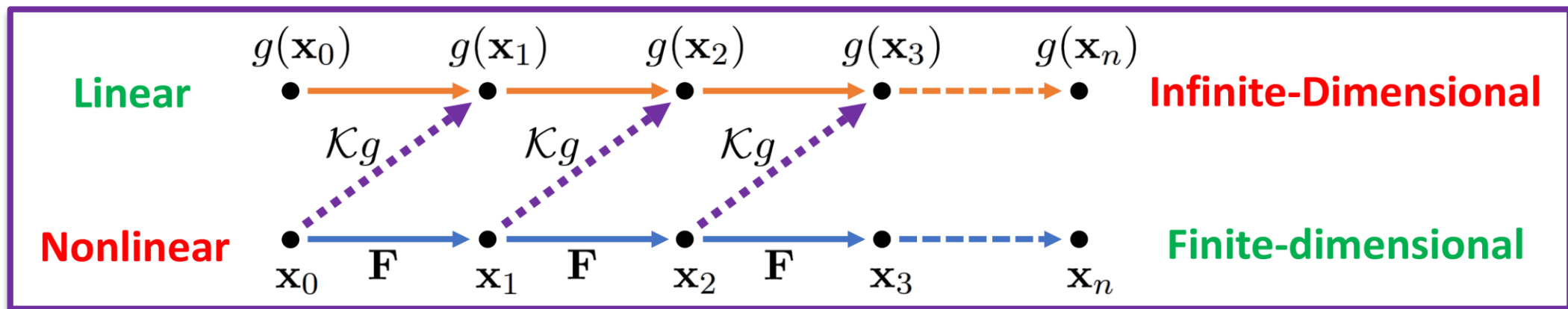


- \mathcal{K} acts on functions $g: \Omega \rightarrow \mathbb{C}$, $[\mathcal{K}g](x) = g(F(x))$.
- **Function space:** $L^2(\Omega, \omega)$, positive measure ω , inner product $\langle \cdot, \cdot \rangle$.

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Koopman mode decomposition

$$x_{n+1} = F(x_n)$$

$$[\mathcal{K}g](x) = g(F(x))$$

$$g(x) = \sum_{\text{eigenvalues } \lambda_j} c_{\lambda_j} \underbrace{\varphi_{\lambda_j}(x)}_{\text{eigenfunction of } \mathcal{K}} + \int_{-\pi}^{\pi} \underbrace{\phi_{\theta,g}(x)}_{\text{continuous spectrum}} d\theta$$

$$g(x_n) = [\mathcal{K}^n g](x_0) = \sum_{\text{eigenvalues } \lambda_j} c_{\lambda_j} \lambda_j^n \varphi_{\lambda_j}(x_0) + \int_{-\pi}^{\pi} e^{in\theta} \phi_{\theta,g}(x_0) d\theta$$

Encodes: geometric features, invariant measures, transient behavior, long-time behavior, coherent structures, quasiperiodicity, etc.

GOAL: Data-driven approximation of \mathcal{K} and its spectral properties.

Koopman mode decomposition

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eigenfunction of \mathcal{K}

continuous spectrum

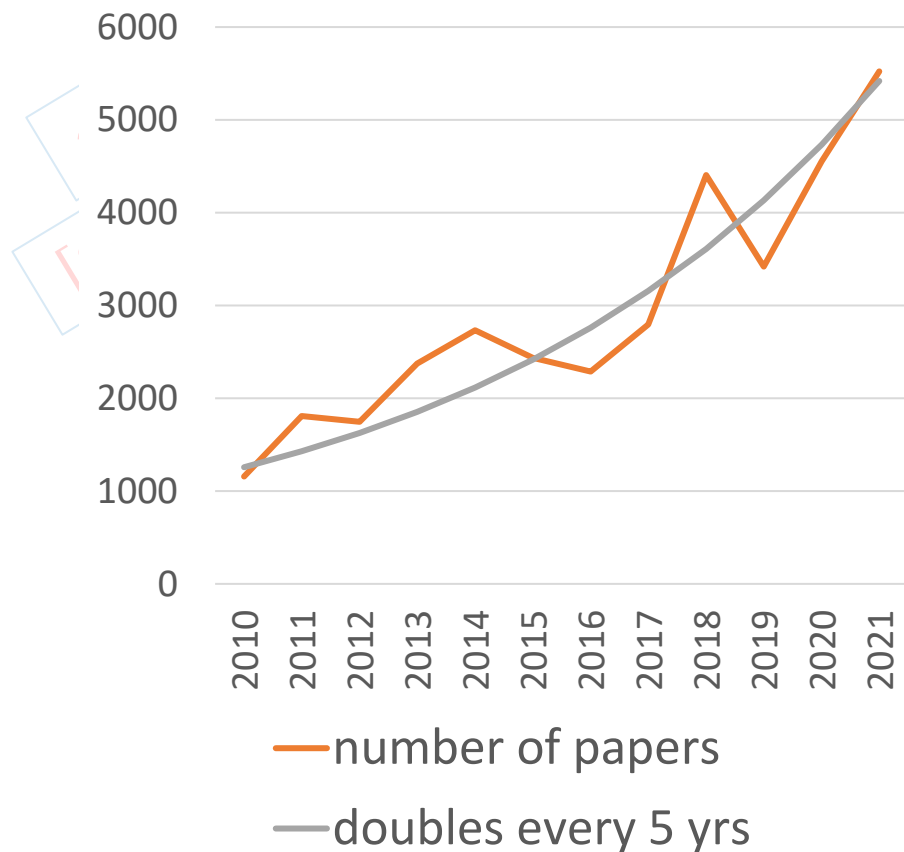
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Koopman mode decomposition

**New Papers on
“Koopman Operators”**



$$\sum_{\text{values } \lambda_j} c_{\lambda_j} \varphi_{\lambda_j}(x) + \int_{-\pi}^{\pi} \phi_{\theta,g}(x) d\theta$$

eigenfunction of \mathcal{K} (pointing to $\varphi_{\lambda_j}(x)$)

continuous spectrum (pointing to $\int_{-\pi}^{\pi} \phi_{\theta,g}(x) d\theta$)

$$\sum_{\text{values } \lambda_j} c_{\lambda_j} \lambda_j^n \varphi_{\lambda_j}(x_0) + \int_{-\pi}^{\pi} e^{in\theta} \phi_{\theta,g}(x_0) d\theta$$

s, invariant measures, transient behavior, it structures, quasiperiodicity, etc.

GOAL: Data-driven approximation of \mathcal{K} and its **spectral properties.**

Setting: Measure-preserving systems

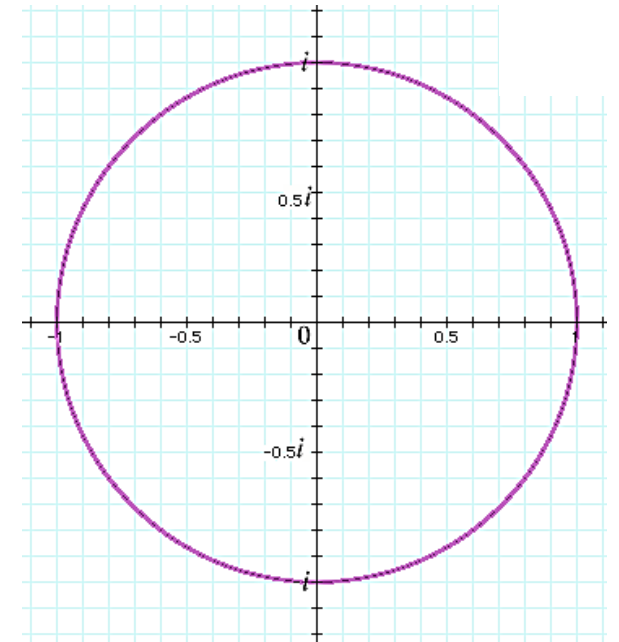
$$[\mathcal{K}g](x) = g(F(x)), \quad g \in L^2(\Omega, \omega)$$

F preserves $\omega \iff \|\mathcal{K}g\| = \|g\|$ (isometry)

$$\iff \mathcal{K}^* \mathcal{K} = I$$

$$\implies \text{Sp}(\mathcal{K}) \subseteq \{z: |z| \leq 1\}$$

(NB: unitary extensions of \mathcal{K} via Wold decomposition.)



Shift example (on $\ell^2(\mathbb{Z})$)

$$\begin{pmatrix} \ddots & & & & & & \\ & \ddots & & & & & \\ & & 0 & 1 & & & \\ & & & 0 & 1 & & \\ & & & & 0 & 1 & \\ & & & & & 0 & 1 \\ & & & & & & 0 & \ddots \\ & & & & & & & \ddots \end{pmatrix} \xrightarrow{\text{Two-way infinite}} \begin{pmatrix} 0 & 1 & & & & \\ & \ddots & \ddots & & & \\ & & \ddots & \ddots & & \\ & & & \ddots & & \\ & & & & 1 & \\ & & & & & 0 \end{pmatrix} \in \mathbb{C}^{N \times N}$$

- Spectrum is unit circle.
- Spectrum is stable.
- Continuous spectra.
- Unitary evolution.

- Spectrum is $\{0\}$.
- Spectrum is unstable.
- Discrete spectra.
- Nilpotent evolution.



Caution

Lots of Koopman operators are built up from operators like these!

How to fix a Jordan block

$$\begin{pmatrix} 0 & 1 & & \\ & \ddots & \ddots & \\ & & 0 & 1 \\ & & & 0 \end{pmatrix} \in \mathbb{C}^{N \times N}$$

polar
decomposition



Circulant matrix

$$\begin{pmatrix} 0 & 1 & & \\ & \ddots & \ddots & \\ & & 0 & 1 \\ 1 & & & 0 \end{pmatrix} \in \mathbb{C}^{N \times N}$$

- Spectrum is $\{0\}$.
- Spectrum is unstable.
- Nilpotent evolution.

- Spectrum converges to unit circle as $N \rightarrow \infty$.
- Spectrum is stable.
- Unitary evolution.

Extended Dynamic Mode Decomposition (EDMD)

$$\Psi(x) = [\psi_1(x) \quad \dots \quad \psi_N(x)], \quad g = \sum_{j=1}^N \mathbf{g}_j \psi_j = \Psi \mathbf{g} \in \text{span} \{\psi_1, \dots, \psi_N\}$$

$$\min_{\mathbb{K} \in \mathbb{C}^{N \times N}} \left\{ \int_{\Omega} \max_{\|\mathbf{g}\|_2=1} |\Psi(x) \mathbb{K} \mathbf{g} - [\mathcal{K}g](x)|^2 d\omega(x) = \int_{\Omega} \|\Psi(x) \mathbb{K} - \Psi(F(x))\|_2^2 d\omega(x) \right\}$$

quadrature

$$\{x^{(m)}, y^{(m)} = F(x^{(m)})\}_{m=1}^M$$

$$\min_{\mathbb{K} \in \mathbb{C}^{N \times N}} \sum_{m=1}^M w_m \|\Psi(x^{(m)}) \mathbb{K} - \Psi(y^{(m)})\|_2^2$$

\mathbb{K} : Galerkin method on $V_N = \text{span} \{\psi_1, \dots, \psi_N\}$

- Schmid, "Dynamic mode decomposition of numerical and experimental data," *J. Fluid Mech.*, 2010.
- Rowley, Mezić, Bagheri, Schlatter, Henningson, "Spectral analysis of nonlinear flows," *J. Fluid Mech.*, 2009.
- Kutz, Brunton, Brunton, Proctor, "Dynamic mode decomposition: data-driven modeling of complex systems," *SIAM*, 2016.
- Williams, Kevrekidis, Rowley "A data-driven approximation of the Koopman operator: Extending dynamic mode decomposition," *J. Nonlinear Sci.*, 2015.

A simple alteration

$$G_{jk} = \sum_{m=1}^M w_m \overline{\psi_j(x^{(m)})} \psi_k(x^{(m)}) \approx \langle \psi_k, \psi_j \rangle$$

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Measure-preserving: $\mathbf{g}^* G \mathbf{g} \approx \|g\|^2 = \|\mathcal{K}g\|^2 \approx \mathbf{g}^* \mathbb{K}^* G \mathbb{K} \mathbf{g}$

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Enforce: $G = \mathbb{K}^* G \mathbb{K}$

A simple alteration

$$G_{jk} = \sum_{m=1}^M w_m \overline{\psi_j(x^{(m)})} \psi_k(x^{(m)}) \approx \langle \psi_k, \psi_j \rangle$$

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Enforce: $G = \mathbb{K}^* G \mathbb{K}$

quadrature

**Orthogonal
Procrustes problem**

$$\min_{\substack{\mathbb{K} \in \mathbb{C}^{N \times N} \\ G = \mathbb{K}^* G \mathbb{K}}} \sum_{m=1}^M w_m \|\Psi(x^{(m)}) \mathbb{K} G^{-1/2} - \Psi(y^{(m)}) G^{-1/2}\|_2^2$$

The mpEDMD algorithm

Algorithm 4.1 The mpEDMD algorithm

Input: Snapshot data $\mathbf{X} \in \mathbb{C}^{d \times M}$ and $\mathbf{Y} \in \mathbb{C}^{d \times M}$, quadrature weights $\{w_m\}_{m=1}^M$, and a dictionary of functions $\{\psi_j\}_{j=1}^N$.

- 1: Compute the matrices Ψ_X and Ψ_Y and $\mathbf{W} = \text{diag}(w_1, \dots, w_M)$.
- 2: Compute an economy QR decomposition $\mathbf{W}^{1/2} \Psi_X = \mathbf{Q}\mathbf{R}$, where $\mathbf{Q} \in \mathbb{C}^{M \times N}$, $\mathbf{R} \in \mathbb{C}^{N \times N}$.
- 3: Compute an SVD of $(\mathbf{R}^{-1})^* \Psi_Y^* \mathbf{W}^{1/2} \mathbf{Q} = \mathbf{U}_1 \Sigma \mathbf{U}_2^*$.
- 4: Compute the eigendecomposition $\mathbf{U}_2 \mathbf{U}_1^* = \hat{\mathbf{V}} \Lambda \hat{\mathbf{V}}^*$ (via a Schur decomposition).
- 5: Compute $\mathbb{K} = \mathbf{R}^{-1} \mathbf{U}_2 \mathbf{U}_1^* \mathbf{R}$ and $\mathbf{V} = \mathbf{R}^{-1} \hat{\mathbf{V}}$.

Output: Koopman matrix \mathbb{K} with eigenvectors \mathbf{V} and eigenvalues Λ .

$V_N = \text{span} \{\psi_1, \dots, \psi_N\}$
 $\mathcal{P}_{V_N}: L^2(\Omega, \omega) \rightarrow V_N$
 orthogonal projection

As $M \rightarrow \infty$, **unitary part** of polar decomposition of $\mathcal{P}_{V_N} \mathcal{K} \mathcal{P}_{V_N}^*$.

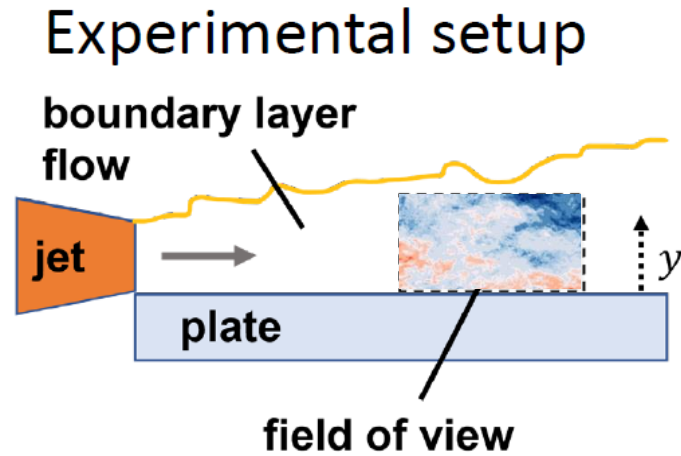
- C., "The mpEDMD Algorithm for Data-Driven Computations of Measure-Preserving Dynamical Systems," **SINUM**, 2023.
- Code: <https://github.com/MColbrook/Measure-preserving-Extended-Dynamic-Mode-Decomposition>

Convergence properties (theorems in paper)

- Spectral measures.
- Functional calculus, L^2 forecasting etc.
- Koopman mode decomposition.
- Spectrum.
- Resolvent (see later!)

Key ingredient: **unitary** discretization.

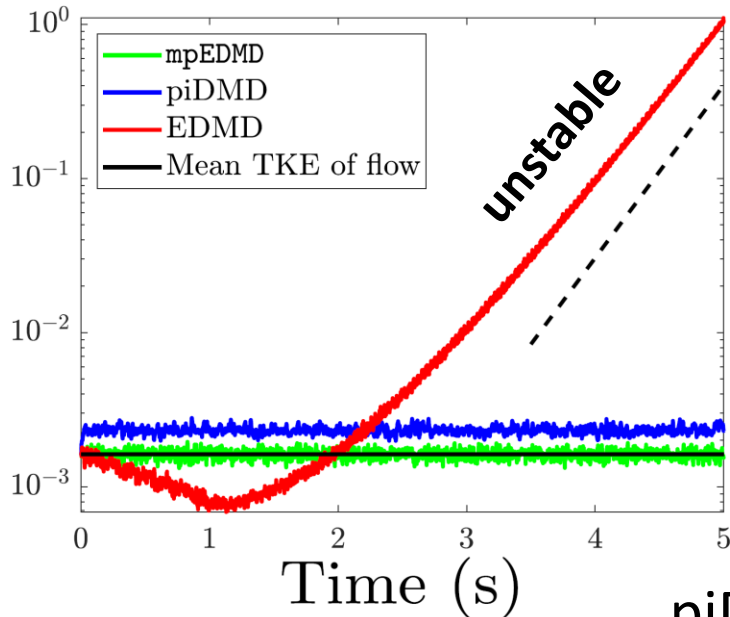
Turbulence (real data)



- Reynolds number $\approx 6.4 \times 10^4$
- Ambient dimension (d) $\approx 100,000$ (velocity at measurement points)

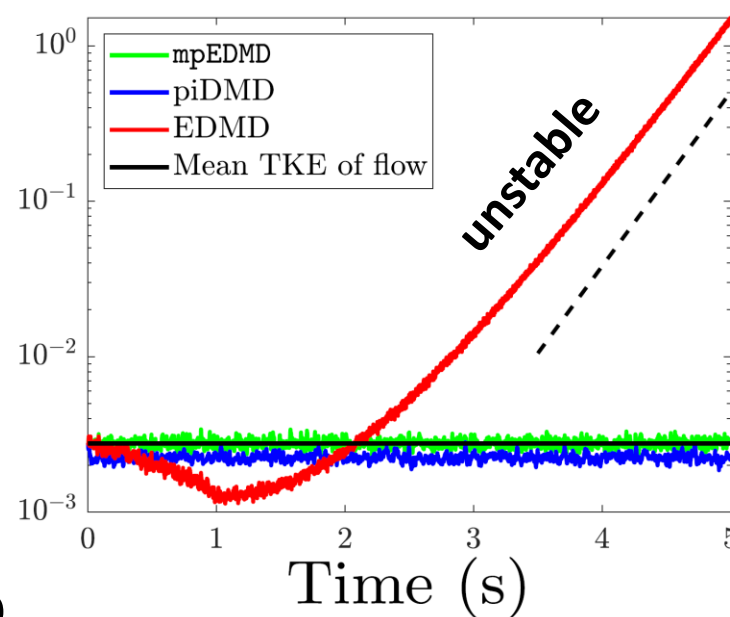
*PIV data provided by Máté Szőke (Virginia Tech)

Turbulent K.E. $y=5\text{mm}$



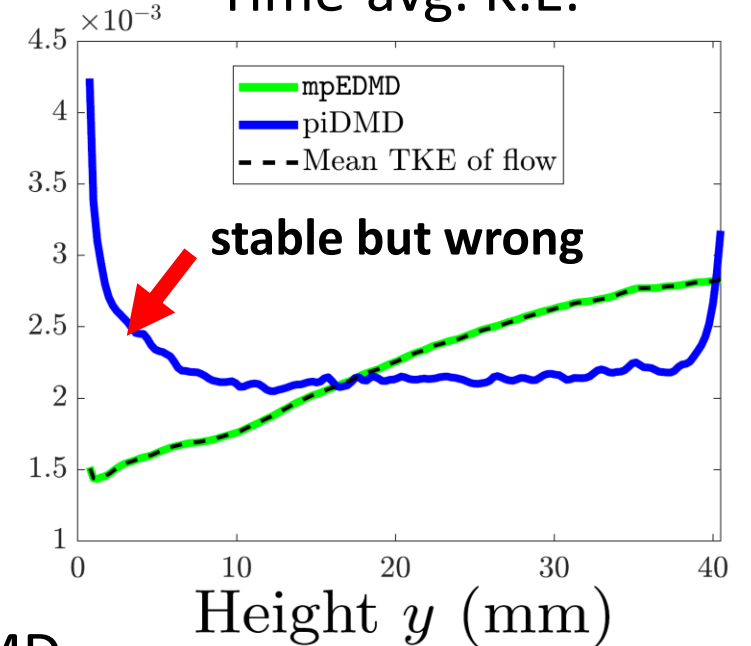
piDMD

Turbulent K.E. $y=35\text{mm}$



EDMD

Time-avg. K.E.



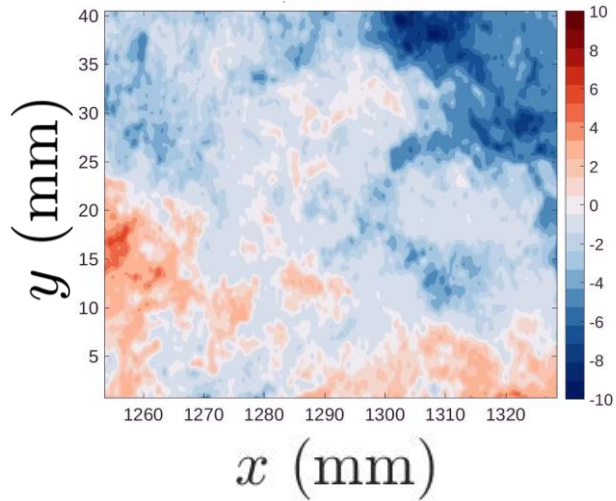
• Baddoo, Herrmann, McKeon, Kutz, Brunton, "Physics-informed dynamic mode decomposition (piDMD)," preprint.

• Williams, Kevrekidis, Rowley "A data-driven approximation of the Koopman operator: Extending dynamic mode decomposition," *J. Nonlinear Sci.*, 2015.

Turbulence statistics

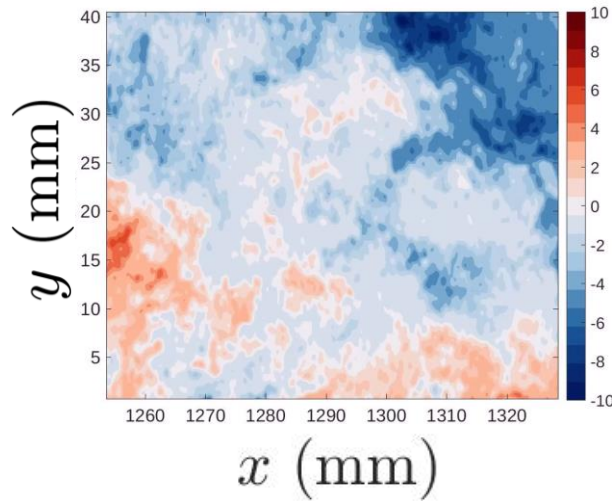
Flow

time=0.001000



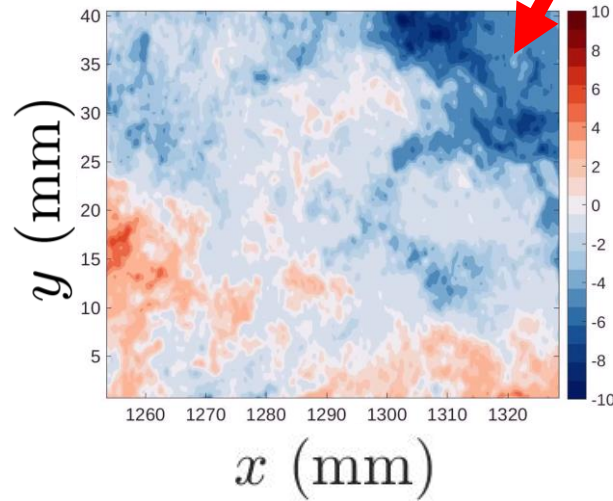
mpEDMD

time=0.001000



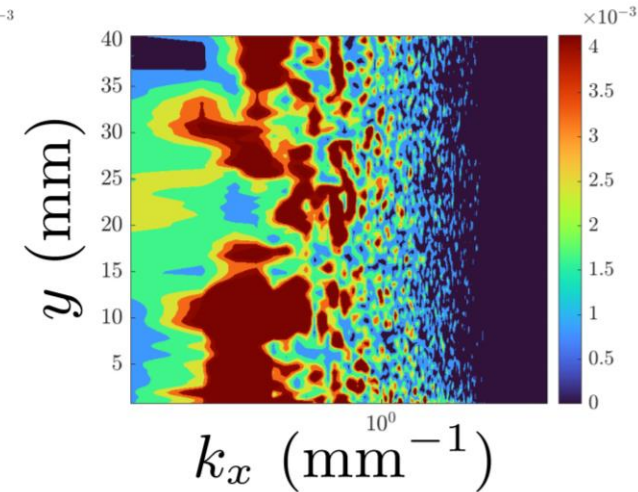
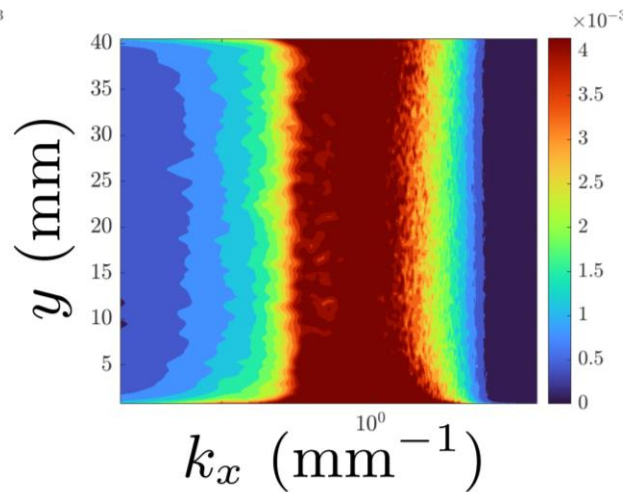
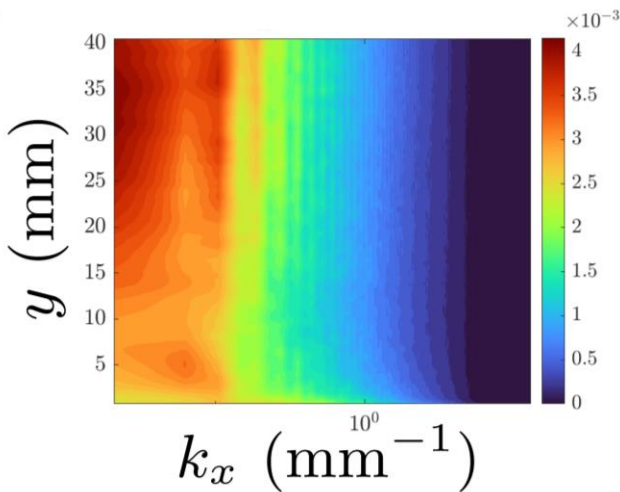
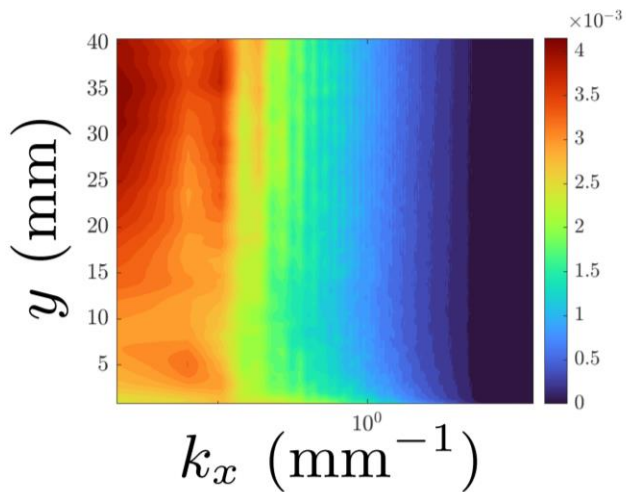
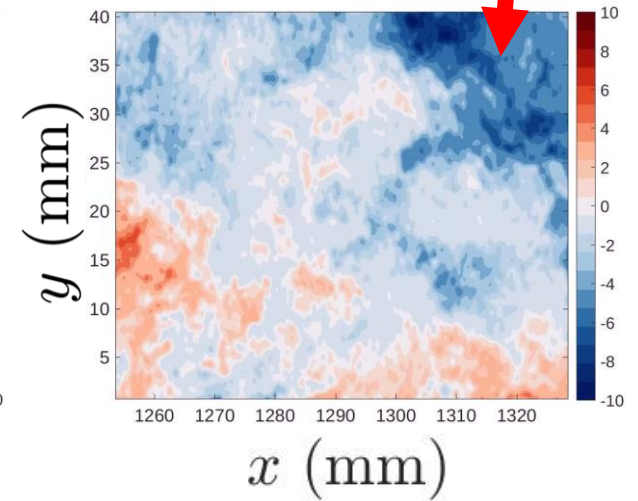
piDMD

time=0.001000



EDMD

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Back to the shift!

“Solve” $(U - zI)u_z = 0$

$$u_z = \sum_{j=-\infty}^{\infty} z^j e_j$$

$e_j \rightarrow e_{j-1}$

$$U = \begin{pmatrix} \ddots & & & & & & \\ & \ddots & & & & & \\ & & 0 & 1 & & & \\ & & & 0 & 1 & & \\ & & & & 0 & 1 & \\ & & & & & 0 & 1 & \ddots \\ & & & & & & & 0 & 1 & \ddots \end{pmatrix}$$

Two-way infinite

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Doesn't live in $\ell^2(\mathbb{Z})!!!$

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Let $|z| = 1$, $\phi = \sum_{j=-\infty}^{\infty} \phi_j e_j$ where ϕ_j decay faster than any inverse polynomial.

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Test functions

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Test functions

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Generalised eigenfunctions u_z and generalised eigenvalues $\{z: |z| = 1\}$

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$e_j \rightarrow e_{j-1}$

Two-way infinite

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Generalised eigenfunctions u_z and generalised eigenvalues $\{z: |z| = 1\}$

RIGGED HILBERT SPACE

Example: Nonlinear pendulum


$$\dot{x}_1 = x_2, \quad \dot{x}_2 = -\sin(x_1)$$

$$\Omega = [-\pi, \pi]_{\text{per}} \times \mathbb{R},$$

$$\Delta_t = 1,$$

$\omega =$ Lebesgue measure

**Considered a challenge in
Koopman theory!**



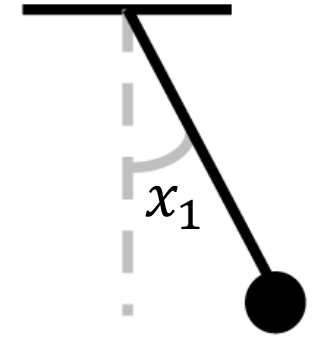
ARTICLE

DOI: [10.1038/s41467-018-07210-0](https://doi.org/10.1038/s41467-018-07210-0) OPEN

Deep learning for universal linear embeddings of nonlinear dynamics

Bethany Lusch^{1,2}, J. Nathan Kutz¹ & Steven L. Brunton^{1,2}

Identifying coordinate transformations that make strongly nonlinear dynamics approximately linear has the potential to enable nonlinear prediction, estimation, and control using linear theory. The Koopman operator is a leading data-driven embedding, and its eigenfunctions provide intrinsic coordinates that globally linearize the dynamics. However, identifying and representing these eigenfunctions has proven challenging. This work leverages deep learning to discover representations of Koopman eigenfunctions from data. Our network is parsimonious and interpretable by construction, embedding the dynamics on a low-dimensional manifold. We identify nonlinear coordinates on which the dynamics are globally linear using a



Explicit diagonalization using Radon transform!

- Action-angle coordinates (n degrees of freedom):

$$\dot{\mathbf{I}} = \mathbf{0}, \quad \dot{\boldsymbol{\theta}} = \mathbf{I}, \quad \Omega = \mathbb{R}^n \times [-\pi, \pi]_{\text{per}}^n$$

Explicit diagonalization using Radon transform!

- Action-angle coordinates (n degrees of freedom):

$$\dot{\mathbf{I}} = \mathbf{0}, \quad \dot{\boldsymbol{\theta}} = \mathbf{I}, \quad \Omega = \mathbb{R}^n \times [-\pi, \pi]_{\text{per}}^n$$

- g in Schwartz space,

$$g(\mathbf{I}, \boldsymbol{\theta}) = \sum_{\mathbf{k} \in \mathbb{Z}^n} \hat{g}_{\mathbf{k}}(\mathbf{I}) e^{i\mathbf{k} \cdot \boldsymbol{\theta}}, \quad \hat{g}_{\mathbf{k}}(\mathbf{I}) = \frac{1}{(2\pi)^n} \int_{[-\pi, \pi]_{\text{per}}^n} g(\mathbf{I}, \boldsymbol{\theta}) e^{-i\mathbf{k} \cdot \boldsymbol{\theta}} d\boldsymbol{\theta}$$

Explicit diagonalization using Radon transform!

- Action-angle coordinates (n degrees of freedom):

$$\dot{\mathbf{I}} = \mathbf{0}, \quad \dot{\boldsymbol{\theta}} = \mathbf{I}, \quad \Omega = \mathbb{R}^n \times [-\pi, \pi]_{\text{per}}^n$$

- g in Schwartz space,

$$g(\mathbf{I}, \boldsymbol{\theta}) = \sum_{\mathbf{k} \in \mathbb{Z}^n} \hat{g}_{\mathbf{k}}(\mathbf{I}) e^{i\mathbf{k} \cdot \boldsymbol{\theta}}, \quad \hat{g}_{\mathbf{k}}(\mathbf{I}) = \frac{1}{(2\pi)^n} \int_{[-\pi, \pi]_{\text{per}}^n} g(\mathbf{I}, \boldsymbol{\theta}) e^{-i\mathbf{k} \cdot \boldsymbol{\theta}} d\boldsymbol{\theta}$$

$$\hat{g}_{\mathbf{k}}(\mathbf{I}) = \sum_{j=1}^{\infty} \sum_{m=-\infty}^{\infty} \int_{[-\pi, \pi]_{\text{per}}} \left\langle g_{\theta}^{(\mathbf{k}, m, j)*} \mid g \right\rangle g_{\theta}^{(\mathbf{k}, m, j)} d\theta$$

$$g_{\theta}^{(\mathbf{k}, m, j)} = \frac{1}{(2\pi)^n} \underbrace{\delta(\theta + 2\pi m - \Delta t \mathbf{k} \cdot \mathbf{I})}_{\text{Supported on hyperplane}} \underbrace{\psi_j^{(k)} e^{i\mathbf{k} \cdot \boldsymbol{\theta}}}_{\text{Orthonormal basis of hyperplane}}$$

Plane wave

Generalised eigenfunctions

Supported on
hyperplane

Orthonormal basis of
hyperplane

Gelfand's theorem → diagonalisation

- **Finite matrix:** $B \in \mathbb{C}^{n \times n}$, $B^*B = BB^*$, orthonormal basis of e-vectors $\{v_j\}_{j=1}^n$

$$v = \sum_{j=1}^n (v_j^* v) v_j, \quad Bv = \sum_{j=1}^n \lambda_j (v_j^* v) v_j \quad \forall v \in \mathbb{C}^n$$

- **Infinite dimensions:** Unitary \mathcal{K} . Typically, **no basis of eigenfunctions!**
Some technical assumptions (can always be realized):

$$g = \int_{[-\pi, \pi]_{\text{per}}} \underbrace{\langle g_\theta^* | g \rangle}_{\text{Koopman modes}} g_\theta dv(\theta), \quad \mathcal{K}g = \int_{[-\pi, \pi]_{\text{per}}} e^{i\theta} \langle g_\theta^* | g \rangle g_\theta dv(\theta)$$

$g \in S \subset L^2(\Omega, \omega)$

generalized eigenfunctions
distributions $\in S^*$

$e^{i\theta} = \lambda$

Koopman Mode Decomposition

Rigged DMD: Smoothing

Carathéodory function:

$$F_g(z) = (\mathcal{K} + zI)(\mathcal{K} - zI)^{-1}g = \int_{[-\pi, \pi]_{\text{per}}} \boxed{\frac{e^{i\theta} + z}{e^{i\theta} - z}} \langle g_\theta^* | g \rangle g_\theta d\nu(\theta)$$

Rigged DMD: Smoothing

Carathéodory function:

$$F_g(z) = (\mathcal{K} + zI)(\mathcal{K} - zI)^{-1}g = \int_{[-\pi, \pi]_{\text{per}}} \frac{e^{i\theta} + z}{e^{i\theta} - z} \langle g_{\theta}^* | g \rangle g_{\theta} d\nu(\theta)$$

Let $r = 1 + \varepsilon > 1$, $\theta_0 \in [-\pi, \pi]_{\text{per}}$,

$$\frac{1}{4\pi} [F_g(r^{-1}e^{i\theta_0}) - F_g(re^{i\theta_0})]$$

Rigged DMD: Smoothing

Carathéodory function:

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Let $r = 1 + \varepsilon > 1$, $\theta_0 \in [-\pi, \pi]_{\text{per}}$,

$$\begin{aligned} & \frac{1}{4\pi} [F_g(r^{-1}e^{i\theta_0}) - F_g(re^{i\theta_0})] \\ &= \frac{1}{2\pi} \int_{[-\pi, \pi]_{\text{per}}} \frac{r^2 - 1}{1 + r^2 - 2r\cos(\theta_0 - \theta)} \langle g_{\theta}^* | g \rangle g_{\theta} d\nu(\theta) \end{aligned}$$

Rigged DMD: Smoothing

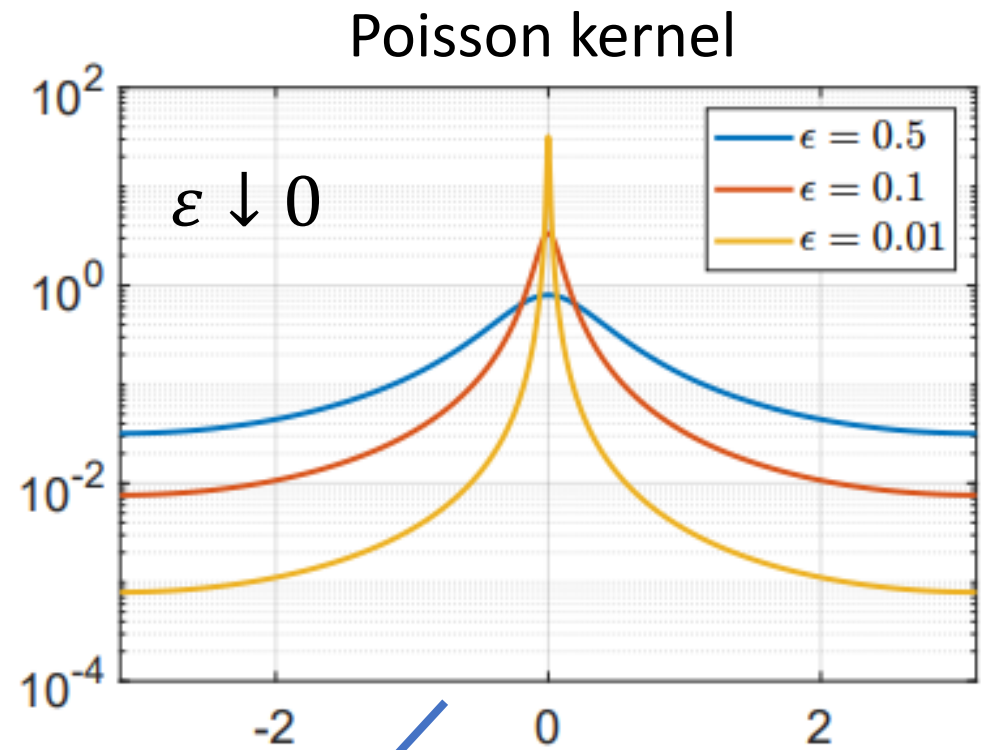
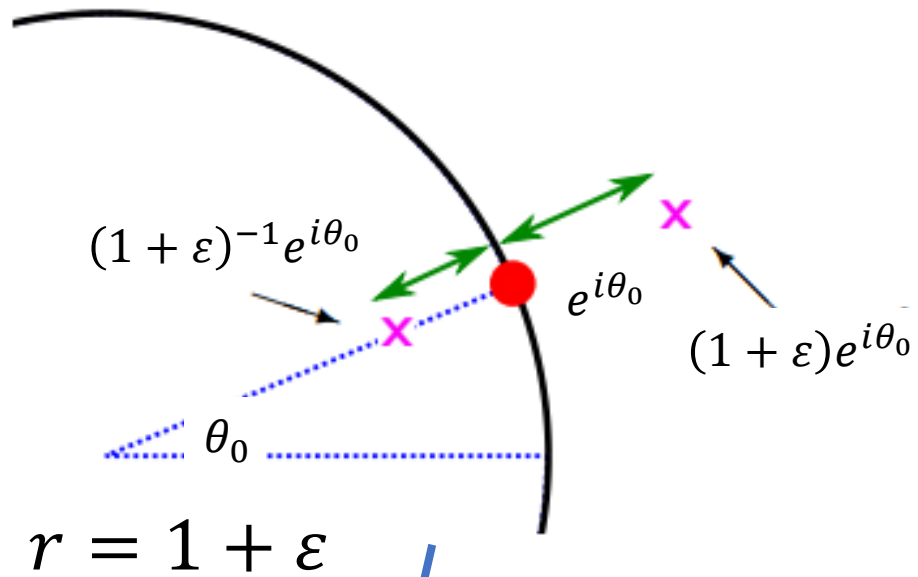
Carathéodory function:

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Smoothed generalized eigenfunction



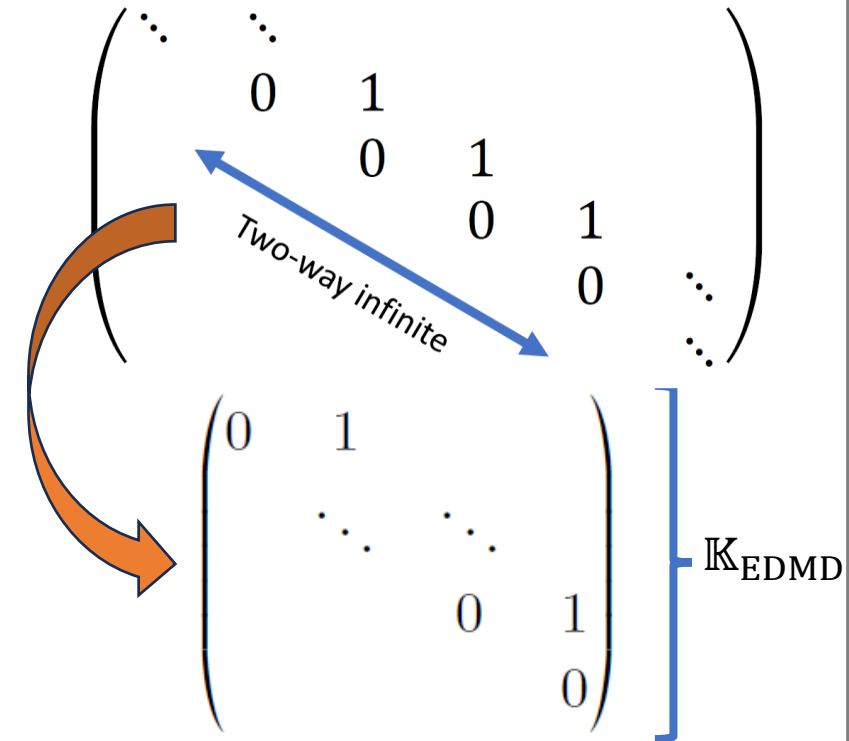
$$\frac{1}{4\pi} [F_g(r^{-1}e^{i\theta_0}) - F_g(re^{i\theta_0})]$$

$$= \frac{1}{2\pi} \int_{[-\pi, \pi]_{\text{per}}} \underbrace{\frac{r^2 - 1}{1 + r^2 - 2r\cos(\theta_0 - \theta)}}_{\text{Poisson kernel}} \langle g_\theta^* | g \rangle g_\theta dv(\theta)$$

Smoothed generalized eigenfunction

F_g requires $(\mathcal{K} - zI)^{-1}$

EDMD diverges:



Acts on $\text{span}\{e_{-N}, \dots, e_N\}$

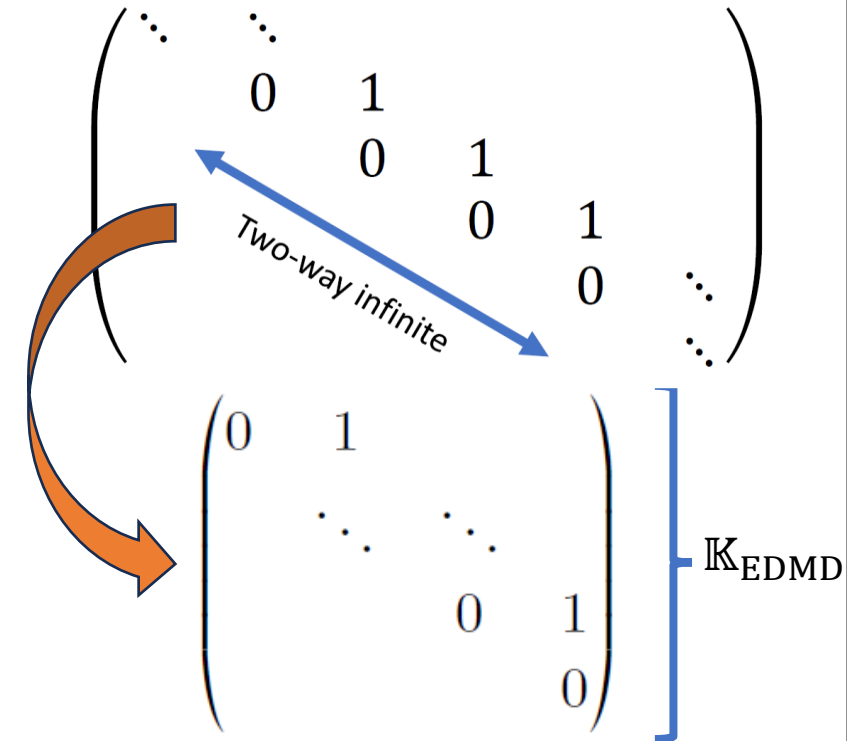
Exponential
blowup
as $N \rightarrow \infty$.

E.g., if $|z| < 1$,

$$(\mathbb{K}_{\text{EDMD}} - zI)^{-1}e_0 = \sum_{j=1}^N \frac{-1}{z^j} e_{-j}$$

F_g requires $(\mathcal{K} - zI)^{-1}$

EDMD diverges:



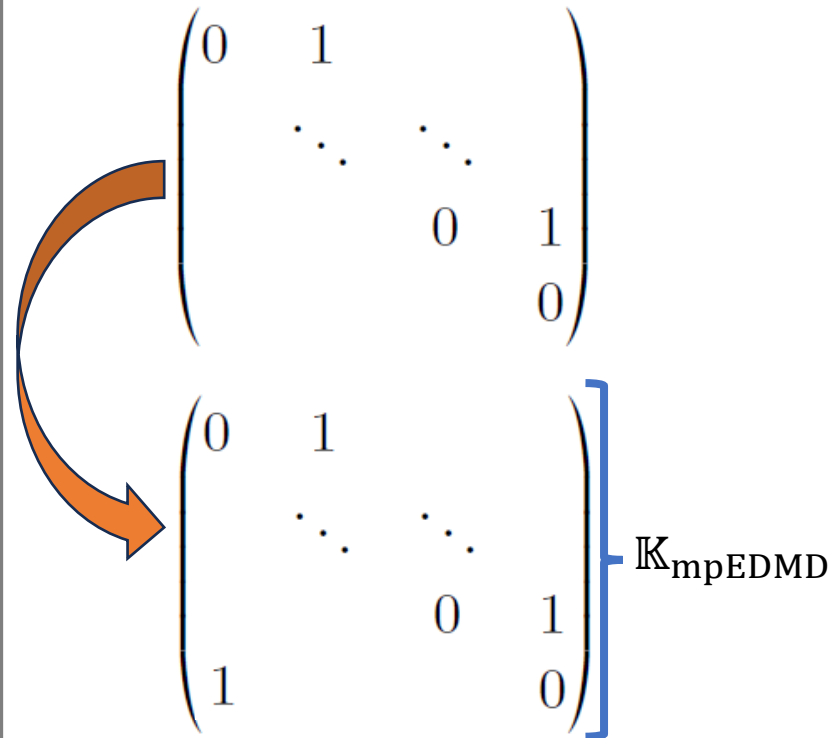
Acts on $\text{span}\{e_{-N}, \dots, e_N\}$

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Exponential
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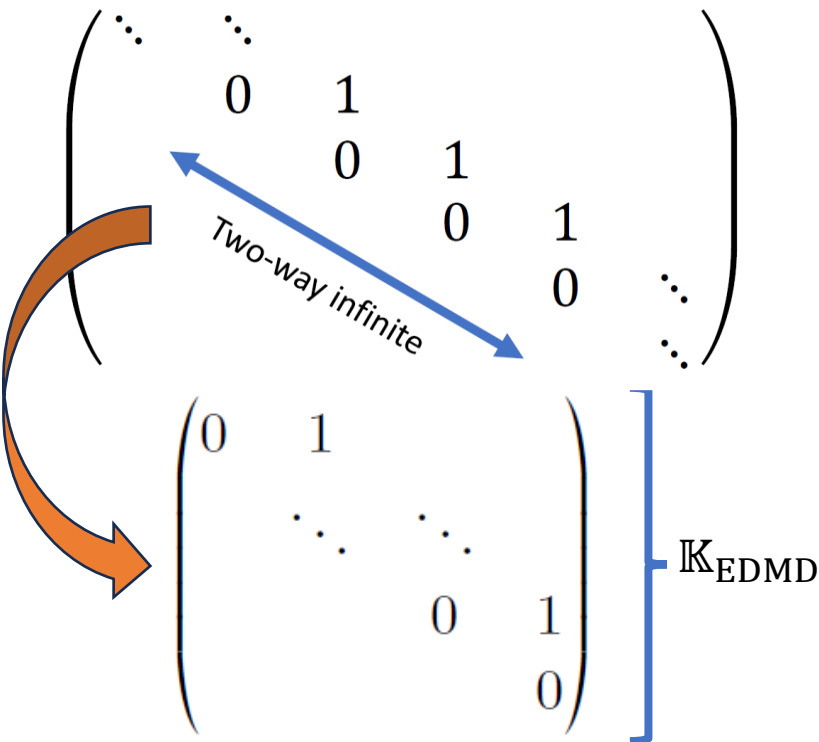
mpEDMD converges:



General method: unitary part of a
polar decomposition of EDMD!

F_g requires $(\mathcal{K} - zI)^{-1}$

EDMD diverges:



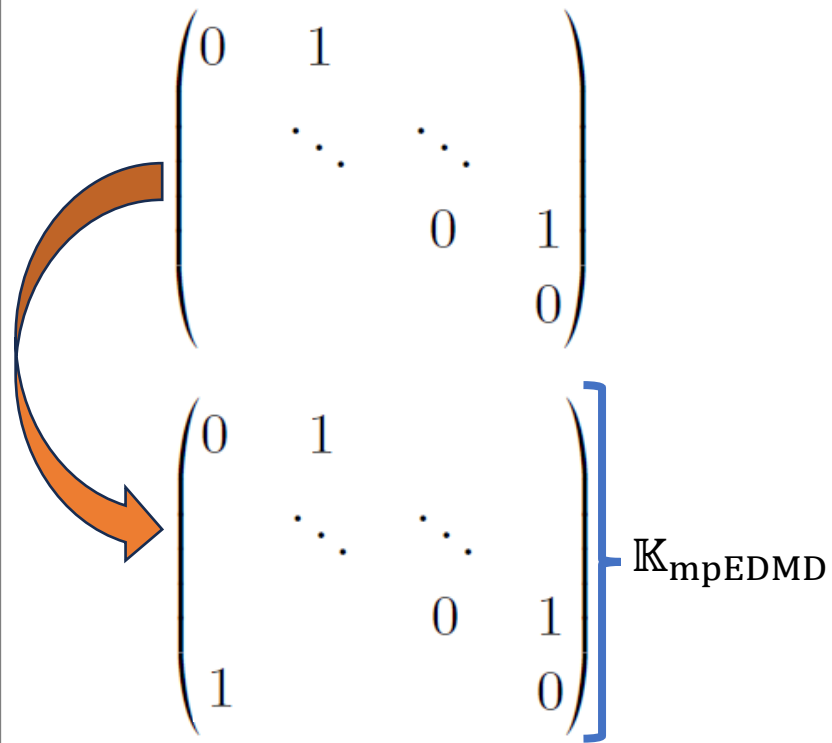
Acts on $\text{span}\{e_{-N}, \dots, e_N\}$

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mpEDMD converges:



General method: unitary part of a **polar decomposition** of EDMD!

Rigged DMD converges:

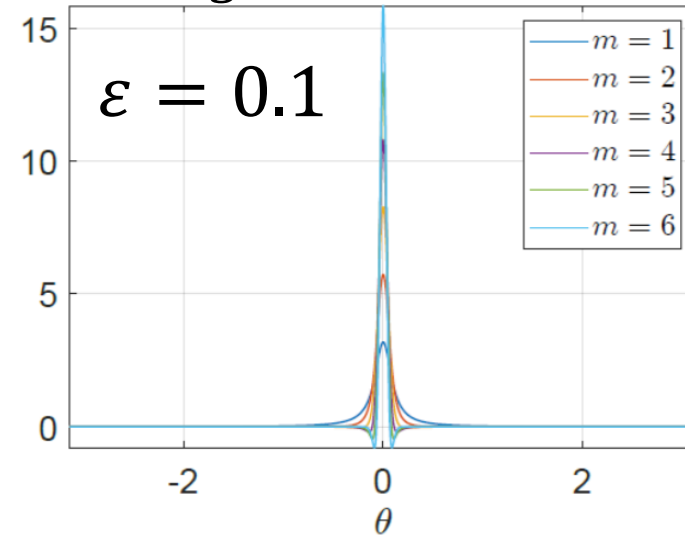
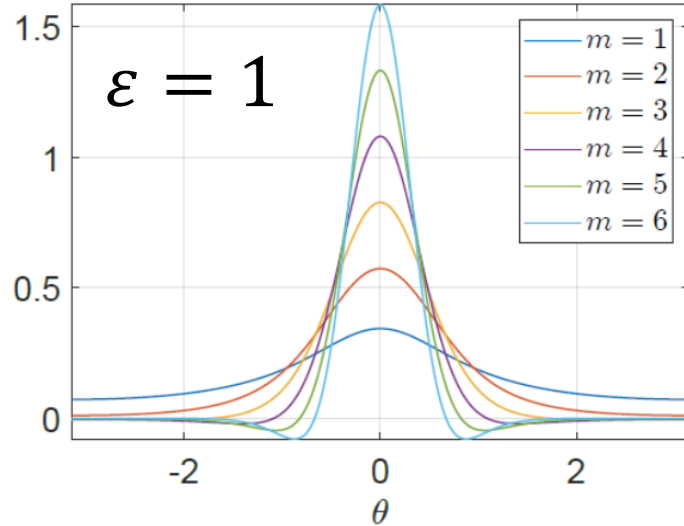
- For general \mathcal{K} : $(\mathbb{K}_{\text{mpEDMD}} - zI)^{-1} \mathbf{g}$ converges to $(\mathcal{K} - zI)^{-1} g$ as $\lim_{N \rightarrow \infty} \lim_{M \rightarrow \infty}$
- Hence, Rigged DMD converges as $\lim_{\varepsilon \downarrow 0} \lim_{N \rightarrow \infty} \lim_{M \rightarrow \infty}$
- ResDMD allows us to select $\varepsilon = \varepsilon(N)$ adaptively (convergence in **2 limits**)



Better smoothing kernels as $\varepsilon \downarrow 0$

- Poisson kernel: **slow** convergence $\mathcal{O}(\varepsilon \log(1/\varepsilon))$.
- Construct high-order *rational* kernels using $F_g(z)$.

High-Order
Kernels

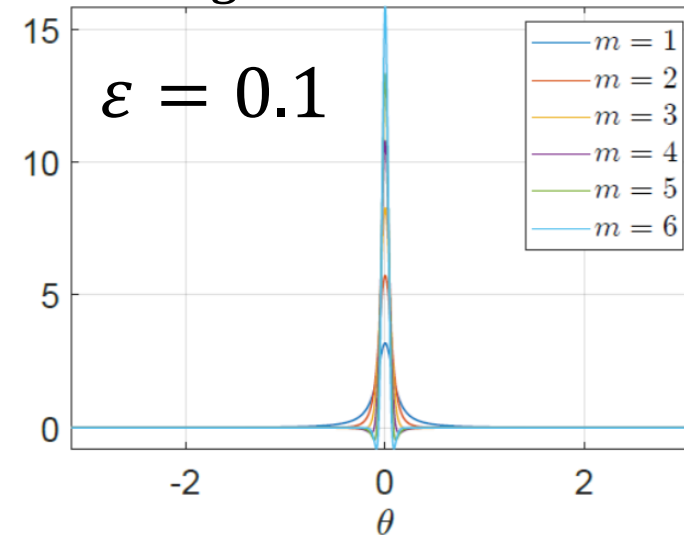
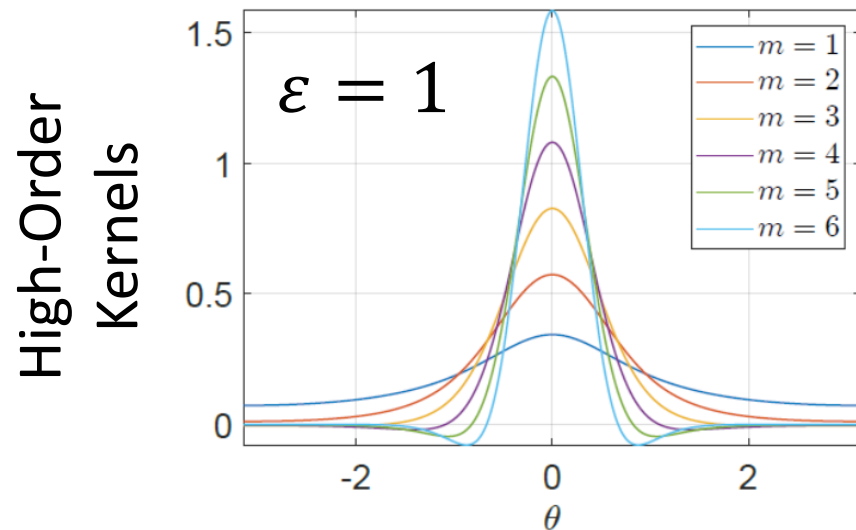


Smaller ε
requires
more data

Better smoothing kernels as $\varepsilon \downarrow 0$

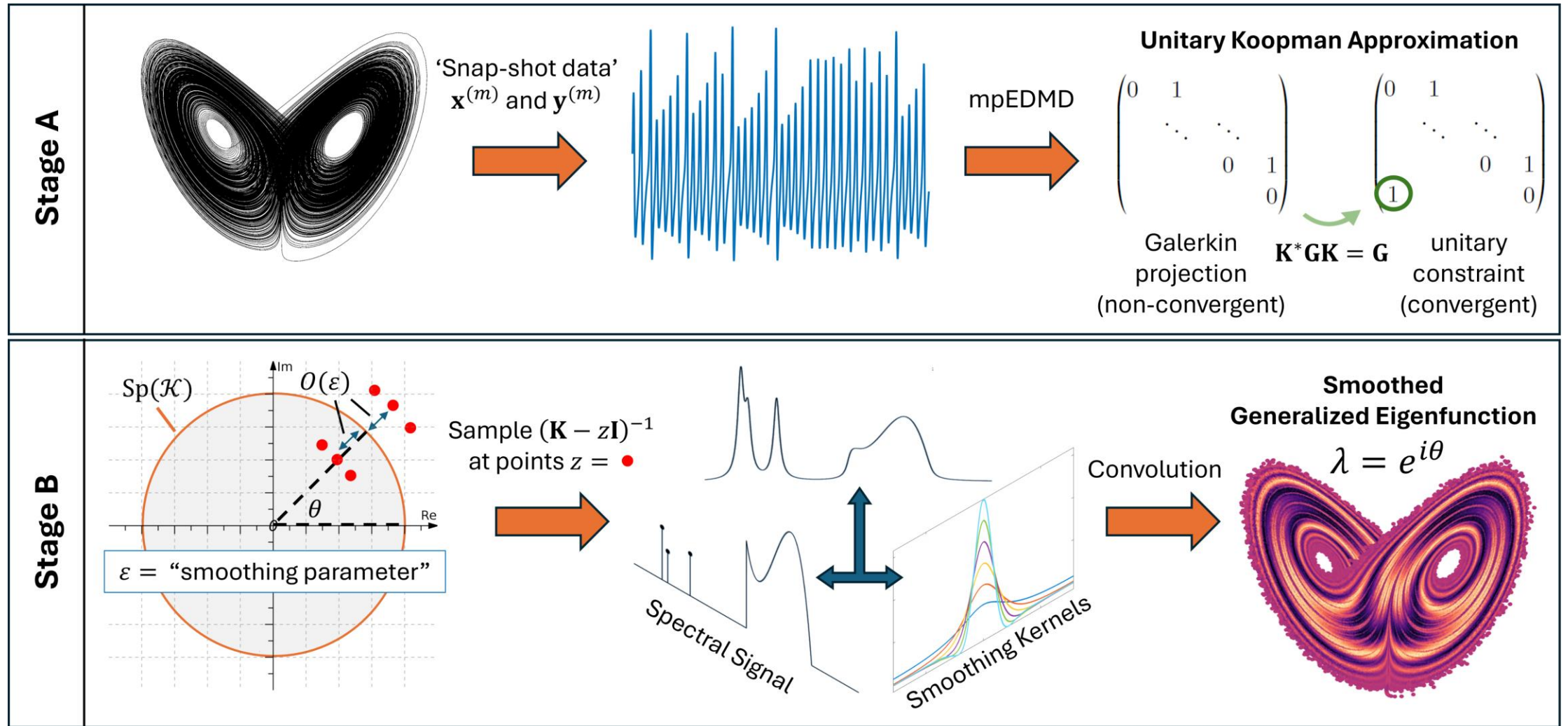
- Poisson kernel: **slow** convergence $\mathcal{O}(\varepsilon \log(1/\varepsilon))$.
- Construct high-order *rational* kernels using $F_g(z)$.

Smaller ε
requires
more data



- Theorem:** Suppose quadrature rule converges & $\lim_{N \rightarrow \infty} \text{dist}(h, V_N) = 0$ for any $h \in L^2(\Omega, \omega)$. Choosing $N = N(\varepsilon)$, **fast** $\mathcal{O}(\varepsilon^m \log(1/\varepsilon))$ convergence for:
- Generalized eigenfunctions (topology of \mathcal{S}^*).
 - Spectral measures (gen. efun. projections): pointwise, L^p , weak,...
 - Forecasting (i.e., iterating Koopman mode decomposition), coherency etc.

Rigged DMD

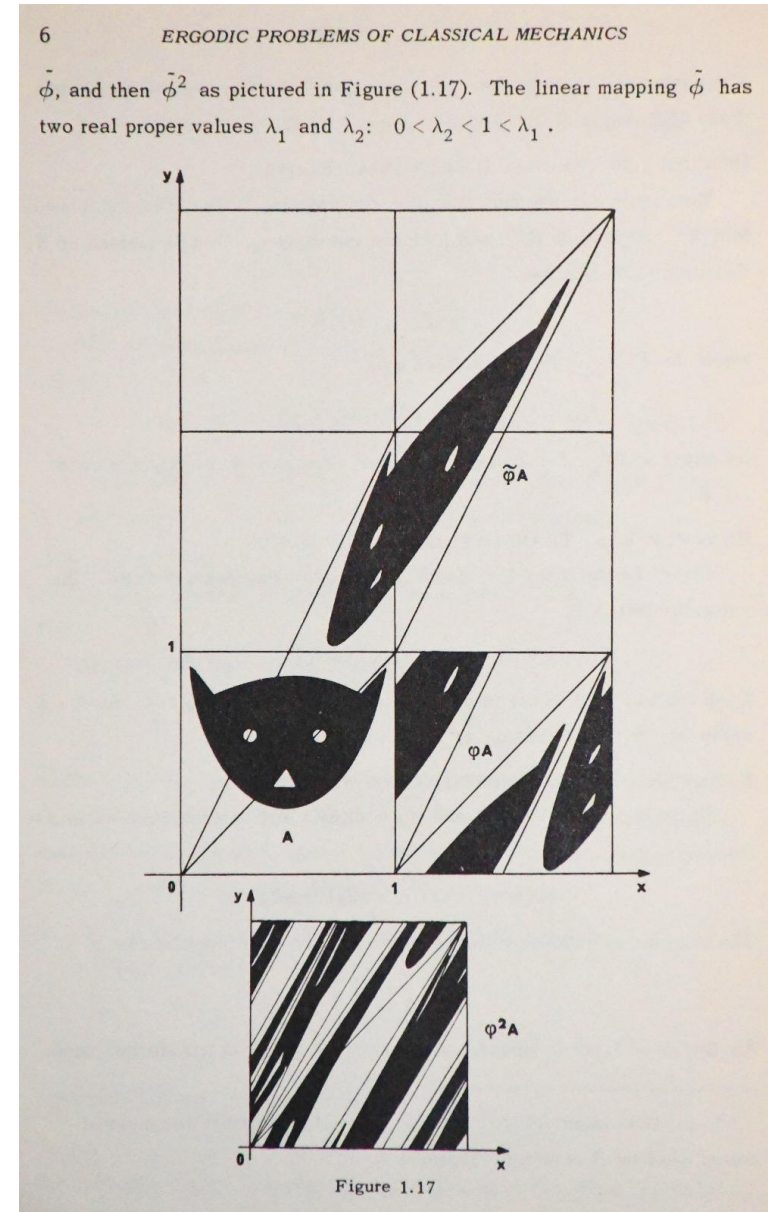


- C., Drysdale, Horning, "Rigged Dynamic Mode Decomposition: Data-Driven Generalized Eigenfunction Decompositions for Koopman Operators", arxiv preprint.
- Code: <https://github.com/MColbrook/Rigged-Dynamic-Mode-Decomposition>

Example: Arnold's cat map

$$F(x, y) = (2x + y, x + y) \bmod 2\pi$$

$$\Omega = [-\pi, \pi]_{\text{per}}^2, \quad \omega = \text{Lebesgue measure}$$



Arnold's "Ergodic Problems of Classical Mechanics"

Example: Arnold's cat map

Experimental details

Length-one trajectories, $M = 50 \times 50, N = 500$

$$g(x, y) = \sin(x) + \frac{1}{2} \sin(2x + y) + \frac{i}{4} \sin(5x + 3y)$$

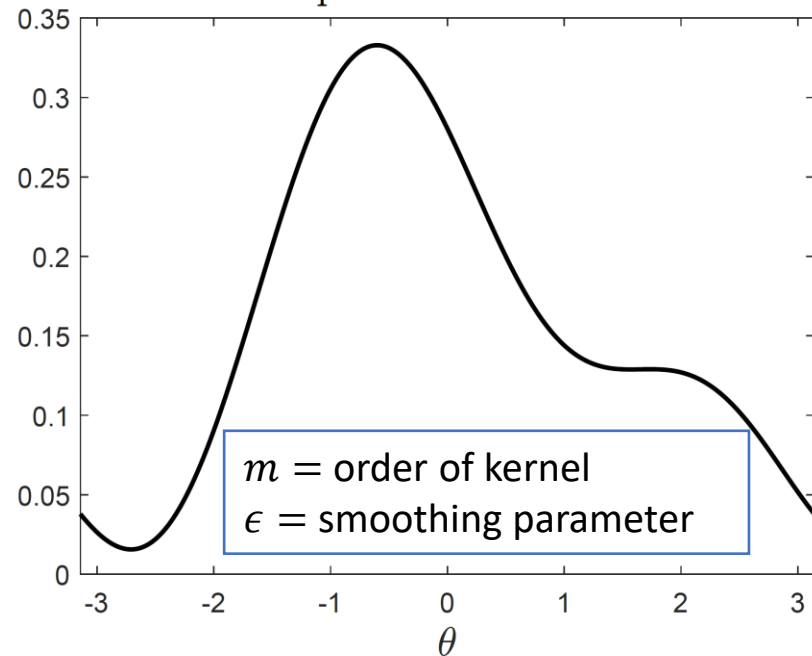
Krylov subspace: $V_N = \{g, \mathcal{K}g, \dots, \mathcal{K}^{N-1}g\}$

$$F(x, y) = (2x + y, x + y) \bmod 2\pi$$

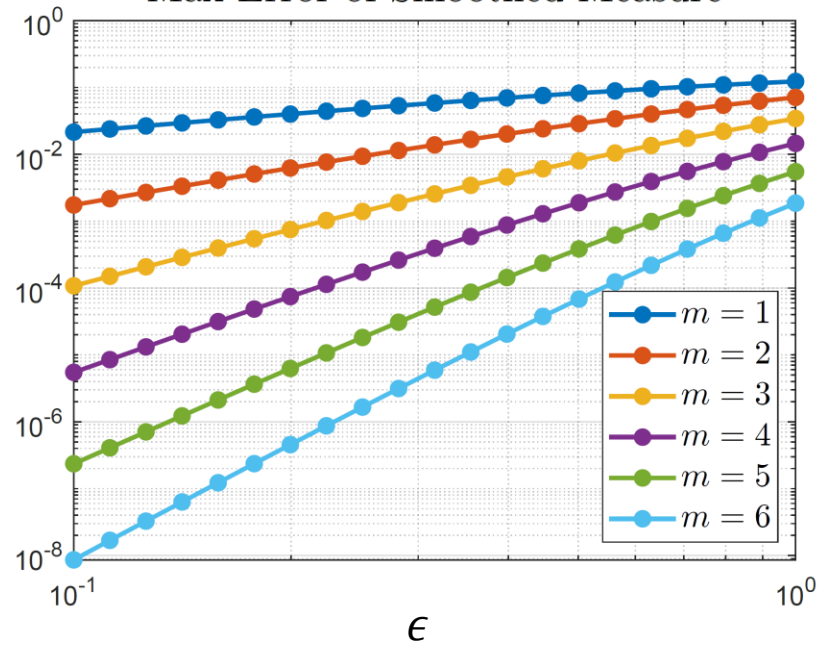
$$\Omega = [-\pi, \pi]_{\text{per}}^2, \quad \omega = \text{Lebesgue measure}$$

Explicit formula: g_θ become more oscillatory as $\epsilon \downarrow 0$ (non-decaying Fourier series)

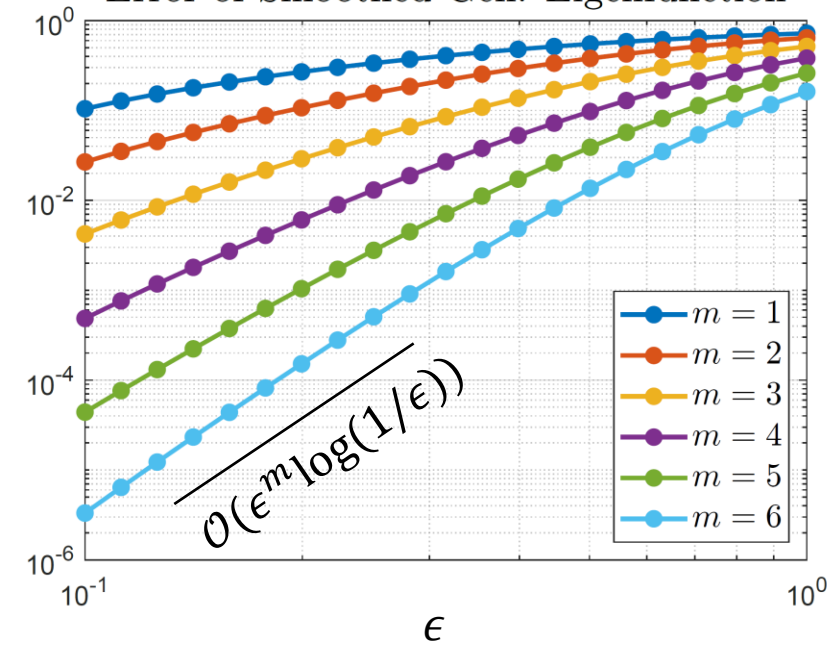
Spectral Measure



Max Error of Smoothed Measure



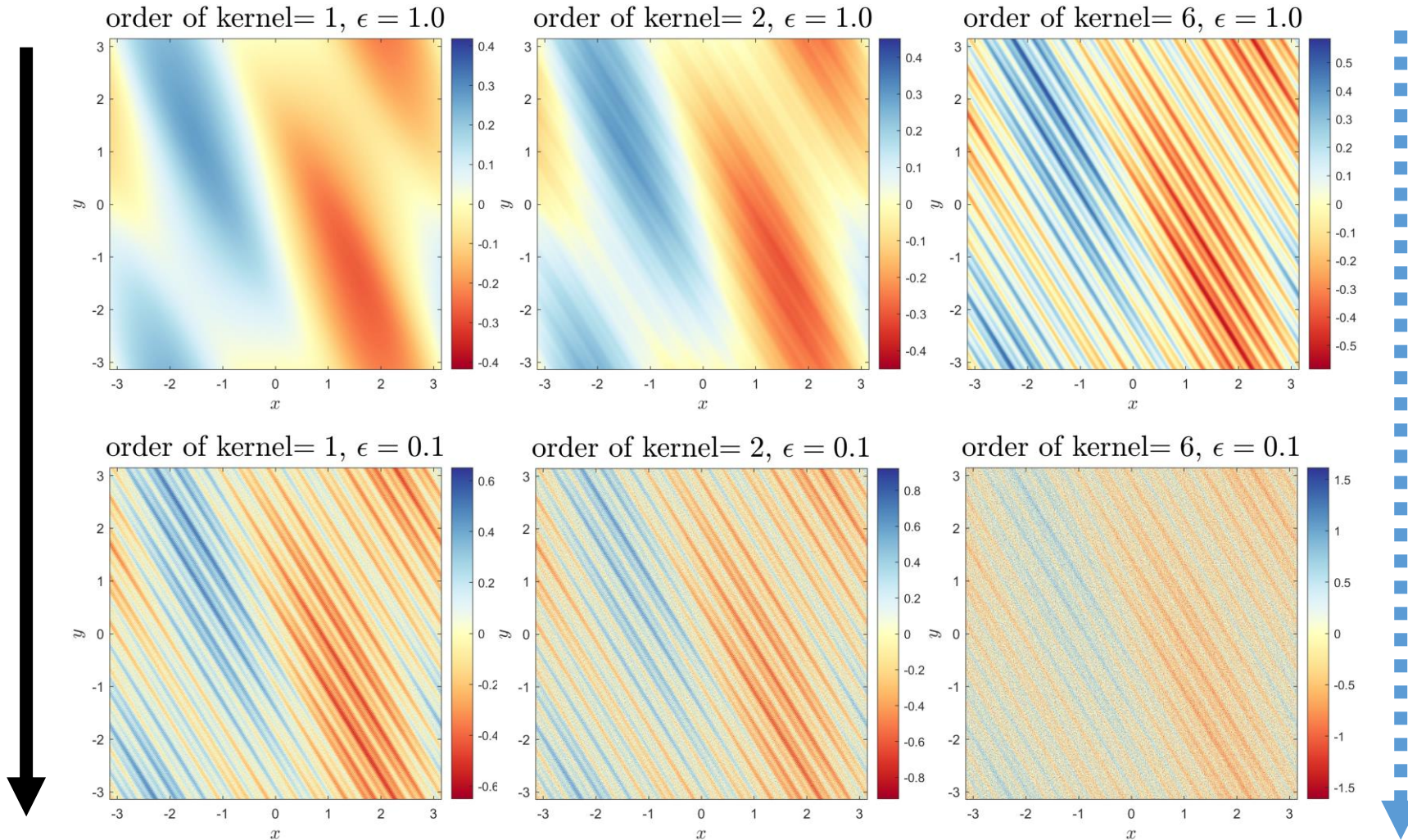
Error of Smoothed Gen. Eigenfunction



Higher kernel order (accuracy)



Higher resolution ($\epsilon \downarrow 0$)

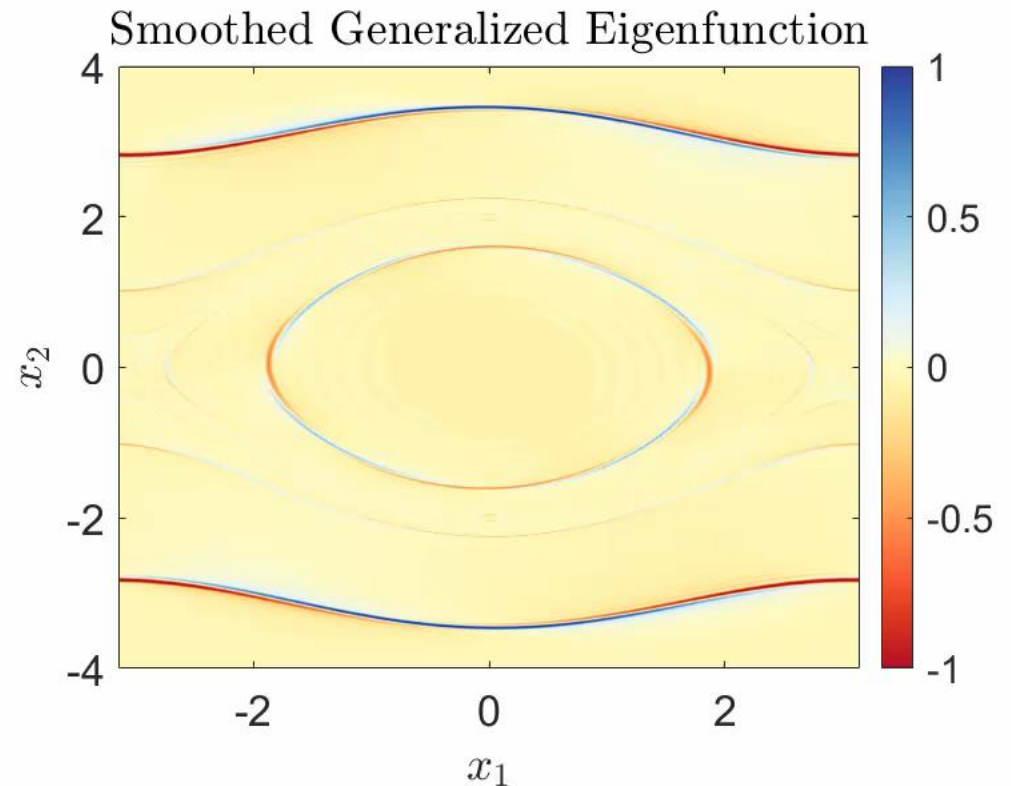
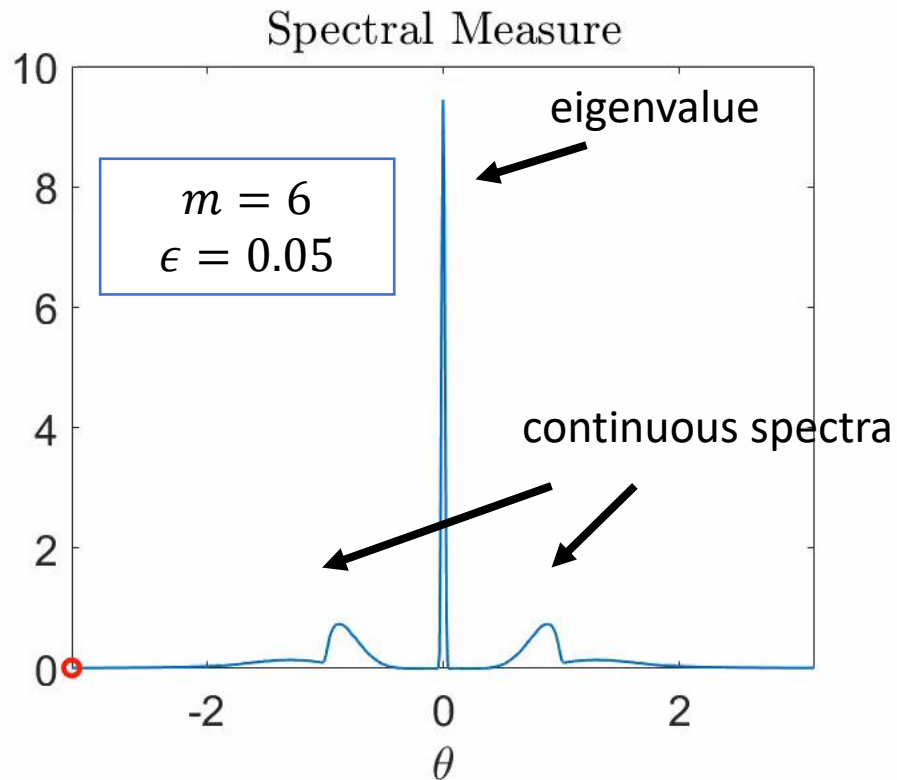


Example: Nonlinear pendulum

Experimental Details
 Length-one trajectories over grid
 $M = 500 \times 500, N = 300$
 $g(x_1, x_2) = \exp(ix_1) / \cosh(x_2)$
 Krylov subspace: $V_N = \{g, \mathcal{K}g, \dots, \mathcal{K}^{N-1}g\}$

$$\dot{x}_1 = x_2, \quad \dot{x}_2 = -\sin(x_1), \quad \Omega = [-\pi, \pi]_{\text{per}} \times \mathbb{R}, \quad \Delta_t = 1, \quad \omega = \text{Lebesgue measure}$$

Explicit formula: g_θ become plane waves concentrated on unions of lines of constant energy as $\epsilon \downarrow 0$.



\mathcal{S} constructed from Krylov subspace

- If \mathcal{K} is represented by an infinite matrix with finitely many non-zero entries in each column, can build \mathcal{S} using weighted sequence spaces.
- Always possible using time-delay embedding:
 - $\{ \text{Unions (different } g) \text{ of spaces } \text{span}\{g, \mathcal{K}g, \dots, \mathcal{K}^{N-1}g, \dots\} \} \subset \mathcal{S}$
- Generalises shift example: in coefficient space w.r.t. Krylov subspaces.

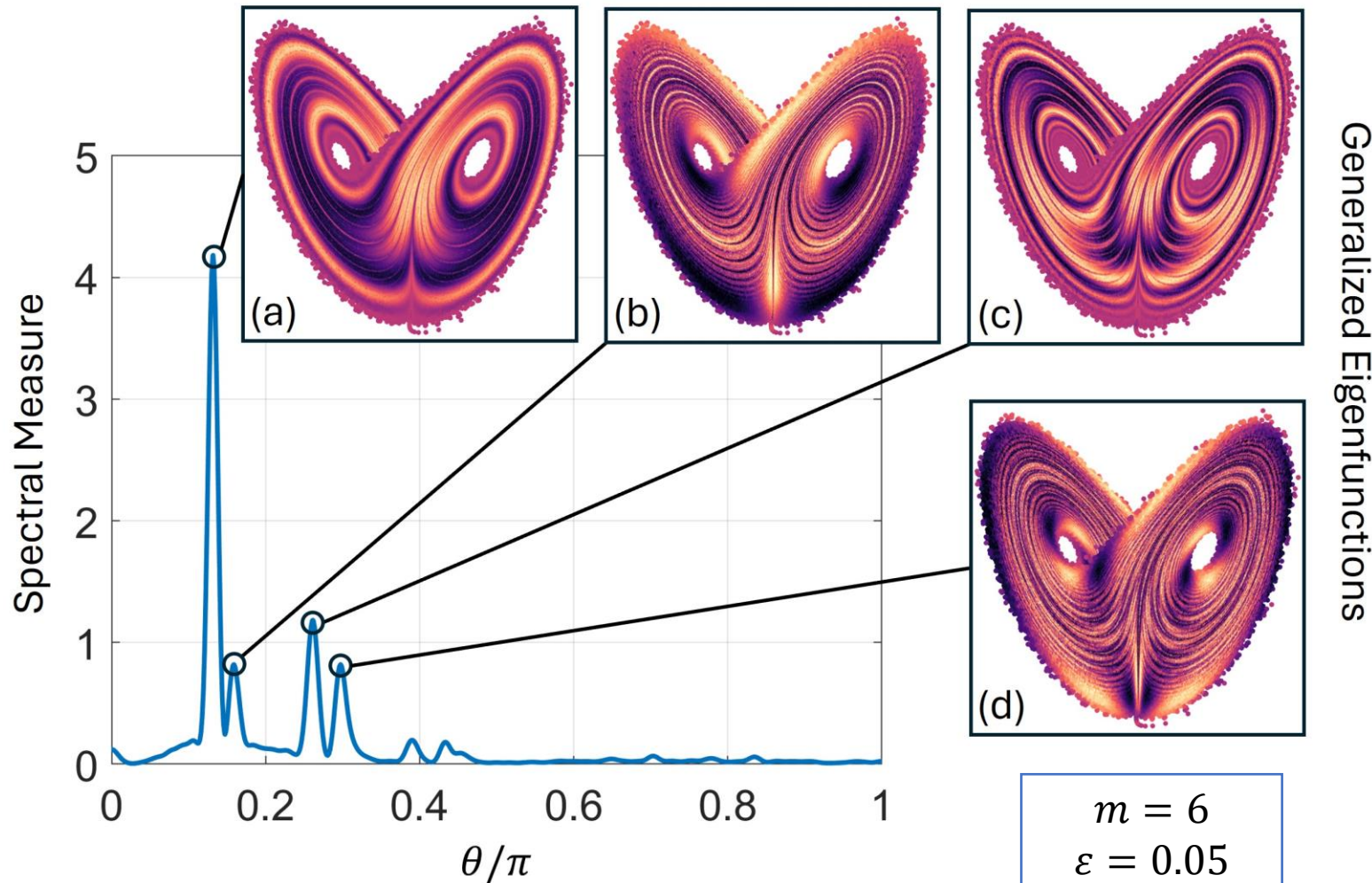
\mathcal{S} constructed from Krylov subspace

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- Generalises shift example: in coefficient space w.r.t. Krylov subspaces.

Let's do this for Lorenz...

Example: Lorenz system

$$\dot{x}_1 = 10(x_2 - x_1), \quad \dot{x}_2 = x_1(28 - x_3) - x_2, \quad \dot{x}_3 = x_1x_2 - 8/3 x_3, \quad \Delta_t = 0.05, \quad \Omega = \text{attractor}, \quad \omega = \text{SRB measure}$$



Generalized Eigenfunctions

No formula for
generalized eigenfunctions!!

Experimental Details

Single trajectory (ergodic system)

$$M = 10000, N = 1000$$

$$g(x_1, x_2, x_3) = \tanh\left(\frac{x_1x_2 - 5x_3}{10}\right) - c$$

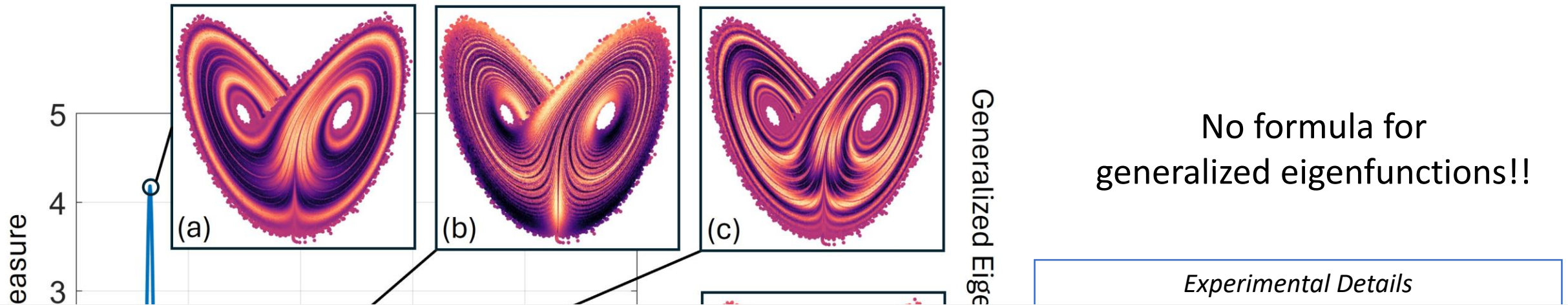
Krylov subspace: $V_N = \{g, \mathcal{K}g, \dots, \mathcal{K}^{N-1}g\}$

$$m = 6$$

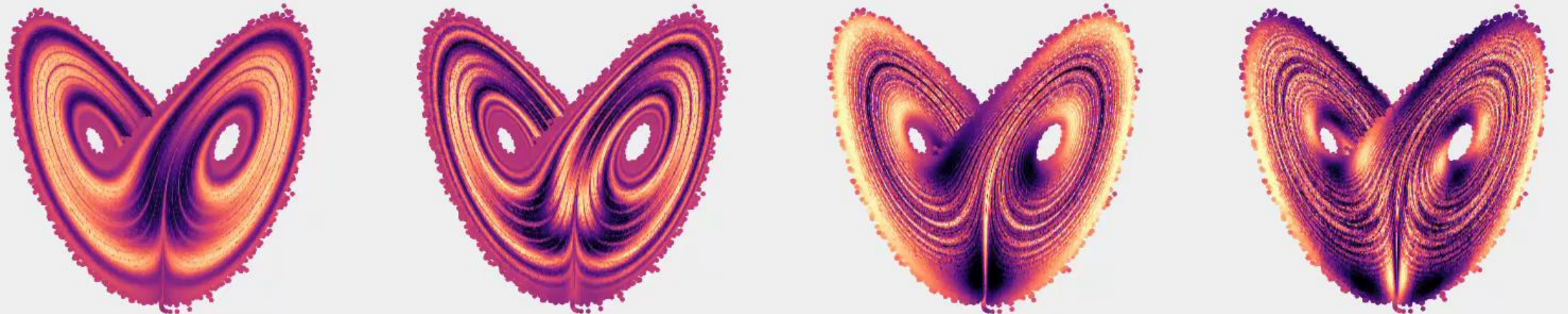
$$\varepsilon = 0.05$$

Example: Lorenz system

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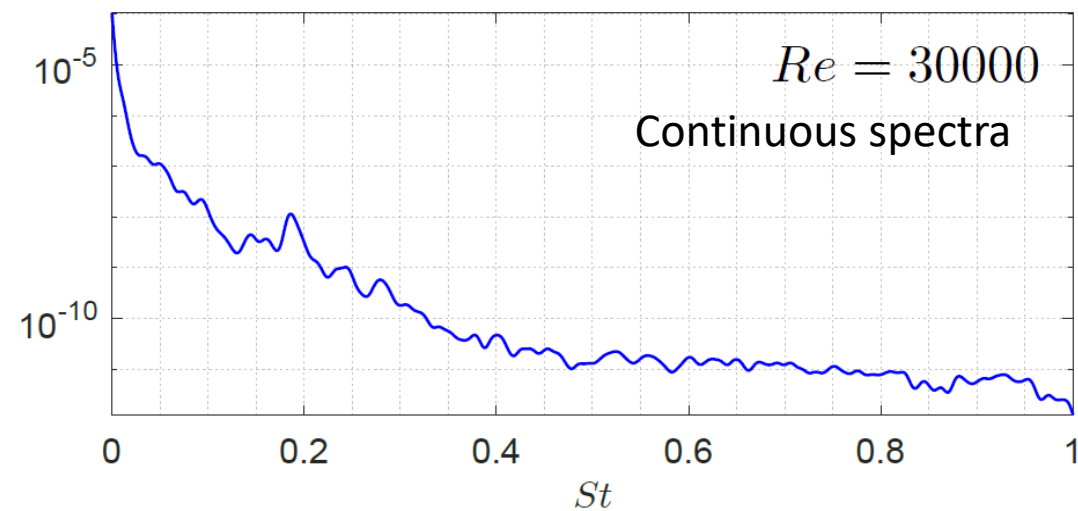
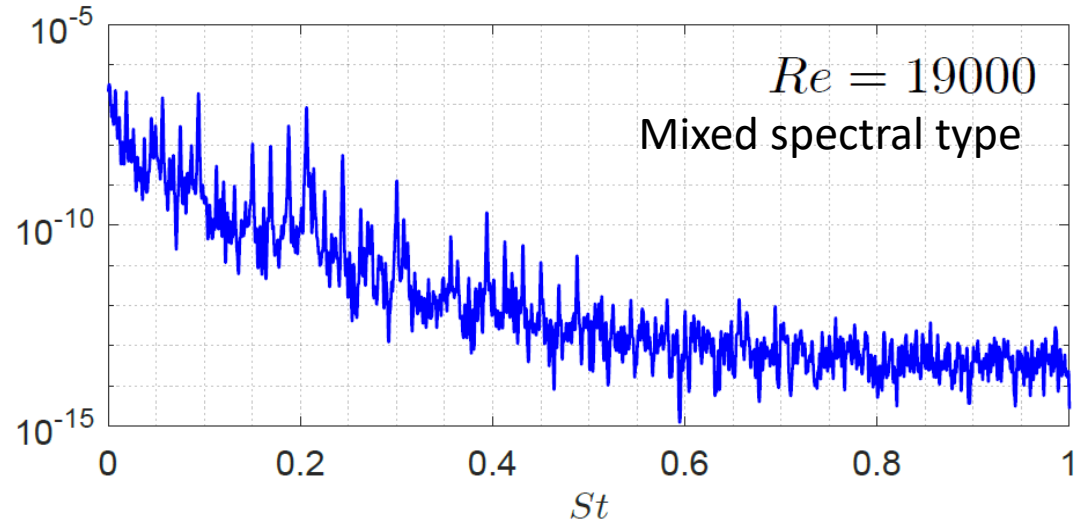
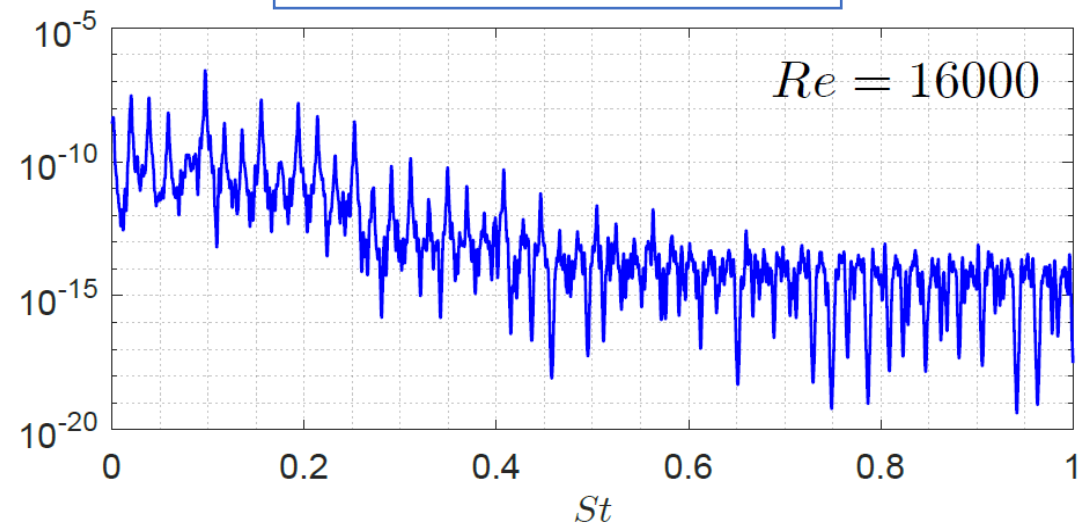
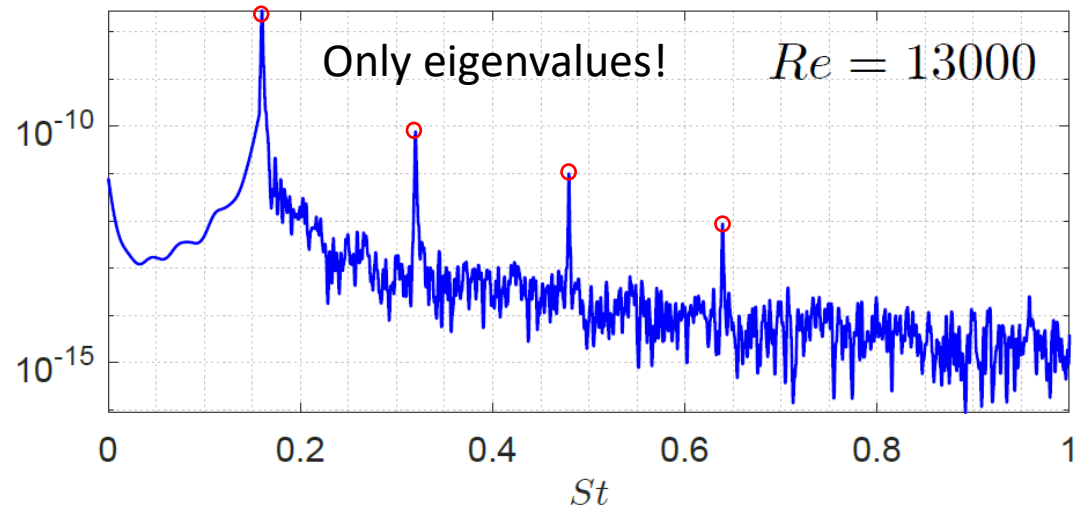


No formula for
generalized eigenfunctions!!



Example: Noisy cavity flow (spectral measures)

Single trajectory
 $M = 10000, N$ varies
Basis: POD modes
20% Gaussian noise
*Raw measurements provided
Arbabi and Mezić (PRF 2017)

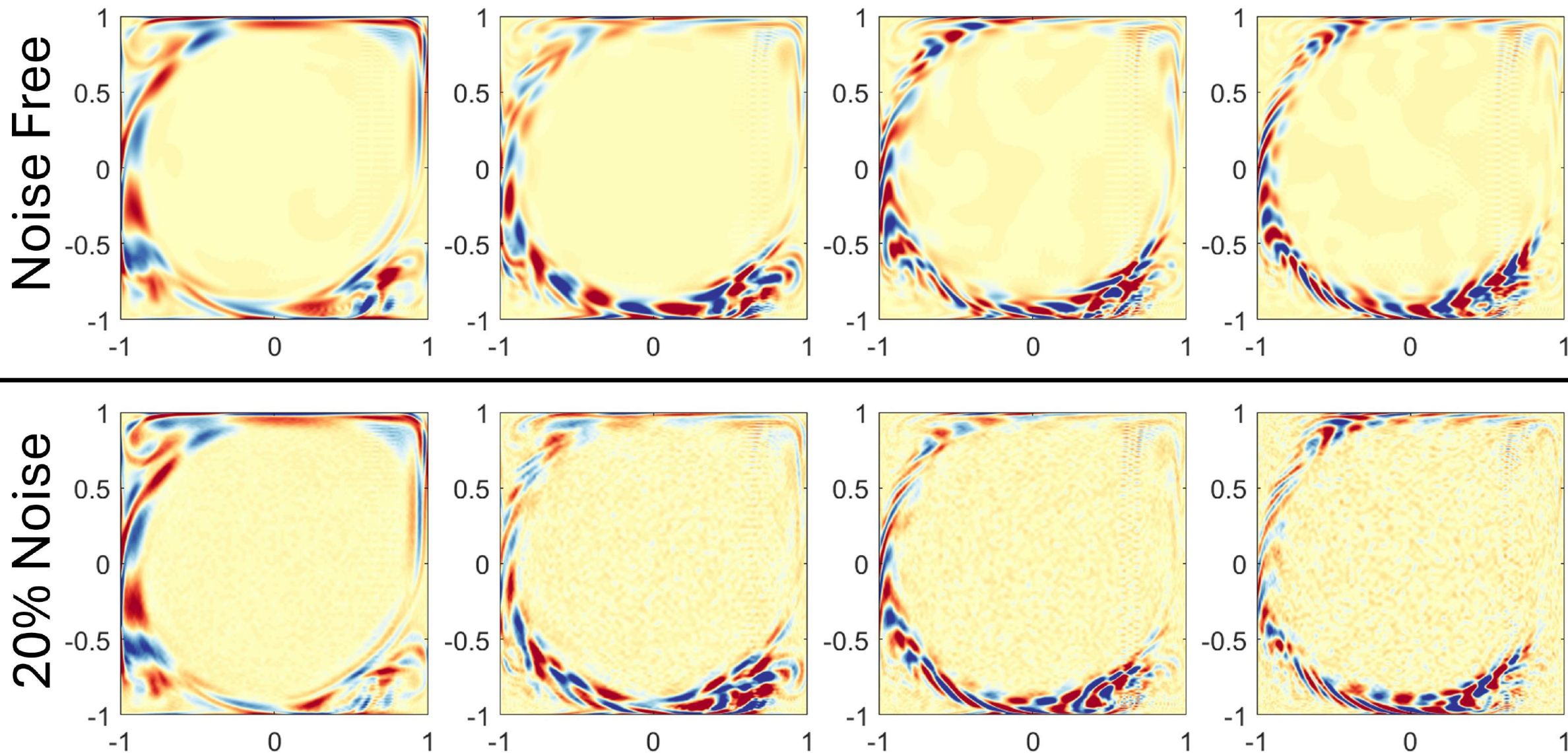


Example: Noisy cavity flow (generalized Koopman modes)

Re=30000

52

Deep in the continuous spectrum!!!



Practical + dictionary agnostic
+ theoretical guarantees

Summary

Interest in Koopman boils down to a data-driven inf-dim spectral problem.

• mpEDMD

- EDMD + enforcing measure-preserving (polar decomposition of Galerkin)
- Convergence of spectral measures, spectra, Koopman mode decomposition.
- Long-time stability, improved qualitative behavior, increased stability to noise.

• Rigged DMD

- Continuous spectra and generalized eigenfunctions.
- Smoothing kernels + resolvent (using mpEDMD).
- High-order convergence.

Future work

- Use in control.
- Other function spaces? E.g., RKHS

[General (non-measure-preserving) systems: ResDMD]

Brief Summaries

Journal of the Society for Industrial and Applied Mathematics siamnews.org
siam news Volume 56 Issue 1
 January/February 2023

Resilient Data-driven Dynamical Systems with Koopman: An Infinite-dimensional Numerical Analysis Perspective
 By Steven L. Brunton and Matthew J. Colbrook

on the local analysis of...
 D...
 Consider a discrete-time dynamical system with state x in a state space $\Omega \subset \mathbb{R}^n$ that is governed by an unknown and typically nonlinear function $F: \Omega \rightarrow \Omega$.

$x_{n+1} = F(x_n), \quad n \geq 0. \quad (1)$

The classical, geometric way to analyse such systems—which dates back to the seminal work of Henri Poincaré—is based

Measure-preserving Extended Dynamic Mode Decomposition

Residual Dynamic Mode Decomposition

YouTube

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